# On the boundary values of harmonic functions in $\mathbf{R}^{3}$ 

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## 1. Introduction

The purpose of this paper is to exhibit three theorems about the boundary values of harmonic functions, defined in $R^{3}$ regions which are bounded by Liapunov surfaces. Theorem 1 shows the existence almost everywhere of nontangential boundary values of positive harmonic functions. A full proof of this theorem is given. Theorem 2 assures the existence almost everywhere of nontangential boundary values for functions bounded in cones with vertex on the boundary and lying in the region. Theorem 3, finally, gives a necessary and sufficient condition for the existence of non-tangential boundary values, originally derived by Marcinkiewicz and Zygmund and later generalized by Stein. As the proofs of the two latter theorems differ from proofs published elsewhere only in the technical aspect, these are not included here.

## 2. Definitions

We consider an open region $\Omega_{1}$, bounded by a Liapunov surface $S_{1}$. By Liapunov surface we mean a closed, bounded surface with the following properties:
$1^{0}$. At every point of $S_{1}$ there exists a uniquely defined tangent plane, and thus also a normal.
$2^{\circ}$. There exist two constants $C^{\prime}>0$ and $\lambda, 0<\lambda \leqslant 1$, such that if $\theta$ is the angle between two normals, and $r$ is the distance between their foot points, the following inequality holds $\theta<C^{\prime} \cdot r^{\lambda}$.
$3^{\circ}$. There exists a constant $d>0$ such that if $\sum$ is a sphere with radius $d$ and center $Q_{0}$ on the surface, a line parallel to the normal at $Q_{0}$ meets $S_{1}$ at most once inside $\Sigma$. It is easily realized that $d$ may be chosen arbitrarily small.

For the properties of Liapunov surfaces see Gunther [5]. In the sequel we shall consider only inner normals, which will simply be referred to as normals. We denote by $V(Q, \alpha, h)$ a right circular cone having vertex at $Q \in S_{1}$, axis along the normal at $Q$, altitude $h$, generating angle $=\alpha$ and being contained in $\Omega_{1}$. Non-tangential approach to the boundary means approach inside some $V(Q, \alpha, h)$. $r(P, Q)$ will be the distance between $P \in \Omega_{1}$ and the tangent plane at $Q \in S_{1}$. The volume element in $R^{3}$ will be denoted by $d v$ and the surface element by $d S$.
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## 3. Lemmata

If $Q \in S_{1}$ we introduce "the local coordinate system" with the origin at $Q$, with the $(x y)$-plane in the tangent plane at $Q$, and with the positive $z$-axis along the normal at $Q$. Inside the Liapunov sphere, $S_{1}$ may be represented on the form $z=f(x, y)$.

Lemma 1. If $C^{\prime} d^{\lambda}<1$ we have
and

$$
\begin{aligned}
& |f(x, y)| \leqslant 2 C^{\prime}\left(x^{2}+y^{2}\right)^{\frac{\lambda+1}{2}} \\
& |f(x, y)| \leqslant 2 C^{\prime}|P-Q|^{1+\lambda}
\end{aligned}
$$

where $P=(x, y, f(x, y))$.
Proof: Cf. Smirnov [7], p. 490.
Lemma 2. Let $\omega$ be $<\pi / 2$. Then there is a constant $d_{1}$ such that if $d<d_{1}$, a line making an angle $\leqslant \omega$ with the normal at an arbitrary point $Q \in S_{1}$, will meet $S_{1}$ at most once inside the sphere with center $Q$ and radius d.

Proof: See Gunther [5], p. 6.
Denote by $G(P, Q)$ the Green's function of $\Omega_{1}$.
Lemma 3. There is a constant $c_{1}$, depending on $\Omega_{1}$ only, such that for any points $P, Q$ we have

$$
\left|\frac{\partial G}{\partial x_{\theta}}(P, Q)\right| \leqslant \frac{c_{1}}{|P-Q|^{2}},
$$

where $\left(\partial G / \partial x_{\theta}\right)(P, Q)$ denotes differentiation in an arbitrary direction at $P$. In particular, for $P \in S_{1}$ we have

$$
\left|\frac{\partial G}{\partial n_{P}}(P, Q)\right| \leqslant \frac{c_{1}}{|P-Q|^{2}},
$$

$\partial G / \partial n_{P}$ denoting the normal derivative at $P$.
Proof: See Eidus [3].
Lemma 4. The derivatives of $G(P, Q)$ are continuous in $\Omega_{1} \cup S_{1}$.
Proof: According to Schauder [6], a harmonic function in $\Omega_{1}$, having tangential derivatives along the boundary which satisfy a Hölder condition, has derivatives of the first order that are continuous in $\Omega_{1} \cup S_{1}$. Thus it suffices to prove that $1 /|P-Q|, Q \in \Omega_{1}, P \in S_{1}$, has tangential derivatives of the abovementioned kind. But that is clear, because the tangential derivatives may be expressed as a linear combination of the partial derivatives of $1 /|P-Q|$ with functions which themselves satisfy Hölder conditions, as coefficients.

