# Estimates of the age of a heat distribution 

By Johan Philip

## ABSTRACT


#### Abstract

The paper deals with the possibility to solve the heat equation backwards in time. More specifically, we treat the following problem. Given the temperature at a finite number of points of a homogeneous bar, how old can the heat distribution be? In the case that the temperature is given at equidistant points $x_{i}$, the problem is completely solved. In the case of nonequidistant $x_{i}$ we find an upper bound for the age. Such a bound is also obtained when the information about the heat distribution is given by the value of a finite number of linear functionals.


## I. Introduction

We consider the heat distribution (temperature distribution) in a homogeneous bar of infinite length (coordinate $x$ ) as a function of time ( $t$ ). Our heat distributions will be considered as positive measures $u_{t}(x)$.

The fundamental solution of the heat equation ( $\left.\partial^{2} u / \partial x^{2}=\partial u / \partial t\right)$ is

$$
\begin{equation*}
\psi_{t}(x)=(\mathbf{1} / 2 \sqrt{\pi t}) \exp \left(-x^{2} / 4 t\right) \quad(t>0) . \tag{1}
\end{equation*}
$$

An "initial heat distribution" $u_{0}$ at $t=0$ gives the following distribution at the time $t$ :

$$
\begin{equation*}
u_{i}=\psi_{t} * u_{\mathbf{0}} . \tag{2}
\end{equation*}
$$

We shall be concerned with problems connected with solving the heat equation backwards in time, viz. with the following problem: If $v$ is a bounded positive measure, for which $t$ does there exist a bounded positive measure $u_{0}$ satisfying

$$
\begin{equation*}
v=\psi_{t} * u_{0} ? \tag{3}
\end{equation*}
$$

When $t \rightarrow 0, \psi_{t}$ approaches the Dirac measure at the origin, so for $t=0$, (3) has the solution $u_{0}=v$. When $t \rightarrow \infty, \psi_{t} * u_{0} \rightarrow 0$ for every $x$, so (3) has no solution for large $t$ if $v \neq 0$. Furthermore, we have

$$
\begin{equation*}
\psi_{t_{1}} * \psi_{t_{3}}=\psi_{t_{1}+t_{2}} \tag{4}
\end{equation*}
$$

so if (3) has a solution $u_{0}$ for $t=\tau$, it has the solution $u_{0} * \psi_{\tau-x}$ for a time $x \leqslant \tau$.
Thus, it is meaningful to ask for the largest interval (0, $t$ ) in which (3) has a solution.

## J. Philip, Age of a heat distribution

In our problems, however, the information about $v$ will be incomplete and given by $n$ real numbers only. In sections II-IV the information is the values $v\left(x_{i}\right)$ at $n$ points $x_{i}$. In section $V$, we assume the values of $n$ linear functionals of $v$ to be known, which is a more realistic situation from a physical point of view.

## II. Formulation of Problem I

Let the information about $v$ be given by the values $a_{i}=v\left(x_{i}\right)$ at $x_{1}<x_{2}<x_{3}<\ldots<x_{n}$, so that we have the equations

$$
\begin{equation*}
a_{i}=\psi_{t} * u_{0}\left(x_{i}\right) \quad(1 \leqslant i \leqslant n) . \tag{5}
\end{equation*}
$$

If $u_{0} \equiv 0$ we have $a_{i}=0(1 \leqslant i \leqslant n)$. If $u_{0} \neq 0$ and $t>0$ we have $a_{i}>0,(1 \leqslant i \leqslant n)$, since $\psi_{t}>0$ for $t>0$. Thus, if $a_{i}=0$ for some $i$ but not for all we must have $t=0$. Further, for $n=2$, (5) has a solution for any $t \geqslant 0$ consisting of a single Dirac measure of suitable size and position (both depending on $t$ ). We pose

Problem I. For given $v\left(x_{i}\right)=a_{i}>0(1 \leqslant i \leqslant n ; n \geqslant 3)$, find the supremum $T$ of all $t$ for which there exists a positive bounded $u_{0}$ satisfying (5).

