# Spaces of Carleson measures: duality and interpolation 

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## 1. Introduction

In [12], [13] Coifman, Meyer and Stein have developed a theory of "tent spaces" with interesting applications. Their theory has led to a unification and simplification of some basic techniques in harmonic analysis.

The theory of "tent spaces" is closely related to the one of Hardy spaces. In this paper we consider the relationship of these tent spaces with spaces of Carleson measures. In particular we identify the spaces of Carleson measures as the duals of certain tent spaces.

We also compute the real interpolation spaces of some extreme tent spaces by means of computing the corresponding $K$ functionals of Peetre. As pointed out in [13] the interpolation theory of tent spaces can be used to derive the corresponing theory for $H^{p}$ spaces and Lipschitz spaces.

Let us briefly explain the motivation of this paper (we refer to $\S 2$ for detailed definitions). A basic inequality valid for Carleson measures on $\mathbf{R}_{+}^{n+1}$ is

$$
\begin{equation*}
|\mu|\left\{(x, t)||f(x, t)|>\lambda\} \leqq c\left|\left\{x / A_{\infty}(f)(x)>\lambda\right\}\right|\right. \tag{1.1}
\end{equation*}
$$

where $A_{\infty}(f)(x)=\sup _{\Gamma(x)}|f(y, t)|, \Gamma(x)=$ cone with vertex $x$. Then

$$
\begin{equation*}
\left|\int f(x, t) d \mu(x, t)\right| \leqq c\left(\sup _{B} \frac{1}{|B|} \int_{T(B)} d|\mu|\right)\left\|A_{\infty}(f)\right\|_{1} \tag{1.2}
\end{equation*}
$$

where $T(B)=$ "tent with base $B$ ".
Let $\|\mu\|_{V^{1}}=\left\|C_{1}(\mu)\right\|_{\infty}$, where $C_{1}(\mu)(x)=\sup _{B \ni x} \frac{1}{|B|} \int_{T(B)} d|\mu|$. Then (1.2) can

[^0]be interpreted as a duality result between $T_{\infty}^{1}=\left\{f \mid A_{\infty}(f) \in L^{1}\right\}$ and $V^{1}=$ space of Carleson measures on $\overline{\mathbf{R}}_{+}^{n+1}$, i.e., $\left(T_{\infty}^{1}\right)^{*}=V^{1}$ (cf. [2]).

It is of interest to study inequalities of the form (1.2) by means of considering spaces of Carleson measures defined by the boundedness of the functionals $\sup _{B \ni x} \frac{1}{|B|^{\alpha}} \int_{T(B)} d|\mu|, \alpha>0$. These spaces, $V^{\alpha}$, have been considered by several authors: Duren [15], Barker [3], Amar and Bonami [2], Johnson [18].

In this paper we develop this remark systematically and use it to compute the duals of certain "tent spaces" as spaces of Carleson measures. These extreme "tent spaces" therefore play an important role in the study of maximal operators. We also develop the interpolation theory of these extreme spaces completing the results of Coifman, Meyer and Stein [13] as well as those of Amar and Bonami [2].

We should mention that while this paper was being prepared for publication we received a preprint by Bonami and Johnson [5] that contains some of our results. Our approach is, however, different.

We wish to thank Professor E. Stein for several useful conversations on the subject of our work. We are also grateful to the referee for many valuable suggestions including formula (4.5) and a correction to our proof of (7.3) below.

## 2. Preliminaries

We shall work on $\mathbf{R}_{+}^{n+1}$ but most of our results are valid, and useful, in the more general context of homogeneous spaces.

Let $\Omega$ be an open set in $\mathbf{R}^{n}$, we let $T(\Omega)$ be the subset of $\mathbf{R}_{+}^{n+1}$ defined by $T(\Omega)=$ $\{(x, t) \mid B(x, t) \subset \Omega\}$, where $B(x, t)$ denotes the ball with centre $x$ and radius $t$. As usual, $|\Omega|$ will stand for the Lebesgue measure of the set $\Omega$.

Given $\alpha \in \mathbf{R}$, we say that a measure $w$ in $\overline{\mathbf{R}}_{+}^{n+1}$ is a Carleson measure of order $\alpha$, if $\forall \Omega \subset X, \Omega$ open bounded,

$$
\begin{equation*}
|w|(T(\Omega)) \leqq c|\Omega|^{\alpha} \tag{2.1}
\end{equation*}
$$

where $c$ is an absolute constant independent of $\Omega$. In fact, we let

$$
\begin{equation*}
\|w\|_{V^{\alpha}}=\inf \{c:(2.1) \text { holds }\} \tag{2.2}
\end{equation*}
$$

and $V^{\alpha}=\left\{w \mid w\right.$ measure on $\left.\overline{\mathbf{R}}_{+}^{n+1},\|w\|_{V^{\alpha}}<\infty\right\}$. Notice, in particular, that $V^{0}=$ \{finite measures on $\overline{\mathbf{R}}_{+}^{n+1}$ \}.