Lemma 5. Let $\omega$ be $<\pi / 2$. Then there exist two constants $\delta>0$ and $c_{13}>0$ such that if $Q \in S_{1}, B \in \Omega_{1},|B-Q|<\delta$ and if the angle between $\overrightarrow{Q B}$ and the normal at $Q$ is $\leqslant \omega$ we have

$$
\frac{\partial G}{\partial n_{Q}}(Q, B) \geqslant \frac{c_{13}}{|Q-B|^{2}} .
$$

Proof: It suffices to prove the lemma in the case $\omega=0$. A simple application of Harnack's inequality then gives the general case. We may write (cf. Gunther [5], p. 202)

$$
G\left(P, P^{\prime}\right)=\frac{1}{\left|P-P^{\prime}\right|}-\frac{1}{4 \pi} \int_{S_{1}} \frac{1}{|P-T|} \cdot \frac{\partial G}{\partial n_{T}}\left(T, P^{\prime}\right) d S(T)
$$

As we may regard the integral as the potential of a simple layer, we get (cf. Gunther [5], p. 62):

$$
\frac{\partial G}{\partial n_{Q}}(Q, B)=\frac{1}{|Q-B|^{2}}-\frac{1}{4 \pi} \int_{S_{1}} \frac{\cos \varphi}{|Q-T|^{2}} \cdot \frac{\partial G}{\partial n_{T}}(T, B) d S(T)+\frac{1}{2} \frac{\partial G}{\partial n_{Q}}(Q, B)
$$

i.e. $\quad \frac{\partial G}{\partial n_{Q}}(Q, B)=\frac{2}{|Q-B|^{2}}-\frac{1}{2 \pi} \int_{S_{1}} \frac{\cos \varphi}{|Q-T|^{2}} \frac{\partial G}{\partial n_{T}}(T, B) d S(T)$,
where $\varphi$ is the angle between $\vec{Q} \vec{T}$ and the normal at $Q$. Choose $\sigma>0$ so that $32 c_{1} \cdot C^{\prime} \cdot \sigma^{\lambda}<1$ and $\sigma<d$. If $T=(\xi, \eta, \zeta)$ we have by Lemma 1, when $\xi^{2}+\eta^{2} \leqslant \sigma^{2}$,

$$
\zeta<2 C^{\prime}|Q-T|^{1+\lambda}
$$

Moreover, we have $\cos \varphi=\frac{\zeta}{|Q-T|}$, which gives us

$$
|\cos \varphi| \leqslant 2 C^{\prime}|Q-T|^{\lambda}
$$

For $\delta$ small enough we evidently have

$$
|B-T| \geqslant \frac{1}{\sqrt{2}}|B-Q|
$$

Hence

$$
\begin{aligned}
\left|\int_{|Q-T| \leqslant \sigma} \frac{\cos \varphi}{|Q-T|^{2}} \frac{\partial G}{\partial n_{T}}(T, \dot{B}) d S(T)\right| & \leqslant \frac{4 c_{1} \cdot C^{\prime}}{|B-Q|^{2}} \int_{|Q-T| \leqslant \sigma}|Q-T|^{2-2} d S(T) \\
& \leqslant \frac{16 c_{1} \cdot C^{\prime}}{|B-Q|^{2}}{ }^{\lambda} \cdot 2 \pi<\frac{\pi}{|B-Q|^{2}}
\end{aligned}
$$

Having fixed $\sigma$ we find

$$
\left|\int_{|Q-T| \geqslant \sigma} \frac{\cos \varphi}{|Q-T|^{2}} \frac{\partial G}{\partial n_{T}}(T, B) d S(T)\right| \leqslant \frac{1}{\sigma^{2}} \cdot 4 \pi .
$$

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If now $\delta$ is chosen small enough, we get

$$
\frac{1}{\sigma^{2}} \cdot 4 \pi<\frac{2 \pi}{2 \delta^{2}}
$$

Thus

$$
\frac{\partial G}{\partial n_{Q}}(Q, B)=\frac{2}{|Q-B|^{2}}-\frac{1}{2 \pi} \int_{S_{i}} \frac{\cos \varphi}{|Q-T|^{2}} \frac{\partial G}{\partial n_{T}}(T, B) d S(T) \geqslant \frac{1}{|Q-B|^{2}}
$$