For $t$ fixed, our problem is a finite moment problem. We take the following condition for existence of a solution to this problem from Rogosinski (1958), Theorem 1 and Corollary 1 :

Theorem R. There exists a positive $u_{0}$ satisfying (5) if and only if the point $a=$ $\left(a_{1}, a_{2}, \ldots, a_{n}\right) \in R^{n}$ is in the hull cone ${ }^{1} P_{t}$ of the curve

$$
\begin{equation*}
p_{t}(x)=\left(\psi_{t}\left(x_{1}-x\right), \psi_{t}\left(x_{2}-x\right), \ldots, \psi_{t}\left(x_{n}-x\right)\right), \quad-\infty<x<+\infty . \tag{6}
\end{equation*}
$$

By this theorem we have

$$
\begin{equation*}
T=\sup \left\{t: a \in P_{t}\right\} . \tag{7}
\end{equation*}
$$

We shall investigate the properties of $P_{t}$.
Lemma 1. $P_{t_{2}} \subset P_{t_{1}}$ if $t_{1} \leqslant t_{2}$.
Proof. If $y \in P_{t_{2}}$ there exists a positive measure $\mu$ such that $y=\mu * p_{t_{2}}$. Then we have by (4) $y=\mu * p_{t_{2}}=\mu * \psi_{t_{2}-t_{1}} * p_{t_{1}}$. Since $\mu * \psi_{t_{2}-t_{1}}$ is a positive measure the lemma is proved.

Since $\psi_{t}(x) \geqslant 0, P_{t}$ is a subset of the positive orthant in $R^{n}$. If $\boldsymbol{x} \neq 0, \psi_{t}(x) \rightarrow 0$ when $t \rightarrow 0$, implying that the ray from the origin through the point $p_{t}\left(x_{i}\right)$ approaches the $y_{i}$-axis of $R^{n}$ as $t \rightarrow 0$. Thus, $P_{t}$ is monotonically increasing to the whole positive orthant when $t \rightarrow 0$. Since we have assumed $a_{i}>0(\mathrm{l} \leqslant i \leqslant n)$, there exists an $\varepsilon>0$ such that $a \in P_{\varepsilon}$.

When $t \rightarrow \infty, P_{t}$ decreases to a subcone, say $P_{\infty}$ of the positive orthant. The cone $P_{\infty}$ is the set of points $a$ for which (5) has a solution for all $t . P_{\infty}$ is described by theorem 1 in the case of equidistant $x_{i}$.
By Lemma 1 we have

[^0]\[

$$
\begin{equation*}
\sup \left\{t: a \in P_{t}\right\}=\inf \left\{t: a \notin P_{t}\right\} \tag{9}
\end{equation*}
$$

\]

if we interpret the right-hand side of (9) as $\infty$ when $a \in P_{\infty}$.

## III. Equidistant data

Assume that the values $a_{i}=v\left(x_{i}\right)$ are obtained at equidistant points, i.e. assume $x_{i}=b+i h(\mathbf{l} \leqslant i \leqslant n)$, where $b$ and $h$ are constants. Since the position of the origin on the $x$-axis is immaterial, we put $b=0$.

Theorem 1. Assume $a_{i}>0(1 \leqslant i \leqslant n)$ and define $\tau=\exp \left(h^{2} / 4 t\right)$ for $t>0$ so that $t=h^{2} / 4 \log \tau$. Consider the quadratic forms

$$
\begin{aligned}
Q_{1}(\tau)=\Sigma_{i} \Sigma_{m} a_{i+m} \tau^{(i+m)^{3}} \xi_{i} \xi_{m}, & 1 \leqslant i, m \leqslant[n / 2], \\
Q_{2}(\tau)=\Sigma_{i} \Sigma_{m} a_{i+m-1} \tau^{(i+m-1)^{2}} \xi_{i} \xi_{m}, & 1 \leqslant i, m \leqslant[(n+1) / 2] .
\end{aligned}
$$

Let $\tau_{0}$ be the smallest $\tau \geqslant 1$ such that the forms $Q_{1}$ and $Q_{2}$ are both positive semidefinite for $\tau_{0} \leqslant \tau<\infty$. If $\tau_{0}>1$, we have $T=h^{2} / 4 \log \tau_{0}$. If $\tau_{0}=1$, then $T=\infty$, that is $a \in P_{\infty}$ and the equations (5) are solvable for all $t \geqslant 0$.