Lemma 6. Let $P_{0}$ be a fixed point of $\Omega_{1}$. There exists a constant $c_{2}>0$ such that

$$
\frac{\partial G}{\partial n_{Q}}\left(Q, P_{0}\right) \geqslant c_{2}
$$

for all $Q \in S_{1}$.
Proof: Let $\delta$ be as in Lemma 5. We may assume that $\left|P_{0}-Q\right|>\delta$ for all $Q \in S_{1}$. The set of points $P$, belonging to $\Omega_{1}$, and for which $|P-Q| \geqslant \delta / 2$ when $Q \in S_{1}$, is compact and contains $P_{0}$. Harnack's inequality gives

$$
\frac{1}{c_{3}} u\left(P_{0}\right) \leqslant u(P) \leqslant c_{3} \cdot u\left(P_{0}\right)
$$

for every positive harmonic function and every $P$ belonging to the set. Let $Q$ be an arbitrary point $\in S_{1} .\left(\partial G / \partial n_{Q}\right)(Q, P)$ is harmonic and positive. Moreover, there is a point $P$ on the normal at $Q$, for which $\delta / 2 \leqslant|P-Q| \leqslant \delta$, and by Lemma 5 we get

$$
\frac{\partial G}{\partial n_{Q}}\left(Q, P_{0}\right) \geqslant \frac{1}{c_{3}} \cdot \frac{\partial G}{\partial n_{Q}}(Q, P) \geqslant \frac{c_{13} \cdot 4}{c_{3} \cdot \delta^{2}}=c_{2}>0 .
$$

Lemma 7. If $Q$ and $Q^{\prime} \in S_{1},\left|Q-Q^{\prime}\right| \geqslant \varrho>0$,

$$
\lim _{P \rightarrow Q^{\cdot}} \frac{\partial G}{\partial n_{Q}}(Q, P)=0
$$

uniformly in $Q$ and $Q^{\prime}$.
Proof: As in the proof of Lemma 5 we have

$$
\frac{\partial G}{\partial n_{Q}}(Q, P)=\frac{\partial}{\partial n_{Q}} \frac{2}{|Q-P|}-\frac{1}{2 \pi} \int_{S_{1}} \frac{\cos \varphi}{|Q-T|^{2}} \frac{\partial G}{\partial n_{T}}(T, P) d S(T)
$$

Choose an arbitrary $\varepsilon>0$. It is easily realized, that $\left(\partial / \partial n_{Q}\right)(1 /|Q-P|)$ is uniformly continuous in $\left\{P||Q-P| \geqslant \varrho / 2\}\right.$. Choose $\sigma_{1}, 0<\sigma_{1}<\varrho / 3$, so small that

$$
\left|\frac{\partial}{\partial n_{Q}} \frac{1}{|Q-P|}-\frac{\partial}{\partial n_{Q}} \frac{1}{\left|Q-P^{\prime}\right|}\right|<\frac{\varepsilon}{6}
$$

if $\left|P-P^{\prime}\right|<2 \sigma_{1}$. Finally choose $\sigma_{2}, 0<\sigma_{2}<\varrho / 3$ so that

$$
\frac{144 \cdot c_{1} \cdot C^{\prime} \cdot \pi}{\varrho^{2}} \sigma_{2}^{\lambda}<\frac{\varepsilon}{3}
$$

For $\left|P-Q^{\prime}\right|<\sigma_{1} / 2$ we now have

$$
\int_{\mid T-Q} \left\lvert\, \geqslant \sigma_{1} \frac{\partial G}{\partial n_{T}}(T, P) d S(T) \leqslant A \cdot \frac{\left|P-Q^{\prime}\right|^{\lambda}}{\sigma_{1}^{2}}\right.
$$

where $A$ is a constant not depending on $Q^{\prime}$ and $\sigma_{1}$. In fact, this inequality is a lemma of Liapunov's. It is proved in Gunther $|5|$, p. 200 for the case when $P$ is on the normal at $Q^{\prime}$, but the proof is easily extended to the more general case. Hence we get

$$
\int_{\left\lvert\, \begin{array}{l}
T-Q \\
T-Q \mid \geqslant \sigma_{2} \\
|Q-T|^{2}
\end{array}\right.} \frac{\cos \varphi}{\partial n_{T}}(T, P) d S(T) \leqslant A \cdot\left|P-Q^{\prime}\right|^{2} \cdot \frac{1}{\sigma_{1}^{2}} \cdot \frac{1}{\sigma_{2}^{2}} .
$$