Proof. A symmetric matrix is positive definite if all its diagonal subdeterminants are positive. For the matrices of the forms $Q_{1}$ and $Q_{2}$ these determinants are polynomials in $\tau$ and it is easily shown that their leading coefficients are positive. By the definition of $\tau_{0}$ we then know that there exists a $\tau_{1}$ such that $Q_{1}(\tau)$ and $Q_{2}(\tau)$ both are strictly positive for $\tau_{0}<\tau<\tau_{1}$.

Now, we write out the equations (5) with $x_{i}=i h$

$$
a_{i}=(1 / 2 V / \pi t) \int_{-\infty}^{+\infty} \exp \left(-(i h-x)^{2} / 4 t\right) u_{0}(d x)
$$

Rearranging (5'), we get

$$
a_{i}=\exp \left(-i^{2} h^{2} / 4 t\right) \int_{-\infty}^{+\infty} \exp (i h x / 2 t)(1 / 2 \sqrt{\pi t}) \exp \left(-x^{2} / 4 t\right) u_{0}(d x)
$$

Now, the measure $w=(1 / 2 \sqrt{\pi t}) \exp \left(-x^{2} / 4 t\right) u_{0}$ is positive if and only if $u_{0}$ is positive, so we have the question: for which $t$ does there exist a positive $w$ satisfying

$$
a_{i}=\exp \left(-i^{2} h^{2} / 4 t\right) \int_{-\infty}^{+\infty} \exp (i h x / 2 t) w(d x)
$$

We make a change of variable in the integral by putting $\exp (h x / 2 t)=\eta$. Since $\eta$ is a monotonic function of $x$, the positive measure $w(d x)$ changes to a positive measure, say $\mu(d \eta)$, and we get
or

$$
\begin{align*}
& a_{i}=\exp \left(-i^{2} h^{2} / 4 t\right) \int_{0}^{+\infty} \eta^{i} \mu(d \eta), \\
& a_{i} \tau^{i}=\int_{0}^{+\infty} \eta^{i} \mu(d \eta) \quad(1 \leqslant i \leqslant n) \tag{10}
\end{align*}
$$

## J. PHILIP, Age of a heat distribution

This setting of the problem is known as Stieltjes' moment problem. A sufficient condition for the possibility of representing the quantities $a_{i} \tau^{i^{2}}$ by a positive measure as in (10) is the strict positivity of the forms $Q_{1}$ and $Q_{2}$. (See e.g. Krein, 1951.) Thus, if $\tau_{0}<\tau<\tau_{1}$ we have a representation (10), proving the solvability of (5) for $\tau$ in the open interval ( $\tau_{0}, \tau_{1}$ ). By Lemma 1, however, (5) is then solvable for every $\tau>\tau_{0}$. Theorem R shows that the solution $u_{0}$ can be taken as a finite sum of Dirac measures. Conversely, if (5) has a solution $\bar{\mu}$ for $\bar{\tau}>\mathbf{l}$ we have
and

$$
\begin{aligned}
& Q_{1}(\bar{\tau})=\int_{0}^{\infty}\left(\sum_{m} \xi_{m} \eta^{m}\right)^{2} \bar{\mu}(d \eta) \geqslant 0, \\
& Q_{2}(\bar{\tau})=\int_{0}^{\infty}\left(\sum_{m} \xi_{m} \eta^{m}\right)^{2}(1 / \eta) \bar{\mu}(d \eta) \geqslant 0 .
\end{aligned}
$$

implying $\bar{\tau} \geqslant \tau_{0}$. In the case $\tau_{0}>\mathbf{l}$ there is consequently no solution for $\tau$ in the open interval ( $1, \tau_{0}$ ), proving that $T=h^{2} / 4 \log \tau_{0}$. In the case $\tau_{0}=1$ the eqs. (5) are solvable for every $t \geqslant 0$ implying $T=\infty$.