As in Lemma 5 we also get

$$
\int_{|T-Q| \leqslant \sigma_{\mathrm{e}}} \frac{\cos \varphi}{|Q-T|^{2}} \frac{\partial G}{\partial n_{T}}(T, P) d S(T) \leqslant \frac{9 \cdot c_{1}}{\varrho^{2}} \cdot 16 C^{\prime} \pi \cdot \sigma_{2}^{\lambda}<\frac{\varepsilon}{3} .
$$

Finally

$$
\begin{aligned}
\left\lvert\, 2 \frac{\partial}{\partial n_{Q}} \frac{1}{|Q-P|}\right. & \left.-\frac{1}{2 \pi} \int_{\left|T-Q^{\prime}\right| \leqslant \sigma_{1}} \cdot \frac{\partial}{\partial n_{Q}} \frac{1}{|Q-T|} \frac{\partial G}{\partial n_{T}}(T, P) d S(T) \right\rvert\, \\
& \leqslant \frac{1}{2 \pi} \int_{\left|T-Q^{\prime}\right| \leqslant \sigma_{1} \mid}\left|\frac{\partial}{\partial n_{Q}} \frac{1}{|Q-P|}-\frac{\partial}{\partial n_{Q}} \frac{1}{|Q-T|}\right| \frac{\partial G}{\partial n_{T}}(T, P) d S(T) \\
& +\frac{\partial}{\partial n_{Q}} \left\lvert\, \frac{1}{|Q-P|} \cdot \frac{1}{2 \pi} \int_{\left|T-Q^{\prime}\right| \geqslant \sigma_{1}} \frac{\partial G}{\partial n_{T}}(T, P) d S(T) \leqslant \frac{\varepsilon}{3}+\frac{4}{\varrho^{2}} \cdot A \cdot \frac{\left|P-Q^{\prime}\right|^{\lambda}}{\sigma_{1}^{2}}\right.
\end{aligned}
$$

provided that $\left|P-Q^{\prime}\right|<\sigma_{1} / 2$.
If $P$ is chosen so close to $Q^{\prime}$ that $A \cdot\left|P-Q^{\prime}\right|^{1} \frac{1}{\sigma_{1}^{2}} \cdot \frac{1}{\sigma_{2}^{2}}<\frac{\varepsilon}{6}$, the lemma is proved by inserting the resulting estimates in the integral representation of ( $\partial G / \partial n_{Q}$ ) $(Q, P)$.

## 4. Theorem 1

If $u_{1}$ is a positive harmonic function in $\Omega_{1}$, then $u_{1}$ has non-tangential boundary values almost everywhere on $S_{1}$. Moreover, the boundary values $\in L^{1}\left(S_{1}\right)$.

Remark. Fatou proved in 1906 that a bounded harmonic function in the unit disk has non-tangential boundary values a.e. [4]. Using conformal mapping, Fatou's theorem may be stated for quite general plane regions. Calderon [1] and Carleson [2]

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have proved similar theorems for regions in $R^{n}$ bounded by planes, but have weaker assumptions on the function. Tsuji [10] and Solomencev [8] have studied regions bounded by curved surfaces, but are assuming stronger regularity conditions than those in this paper. The method of proof used here is strongly related to the paper of prof. Carleson, to whom the author also owes his interest in this problem.

Proof: Let $Q^{\prime}$ be an arbitrary point $\in S_{1}$. It suffices to prove the existence of boundary values in a neighbourhood of $Q^{\prime}$. Assume $d$ to be so small that
$1^{\circ}$. $d<d_{1}=$ the number appearing in Lemma 2 , if $\omega=\operatorname{arctg} 1 / 2$.
$2^{\circ} .8 C^{\prime} d^{\lambda}<1$, i.e. $2 C^{\prime} d^{1+\lambda}<d / 4$.
We introduce the local coordinate system with the origin at $Q^{\prime}$. Let $O=(0,0$, $3 d / 4)$. The condition $2^{\circ}$ shows that a line through $O$ and a point on $S_{1}$ inside the sphere $\sum$ with center $Q^{\prime}$ and radius $d$ will make an angle with the $z$-axis which is $<\operatorname{arctg} 1 / 2=\omega$. Condition $1^{\circ}$ and Lemma 2 show that such a line meets $S_{1}$ at most once inside $\Sigma$. We now introduce polar coordinates with origin = $O, \varphi$ being the angle between radius vector and the plane through $O$ parallel with the $(x, y)$-plane. What we just proved shows that the part of $S_{1}$ which is inside $\Sigma$ may be represented by $\varrho=\varrho_{1}(\theta, \varphi)$. Let $S$ be the part of $S_{1}$ which is inside the cylinder, which in the local system has the equation $x^{2}+y^{2}=d^{2} / 4$. The surface of this cylinder may be represented in polar coordinates: $\varrho=\varrho_{2}(\theta, \varphi)$. The values of the vector $(\theta, \varphi)$ for which $\left(\theta, \varphi, \varrho_{1}(\theta, \varphi)\right) \in S$ will be called " $(\theta, \varphi)$ values belonging to $S^{\prime \prime}$. Evidently, $(\theta,-(\pi / 4))$ do not belong to $S$. Now let $\Psi(\theta, \varphi)$ be an infinitely differentiable function, such that $0 \leqslant \Psi \leqslant 1, \Psi(\theta, \varphi)=0$ for $\varphi \geqslant-(\pi / 4)$, and $\Psi(\theta, \varphi)=1$ for $(\theta, \varphi)$ belonging to $S$. The existence of such a function is easily proved. The surface generated by