Remark. The positivity of $Q_{1}$ or $Q_{2}$ is a condition on an odd number of consecutive $a_{i}$. For $n$ odd, $Q_{1} \geqslant 0$ is a condition on $a_{2}, a_{3}, \ldots, a_{n-1}$ and $Q_{2} \geqslant 0$ a condition on $a_{1}, a_{2}, \ldots, a_{n}$. For $n$ even, $Q_{1} \geqslant 0$ is a condition on $a_{2}, a_{3}, \ldots, a_{n}$ and $Q_{2} \geqslant 0$ a condition on $a_{1}, a_{2}, \ldots, a_{n-1}$.

Example 1. Let $n=3$ so that $a_{1}, a_{2}$ and $a_{3}$ are given. The positivity of $Q_{1}$ and $Q_{2}$ then gives the inequalities $a_{2} \tau^{4} \geqslant 0$ and $a_{1} a_{3} \tau^{10}-a_{2}^{2} \tau^{8} \geqslant 0$. If $a_{2}^{2}>a_{1} a_{3}$, we get the estimate $\tau_{0}=a_{2} / \sqrt{a_{1} a_{3}}$ or $T=h^{2} / 4 \log \left(a_{2} / \sqrt{a_{1}} a_{3}\right)$. If $a_{2}^{2} \leqslant a_{1} a_{3}, T=\infty$.

Example 2. This example shows that there does not necessarily exist a solution of (5) for $t=T$. Let $n=4$ so that $a_{1}, a_{2}, a_{3}$ and $a_{4}$ are given. The positivity of $Q_{1}$ and $Q_{2}$ gives

$$
a_{1} a_{3} \tau^{10}-a_{2}^{2} \tau^{8} \geqslant 0 \quad \text { and } \quad a_{2} a_{4} \tau^{20}-a_{3}^{2} \tau^{18} \geqslant 0 .
$$

Suppose $a_{2}^{2} / a_{1} a_{3}>\max \left(1, a_{3}^{2} / a_{2} a_{4}\right)$ so that $\tau_{0}^{2}=a_{2}^{2} / a_{1} a_{3}$. Krein (1951) shows that if there exists a representation (10), there exists one of the form

$$
\begin{equation*}
a_{i} \tau^{i^{2}}=\varrho_{1} \eta_{1}^{i}+\varrho_{2} \eta_{2}^{i}(1 \leqslant i \leqslant 4), \varrho_{1}>0, \varrho_{2}>0, \eta_{1} \leqslant \eta_{2} . \tag{11}
\end{equation*}
$$

The four eqs. (11) are actually just sufficient to determine the four quantities $\varrho_{1}, \varrho_{2}$, $\eta_{1}$ and $\eta_{2}$. Combining the eqs. ( 11 ) for ( $1 \leqslant i \leqslant 3$ ) we get

$$
\varrho_{1} \varrho_{2} \eta_{1} \eta_{2}\left(\eta_{1}-\eta_{1}\right)^{2}=\tau^{8}\left(a_{1} a_{3} \tau^{2}-a_{2}^{2}\right) .
$$

For $\tau=\tau_{0}$ these expressions equal 0 . Thus, either $\eta_{1}=0$ or $\eta_{1}=\eta_{2}$, so the representation has only one pointmass and has the form $\varrho_{2} \eta_{2}^{i}$. This representation does not satisfy the equation for $i=4$.

## IV. Nonequidistant data

When the points $x_{i}$ are not equidistant, we only give an estimate that takes into account three points $a_{i}=v\left(x_{i}\right)$. Of course, if more than three values are known, those three that give the smallest estimate $T$ should be used.

Theorem 2. Let $a_{i}=v\left(x_{i}\right)>0(1 \leqslant i \leqslant 3)$, be the values of $v$ at three points $x_{1}<x_{2}<x_{3}$. Then,

$$
T=\frac{(1 / 4)\left(x_{3}-x_{1}\right)\left(x_{3}-x_{2}\right)\left(x_{2}-x_{1}\right)}{\left(x_{3}-x_{1}\right) \log a_{2}-\left(x_{3}-x_{2}\right) \log a_{1}-\left(x_{2}-x_{1}\right) \log a_{3}}
$$

if this expression is positive. In other cases $T=\infty$.
Proof. By (7) and (9) $T$ can be characterized by the fact that if $t>T$, then there exists a plane in $R^{3}$ separating $a$ from $P_{t}$. Let

$$
\alpha^{T} y=\alpha_{1} y_{1}+\alpha_{2} y_{2}+\alpha_{3} y_{3}=0
$$

be a plane in $R^{3}$. We normalize it by putting $\alpha_{2}=-1$. We shall try to find $\alpha_{1}$ and $\alpha_{3}$ so that

$$
f(x)=\alpha^{T} p_{t}(x) \geqslant 0 \quad \text { for all } x
$$

and

$$
\alpha^{T} a \leqslant 0 .
$$

If this is possible, we obtain a plane that separates (not strictly) $a$ and $P_{t}$. Further, we shall choose $\alpha_{1}$ and $\alpha_{3}$ so that the plane supports the cone $P_{t}$ along the ray through $p_{t}(\Theta)$ for some $\Theta$, giving the equation

$$
\begin{equation*}
f(\Theta)=\alpha^{T} p_{t}(\Theta)=0 \tag{12}
\end{equation*}
$$

Since we require $f(x) \geqslant 0$, we must also have

$$
\begin{equation*}
f^{\prime}(\Theta)=\alpha^{T} p_{t}^{\prime}(\Theta)=0 \tag{13}
\end{equation*}
$$

The equations (12) and (13) determine $\alpha_{1}$ and $\alpha_{3}$ as functions of $t$ and $\Theta$ :

$$
\begin{aligned}
& \alpha_{1}(t, \Theta)=\left[\left(x_{3}-x_{2}\right) /\left(x_{3}-x_{1}\right)\right] \exp \left(\left(x_{1}-x_{2}\right)\left(x_{1}+x_{2}-2 \Theta\right) / 4 t\right), \\
& \alpha_{3}(t, \Theta)=\left[\left(x_{2}-x_{1}\right) /\left(x_{3}-x_{1}\right)\right] \exp \left(\left(x_{3}-x_{2}\right)\left(x_{3}+x_{2}-2 \Theta\right) / 4 t\right) .
\end{aligned}
$$

To see if there exist $t$ and $\Theta$ so that

$$
\alpha^{T} a=\alpha_{1}(t, \Theta) a_{1}-a_{2}+a_{3}(t, \Theta) a_{3} \leqslant 0
$$

we seek $\min _{\Theta} \alpha^{T} a$ by solving $(\partial / \partial \Theta) \alpha^{T} a=0$. We find a unique solution

$$
\begin{equation*}
\Theta_{0}=\frac{1}{2} x_{1}+\frac{1}{2} x_{3}+\left(2 t \log a_{3} / a_{1}\right) /\left(x_{3}-x_{1}\right) \tag{14}
\end{equation*}
$$

and get

$$
\min _{\Theta} \alpha^{T} a=\left[a_{1}^{\left(x_{3}-x_{2}\right)} a_{3}^{\left(x_{2}-x_{2}\right)}\right]^{1 /\left(x_{3}-x_{1}\right)} \exp \left(\frac{\left(x_{2}-x_{1}\right)\left(x_{3}-x_{2}\right)}{4 t}\right)-a_{2}
$$

For some $a_{1}, a_{2}$ and $a_{3}$, this expression is not negative even when $t \rightarrow \infty$, meaning that there does not exist any plane separating $a$ and $P_{t}$ even for large $t$. From such $a_{i}$, our theorem gives no estimate. For other $a_{i}$, the expression is negative for large $t$. Then $T$ is the value of $t$ for which $\min _{\Theta} \alpha^{T} a=0$.

## J. PHILIP, Age of a heat distribution

Remark. For three equidistant points, Theorem 2 gives the same estimate as Theorem 1. Cf. example 1.