$$
\varrho=\varrho(\theta, \varphi)=\varrho_{\mathbf{1}}(\theta, \varphi) \Psi(\theta, \varphi)+\varrho_{\mathbf{2}}(\theta, \varphi)[1-\Psi(\theta, \varphi)]
$$

satisfies the conditions of a Liapunov surface, apart from not being bounded. If we insert the halfsphere with radius $d / 2$ and center ( $0,0, d / 2$ ) however, we get a new region $\Omega$, bounded by a Liapunov surface. We denote by $\Gamma$ the part of $\partial \Omega$ which is not $S$ and assume that the normals of $S \cup \Gamma$ satisfy $\theta \leqslant C \cdot r^{\lambda}$.

Consider the restriction $u$ (as a function of the local coordinates) of $u_{1}$ to $\Omega$. Of course, $u$ is harmonic and positive. Put $u_{\varepsilon}(x, y, z)=u(x, y, z+\varepsilon)$. For $\varepsilon$ small enough this function is defined in $\Omega \cup S \cup \Gamma$, and thus it may be represented as an integral (cf. Gunther [5], p. 202)

$$
u_{\varepsilon}(P)=\frac{1}{4 \pi} \int_{\Gamma U S} u_{\varepsilon}(Q) \frac{\partial G}{\partial n_{Q}}(Q, P) d S(Q) .
$$

The measures $d \mu_{\varepsilon}(Q)=u_{\varepsilon}(Q) d S(Q)$ have compact support and uniformly bounded total mass. For by Lemma 6 and the fact that $u$ is positive, we get

$$
\int_{S \cup \Gamma}\left|d \mu_{\varepsilon}\right|=\int u_{\varepsilon}(\theta) d S(Q) \leqslant \frac{1}{c_{2}} \int \frac{\partial G}{\partial n_{Q}}(Q, O) u_{\epsilon}(Q) d S(Q)=\frac{4 \pi}{c_{2}} u_{\epsilon}(O) .
$$

But $u_{\varepsilon}(0) \rightarrow u(0)$ and thus the total mass is uniformly bounded. Hence we may
choose a weakly convergent subsequence, $d \mu_{\varepsilon} \rightarrow d \mu$. A well-known theorem of Lebesgue allows us to decompose $d \mu$ :

$$
d \mu(Q)=f(Q) d S(Q)+d \sigma(Q)
$$

where $f(Q) \in L^{1}(S \cup \Gamma)$, and $d \sigma$ is singular. This gives us

$$
u(P)=\frac{1}{4 \pi} \int_{\Gamma \cup S} \frac{\partial G}{\partial n_{Q}}(Q, P)[f(Q) d S(Q)+d \sigma(Q)] .
$$

We assert that $u(P) \rightarrow f(Q)$ when $P \rightarrow Q$ non-tangentially, a.e. on $S$. By another well-known theorem of Lebesgue we have for almost all $T \in S$ :

$$
\int_{|Q-T|<\varepsilon}\{|f(Q)-f(T)| d S(Q)+d \sigma(Q)\}=o\left(\varepsilon^{2}\right)
$$