Remark. For equidistant $x_{i}$, the method of proof for theorem 2 can be generalized to handle 5,7 , etc. points by introducing 2,3 , etc. points of contact like $\Theta$. Such a generalization gives exactly the estimate of theorem 1. Also for nonequidistant $x_{i}$, the generalization of the above method is usable, but explicit expressions corresponding to (14) are not obtained.

## V. Practical measurements

Since it is hard to construct a thermometer that gives the temperature $v\left(x_{i}\right)$ at a single point, we shall consider the situation that the information about $v$ is given by the values of $n(\geqslant 3)$ linear functionals

$$
\begin{equation*}
d_{i}=\int v\left(x_{i}-x\right) v(d x) \quad(1 \leqslant i \leqslant n), \tag{15}
\end{equation*}
$$

where $\boldsymbol{v}$ is a positive measure with total mass equal to 1 . The measure $\boldsymbol{v}$ is thought to describe the measuring instrument and the values $d_{i}$ are obtained when the instrument is "centered" at the points $x_{i}$. Physicist often assume

$$
\begin{equation*}
\nu(d x)=(1 / \sqrt{2 \pi} \sigma) \exp \left(-x^{2} / 2 \sigma^{2}\right) d x=\psi_{\sigma^{2} / 2}(x) d x \tag{16}
\end{equation*}
$$

for some $\sigma>0$. The limit case when $\sigma \rightarrow 0$ corresponds to $\nu \rightarrow \delta$ and $d_{i} \rightarrow v\left(x_{i}\right)=a_{i}$, that is the case in secs II-IV. We pose

Problem II. For given $d_{i}>0(1 \leqslant i \leqslant n, n \geqslant 3)$, find an upper bound for the supremum $\tilde{T}$ of all $t$ for which there exists a positive bounded $u_{0}$ satisfying (15).

Remark. This problem has not always a solution since there may not exist any $t$ for which the $d_{i}(1 \leqslant i \leqslant n)$ is a conceivable set of data. Cf. the corollary of Lemma 3.

The equations (15) can also be written

$$
\begin{equation*}
d_{i}=v * v\left(x_{i}\right) \quad(1 \leqslant i \leqslant n) \tag{15'}
\end{equation*}
$$

or by (3)

$$
d_{i}=u_{0} * \psi_{t} * v\left(x_{i}\right) \quad(1 \leqslant i \leqslant n) .
$$

Let

$$
\begin{equation*}
x_{t}=\psi_{t} * v \tag{17}
\end{equation*}
$$

so that

$$
\begin{equation*}
d_{i}=x_{i} * u_{0}\left(x_{i}\right) \quad(1 \leqslant i \leqslant n) . \tag{18}
\end{equation*}
$$

By Rogosinski (1958), theorem 1 and corollary 1 we get
Theorem $\mathbf{R}^{\prime}$. There exists a positive bounded $u_{0}$ satisfying (18) if and only if the point $d=\left(d_{1}, d_{2}, \ldots, d_{n}\right) \in R^{n}$ is in the hull cone $K_{t}$ of the curve

$$
k_{t}(x)=\left(\varkappa_{t}\left(x_{1}-x\right), \varkappa_{t}\left(x_{2}--x\right), \ldots, \varkappa_{t}\left(x_{n}-x\right)\right), \quad-\infty<x<+\infty .
$$

By this theorem we have

$$
\begin{equation*}
\widetilde{T}=\sup \left\{t: d \in K_{t}\right\} . \tag{19}
\end{equation*}
$$

Most of the results of section II have their counterparts here.
Lemma 2.

$$
\begin{equation*}
K_{t_{2}} \subset K_{t_{1}} \quad \text { if } \quad t_{\mathbf{1}} \leqslant t_{2} . \tag{20}
\end{equation*}
$$

Proof. If $y \in K_{t_{2}}$ there exists a positive measure $\mu$ such that $y=\mu * k_{t_{2}}=\mu * p_{t_{2}} * \nu=$ $\mu * \psi_{t_{2}-t_{1}} * p_{t_{1}} * \nu=\mu * \psi_{t_{2}-t_{1}} * k_{t_{1}}$. Since $\mu * \psi_{t_{2}-t_{1}}$ is a positive measure the lemma is proved.

By this lemma we have, as in section II, that

$$
\begin{equation*}
\widetilde{T}=\sup \left\{t: d \in K_{t}\right\}=\inf \left\{t: d \notin K_{t}\right\} \tag{21}
\end{equation*}
$$

if we interpret the infimum as $\infty$ when $d \in K_{t}$ for all $t \geqslant 0$.
Define

$$
\vartheta=\sup \left\{t: \nu=\psi_{t} * \lambda, \lambda \text { positive measure }\right\} .
$$

The total mass of $\lambda$ equals that of $\nu$ so it is equal to 1 . We can think of $\psi_{t}, \nu$ and $\lambda$ as probability distributions corresponding to random variables $X_{\psi}, X_{\nu}$ and $X_{\lambda}$. Taking variances, we get

$$
\operatorname{Var}\left(X_{\nu}\right)=2 \vartheta+\operatorname{Var}\left(X_{\lambda}\right) .
$$

Since $\psi_{t} \rightarrow \delta$ when $t \rightarrow 0$ and $\operatorname{Var}\left(X_{\lambda}\right) \geqslant 0$ we have

$$
0 \leqslant \vartheta \leqslant \frac{1}{2} \operatorname{Var}\left(X_{\nu}\right)
$$

Of course, $\vartheta>0$ only when $v$ is infinitely differentiable. For the particular $v$ in (16) we get $\vartheta=\frac{1}{2} \operatorname{Var}\left(X_{\nu}\right)=\frac{1}{2} \sigma^{2}$.

## Lemma 3.

$$
K_{t} \subset P_{t+\vartheta}
$$

Proof. If $y \in K_{t}$ there exists a positive measure $\mu$ such that $y=k_{t} * \mu=p_{t} * \nu * \mu=$ $p_{t} * \psi_{\vartheta} * \lambda * \mu=p_{t+\vartheta} * \lambda * \mu$. Since $\lambda * \mu$ is a positive measure, the lemma is proved.

Corollary. In general, $K_{0}=\lim _{t \rightarrow 0} K_{t}$ is not the whole positive orthant. Since Problem II has a solution if and only if $d \in K_{0}$, the condition $d_{i}>0(1 \leqslant i \leqslant n)$ is not in general sufficient for Problem II to have a solution.

Theorem 3. Assume $d \in K_{0}$ and apply Theorem 1 or 2 with $a=d$ so that an estimate $T$ is obtained. Then,

$$
\tilde{T} \leqslant T-\vartheta
$$

Proof. For every $t<\tilde{T}$ we have $d \in K_{t} \subset P_{t+\vartheta}$ by Lemma 3. Then, $t+\vartheta \leqslant T=$ $\sup \left\{t: d \in P_{t}\right\}$ proving the theorem.

The physical interpretation of the last formula is that the dispersion of $u_{0}$ caused by heat conduction during the time $\widetilde{T}$ and the dispersion caused by the measuring instrument, which is larger than if it had been aged $\vartheta$, together form a dispersion corresponding to heat conduction during a time $T$.

## J. PHILIP, Age of a heat distribution

## ACKNOWLEDGEMENT

The problem here treated was suggested by Hans Rådström. He has also made many valuable remarks during the preparation of this paper.

Department of Mathematical Statistics, Royal Institute of Technology, Stockholm, Sweden

## REFERENCES

Krein, M. G., The ideas of P. L. Chebyshev and A. A. Markov in the theory of limiting values of integrals and their further development, Amer. Math. Soc. Translations, ser. 2, no 12, 3-120 (1951).
Rogosinski, W. W., Moments of non-negative mass, Proc. Roy. Soc. (London), A245, 1-27 (1958)


[^0]:    ${ }^{1}$ The hull cone of a set $A$ is defined as the smallest convex cone with vertex 0 that contains $A$.