Assume $Q_{0}$ to be such a point and choose an arbitrary $\alpha<\pi / 2$. Let $h$ be so small that $V\left(Q_{0}, \alpha / 2+\pi / 4, h\right)$ is contained in $\Omega$. Suppose now that $\delta>0$ is chosen so that
$1^{\circ} .4 \delta<d$ and $4 \delta<\inf \left|Q_{0}-Q\right|, Q \in \Gamma$.
$2^{\circ}$. Lemmá 5 is valid for $\omega=\pi / 4$.
$3^{\circ}$. $\sqrt{1+\frac{9}{4}-3 \cos 2 \beta}<\frac{5}{8}$ and $\sqrt{1+\frac{10}{9}-\frac{8}{3} \cos 2 \beta}<\frac{5}{12}$ where $\beta=\operatorname{arctg}\left[2 C \delta^{\lambda}\right]$.
$4^{\text {a }} C \delta^{\lambda}<\frac{\sqrt{15}}{14}<\frac{\pi}{8}$.
We assume $A$ to be an arbitrary point $\in V\left(Q_{0}, \alpha, h\right)$, such that $\left|A-Q_{0}\right|=a<\delta / 2$. Define

$$
\begin{gathered}
L_{0}=\left\{P\left|P \in S,\left|Q_{0}-P\right|<a\right\} .\right. \\
L_{v}=\left\{P\left|P \in S, 2^{v-1} a \leqslant\left|Q_{0}-P\right|<2^{y} a\right\}\right.
\end{gathered}
$$

for $\nu=1,2, \ldots, N$ where $2^{N} a \leqslant \delta / 2<2^{N+1} a$ and

$$
L=\{\Gamma \cup S\}-\left\{\bigcup_{v=0}^{N} L_{v}\right\}
$$

As we can write

$$
\begin{gathered}
f\left(Q_{0}\right)=\frac{1}{4 \pi} \int_{S \cup \Gamma} f\left(Q_{0}\right) \frac{\partial G}{\partial n_{Q}}(Q, A) d S(Q) \text { we get } \\
\left|u(A)-f\left(Q_{0}\right)\right| \leqslant \frac{1}{4 \pi} \int_{S \cup \Gamma}\left|f(Q)-f\left(Q_{0}\right)\right| \frac{\partial G}{\partial n_{Q}}(Q, A) d S(Q)+\frac{1}{4 \pi} \int_{S \cup \Gamma} \frac{\partial G}{\partial n_{Q}}(Q, A) d \sigma(Q) \\
\leqslant \frac{1}{4 \pi} \sum_{v=0}^{N}\left\{\sup _{Q \in L_{v}} \frac{\partial G}{\partial n_{Q}}(Q, A)\right\} \int_{i=0}^{v}\left\{\left|f(Q)-f\left(Q_{0}\right)\right| d S(Q)+d \sigma(Q)\right\}
\end{gathered}
$$

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$$
\begin{align*}
& +O(1) \sup _{Q \in L} \frac{\partial G}{\partial n_{Q}}(Q, A)=\varepsilon\left(\frac{\delta}{2}\right) \sum_{v=0}^{N}\left\{\sup _{Q \in L \nu} \frac{\partial G}{\partial n_{Q}}(Q, A)\right\} 2^{2 v} \cdot a^{2} \\
& +O(1) \sup _{Q \in L} \frac{\partial n_{Q}}{G \partial}(Q, A) . \tag{1}
\end{align*}
$$

To simplify the notations we put $\left(\partial G / \partial n_{Q}\right)(Q, P)=K(Q, P)$. In the sequel $c_{i}$ denote constants not depending on $a, v$ or $\delta$. We need estimates of $K(Q, A)$ for $Q \in L_{v}$. For the two cases $\nu=0$ and $\nu=1$ Lemma 5 gives

$$
\begin{equation*}
K(Q, A) \leqslant \frac{c_{1}}{|Q-A|^{2}} \leqslant \frac{c_{1}}{a^{2} \sin ^{2}\left(\frac{\pi}{4}-\frac{\alpha}{2}\right)}=\frac{c_{4}}{a^{2}} . \tag{2}
\end{equation*}
$$

Choose $\nu, 2<\nu \leqslant N$, let $Q_{\nu}$ be an arbitrary point $\in L_{\nu}$, and let $B$ be the point on the normal at $Q_{\nu}$, for which we have $\left|Q_{\nu}-B\right|=2^{v} a$. As $2^{v} a \leqslant 2^{N} a<\delta<d / 4$, $B$ is in $\Omega$. It follows from Lemma 3 that

$$
\begin{equation*}
K\left(Q_{\nu}, B\right) \leqslant \frac{c_{1}}{2^{2 \nu} \cdot a^{2}} \tag{3}
\end{equation*}
$$

We denote by the ( $\xi \eta \zeta$ )-system the local coordinate system with the origin at $Q_{0}$. Let $M_{v}$ be the point of $L_{v}$, whose projection along the $\zeta$-axis on the $(\xi, \eta)$ plane is on the same line through $Q_{0}$ as $Q_{\nu}$ and which lies at a distance $=$ $3 \cdot 2^{\nu-2} a$ from $Q_{0}$. Letting $\Gamma_{1}$ be the sphere with center $M_{v}$ and radius $3 \cdot 2^{y-3} a$, we find that $B$ does not lie inside $\Gamma_{1}$. Now there exists a constant $c_{5}$ such that $K\left(Q_{v}, P\right) \leqslant c_{5} \cdot K\left(Q_{v}, B\right)$ for $P \notin \Gamma_{1}$. For by Lemma 5 we have

$$
\begin{equation*}
K\left(Q_{\nu}, B\right) \geqslant \frac{c_{13}}{\left|Q_{\nu}-B\right|^{2}}=c_{13} \cdot 2^{-2 v} \cdot a^{-2} \tag{4}
\end{equation*}
$$

From assumption $3^{\circ}$ about $\delta$ it follows that $\left|P-Q_{\nu}\right| \geqslant 2^{\nu-4} a$ for $P \oplus \Gamma_{1}$, and from this fact and Lemma 3 it follows that

$$
\begin{equation*}
K\left(Q_{v}, P\right) \leqslant \frac{c_{1}}{\left|Q_{\nu}-P\right|^{2}} \leqslant c_{1} \cdot 2^{-2 \nu+8} \cdot a^{-2} \tag{5}
\end{equation*}
$$

(4) and (5) give

$$
\begin{equation*}
\frac{K\left(Q_{v}, P\right)}{K\left(Q_{v}, B\right)} \leqslant \frac{c_{1} a^{-2 v+8} \cdot a^{-2}}{c_{13} \cdot 2^{-2 v} \cdot a^{-2}}=\frac{c_{1} \cdot 2^{8}}{c_{13}}=c_{5} \tag{6}
\end{equation*}
$$

Denoting by $h_{\nu}(P)$ the harmonic measure of $L_{v-1} \cup L_{v} \cup L_{v+1}$ we find for $P \in \partial \Gamma_{1} \cap \Omega$

$$
\begin{equation*}
h_{p}(P) \geqslant c_{B}>0 \tag{7}
\end{equation*}
$$

For introduce the $(X, Y, Z)$-system $=$ the local coordinate system with the origin at $M_{v}$. The part of $S$, for which $X^{2}+Y^{2} \leqslant\left(7 \cdot 2^{\nu-4} a\right)^{2}$ is contained in $L_{v-1} \cup L_{v} \cup L_{v+1}$.

This follows from assumption $4^{\circ}$ about $\delta$; if $Q$ is on said part of $S$ we have $\left|Q-M_{\nu}\right|^{2} \leqslant\left(7 \cdot 2^{\nu-4} \cdot a\right)^{2}+\left(2 C\left[7 \cdot 2^{\nu-4} a\right]^{1+\lambda}\right)^{2}<2^{\nu-1} a$, while for $Q \in S$ but $Q \notin L_{\nu-1}$ $\cup L_{v} \cup L_{v+1}$ we evidently have $\left|Q-M_{\nu}\right| \geqslant 2^{v-1} a$. If $P$ is an arbitrary point of $\partial \Gamma_{1} \cap \Omega$, we find for $P \in S h_{p}(P)=1$, and hence we may assume that $P \ddagger S$. Put $P^{\prime}=$ the projection of $P$ on the $(X, Y)$-plane and $\overrightarrow{P^{\prime} P}=l,\left|P^{\prime}-P\right|=|l|$. We proved above that

$$
U=\left\{Q|Q \in S, Q=(X, Y, Z) \Rightarrow|(X, Y, 0)-P^{\prime} \left\lvert\, \leqslant \frac{|l|}{12}\right.\right\} \subset L_{\nu-1} \cup L_{\nu} \cup L_{v+1}
$$

If $P^{\prime \prime} \in U$ the angle between $\overrightarrow{P^{\prime \prime}} \vec{P}$ and $l$ is $<\pi / 8$, and from assumption $4^{\circ}$ about $\delta$ it follows that the angle between $l$ and the normal at $P^{\prime \prime}$ is $<\pi / 8$. Hence the angle between $\overrightarrow{P^{\prime \prime} P}$ and the normal at $P^{\prime \prime}$ is $<\pi / 4$. Lemma 5 and assumption $2^{\circ}$ about $\delta$ then give

$$
\frac{\partial G}{\partial n_{P^{\prime \prime}}}\left(P^{\prime \prime}, P\right) \geqslant \frac{c_{13}}{16|l|^{2}} .
$$

Thus

$$
h_{\nu}(P)=\frac{1}{4 \pi} \int_{L_{p-1} \cup L_{v} \cup L_{p+1}} \frac{\partial G}{\partial n_{Q}}(Q, P) d S(Q) \geqslant \frac{c_{13}}{4 \pi \cdot 16|l|^{2}} \cdot \frac{\pi \cdot|l|^{2}}{144}=c_{6}>0 .
$$

For $P \in \partial \Gamma \cap \Omega$ it follows from (6) and (7)

$$
\frac{K\left(Q_{v}, P\right)}{K\left(Q_{v}, B\right)} \leqslant c_{5}=\frac{c_{5}}{c_{6}} \cdot c_{6} \leqslant c_{7} \cdot h_{\nu}(P)
$$

If $Q$ is any other boundary point of $\Omega \backslash \Gamma_{1}$, we have $\lim _{P \rightarrow Q} h_{\nu}(P) \geqslant 0$, while by Lemma 7 we have

$$
\lim _{P \rightarrow Q} \frac{K\left(Q_{\nu}, P\right)}{K\left(Q_{v}, B\right)}=0
$$

Together with (3), the maximum principle then gives for $P \in \Omega \backslash \Gamma_{1}$

$$
K\left(Q_{v}, P\right) \leqslant c_{7} \cdot h_{\nu}(P) \cdot K\left(Q_{v}, B\right) \leqslant c_{7} \cdot h_{\nu}(P) \cdot c_{1} \cdot 2^{-2 \nu} \cdot a^{-2}
$$

Particularly we get for $P=A$

$$
\begin{equation*}
K\left(Q_{\nu}, A\right) \leqslant c_{7} \cdot h_{\nu}(A) \cdot c_{1} \cdot 2^{-2 \nu} \cdot a^{-2} \tag{8}
\end{equation*}
$$

Inserting (2) and (8) in (1) we find

$$
\left|u(A)-f\left(Q_{0}\right)\right| \leqslant c_{8} \cdot \varepsilon\left(\frac{\delta}{2}\right)+\varepsilon\left(\frac{\delta}{2}\right) \cdot c_{9} \cdot \sum_{\nu=2}^{N} h_{\vartheta}(A)+O(1) \cdot \sup _{Q \in L} \frac{\partial G}{\partial n_{Q}}(Q, A) .
$$

Now $\sum_{y=2}^{N} h_{p}(P) \leqslant 3$, as we may regard $\sum_{v=2}^{N} h_{p}(P)$ as the harmonic measure of a certain subset of $S$, each point being counted at most three times. From Lemma 7 it follows

$$
\varlimsup_{A \rightarrow Q_{0}}\left\{\sup _{Q \in L} \frac{\partial G}{\partial n_{Q}}(Q, A)\right\}=0
$$

Hence

$$
\varlimsup_{A \rightarrow Q_{0}}\left|u(A)-f\left(Q_{0}\right)\right| \leqslant c_{10} \cdot \varepsilon\left(\frac{\delta}{2}\right) .
$$

But $\varepsilon(\delta / 2) \rightarrow 0$ when $\delta \rightarrow 0$, and thus the first part of the theorem is proved.
We remarked above that the boundary values $\in L^{1}(S \cup \Gamma)$ and thus also $\in L^{1}(S)$. The latter part of the theorem now follows from the simple fact that we can cover $S_{1}$ with a finite number of such neighbourhoods $S$.

## 5. Theorem 2

If $u$ is harmonic in $\Omega_{1}$, and if for almost all $Q$ there is a cone $V(Q, \alpha, h)$ in which $u$ is bounded, then $u$ has non-tangential boundary values almost everywhere on $S_{1}$.

Remark. The analogue of this theorem in the case $S_{1}$ is a plane was proved by Calderon [1]. Having proved Theorem 1, Theorem 2 follows as in Calderon's paper, his proof needing only minor modifications to be applicable in this case. We therefore omit the proof.

## 6. Theorem 3

The harmonic function $u$ in $\Omega_{1}$ has non-tangential boundary values if and only if

$$
\int_{V(Q, \alpha, h)} \frac{|\operatorname{grad} u(P)|^{2}}{r(P, Q)} d v(P)<\infty
$$

almost everywhere in $Q$.
Remark. The values of $\alpha$ and $h$ may depend on $Q$. The formulation of this theorem is chosen because of its resemblance to a similar theorem by Stein [9], valid when $S_{1}$ is a plane. Note that in this case the theorem of Carleson [2] contains that of Stein. The proof of Theorem 3 in principle follows Stein's proof, the difference being of a technical kind. However, the technical difficulties are overcome by the use of the methods developed in the lemmas and in Theorem 1. Hence we here are content with the mere statement of the theorem.

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