The case $\alpha=1$ corresponds to the measures studied by Carleson [7].
We consider kernels $P_{t}(x, y)$ on $\mathbf{R}^{n} \times \mathbf{R}^{n}$ and the corresponding integral operators associated with them $P_{t} f(x)=\int P_{t}(x, y) f(y) d y$. The maximal operators
associated with these operators are given by

$$
A_{\infty} P(f)(x)=\sup _{(y, t \in \Gamma(x)}\left|P_{t} f(y)\right|
$$

The kernels we consider always satisfy:

$$
\begin{equation*}
\left\|A_{\infty} P(f)\right\|_{p} \leqq c_{p}\|f\|_{p}, \quad 1<p<\infty \tag{2.3}
\end{equation*}
$$

Important examples are provided by the Poisson kernels in different settings and also by $P_{t}^{0}(x, y)=|B(x, t)|^{-1} \chi_{B(x, t)}(y)$. As is well known conditions of the form (2.1) are useful to estimate the distribution function of $P_{t} f$ in terms of $A_{\infty} P(f)$. In fact, given $\lambda>0$, let $\Omega_{\lambda}=\{x \mid M f(x)>\lambda\}$, where $M$ denotes the maximal operator of Hardy-Littlewood. Then

$$
\begin{equation*}
\left\{(x, t) \mid P_{t} f(x)>\lambda\right\} \subset T\left(\Omega_{\lambda}\right) \tag{2.4}
\end{equation*}
$$

Using (2.4) and the usual notation of Lorentz spaces (cf. [17]) we get, for $\alpha>0$,

$$
\begin{equation*}
\left\|P_{t} f\right\|_{L^{p \cdot \alpha}(d w)} \leqq c\|w\|_{V^{\alpha}}\|f\|_{L(p, p \cdot \alpha)} \tag{2.5}
\end{equation*}
$$

In particular the choice $\alpha=\frac{1}{p}$, gives (cf. [2])

$$
\begin{equation*}
\int\left|P_{t} f(x)\right| d|w|(x, t) \leqq c\|w\|_{V^{1 / p}}\|f\|_{L(p, 1)} \tag{2.6}
\end{equation*}
$$

which can be interpreted as a duality formula.
The tent spaces $T_{q}^{p}, 0<p \leqq \infty, 0<q \leqq \infty$ are defined using two families of functionals. The $A_{q}$ functionals defined by:

$$
\begin{gather*}
A_{q}(f)(x)=\left\{\int_{\Gamma(x)}|f(y, t)|^{q} \frac{d y d t}{t^{n+1}}\right\}^{1 / q},  \tag{2.7}\\
A_{\infty}(f)(x)=\sup _{\Gamma(x)}|f(y, t)| \tag{2.7}
\end{gather*}
$$

where, as usual, $\Gamma(x)$ is the cone with vertex $x,\{(y, t)| | x-y \mid<t\}$.
The $C_{q}$ functionals are given by

$$
\begin{equation*}
C_{q}(f)(x)=\sup \left\{\frac{1}{|B|} \int_{T(B)}|f(y, t)|^{q} \frac{d y d t}{t}\right\}^{1 / q} \tag{2.8}
\end{equation*}
$$

where the sup is taken over all balls containing $x$.
Let us observe that the $A_{1}$ and $C_{1}$ functionals can be defined for measures. This observation will be particularly useful for us in what follows. (See 2.9).

The tent spaces $T_{q}^{p}$ are defined by the condition $A_{q}(f) \in L^{p}, 0<q \leqq \infty, 0<p<\infty$, (if $q=\infty$ we impose some additional restrictions such as the continuity of $f$ (cf. [13] and (7.1) below). One defines similarly the $T_{q}^{\infty}$ spaces using $C_{q}$ functionals (cf. [13]).

In view of (2.5) it is of interest to define tent spaces based on Lorentz spaces; this can be done simply by replacing the $L^{p}$ norm condition by an $L(p, r)$ norm condition
in the appropriate definitions. The spaces thus obtained will be denoted $T_{q}^{p, r}$ and will be shown to appear naturally when we interpolate $T_{q}^{p}$ spaces.

Let us note in passing that this description of the $T_{q}^{p, r}$ spaces as interpolation spaces, proves their completeness and provides the duality theory for them, at least when $1<p<\infty, 1 \leqq q<\infty$, from the duality theory developed for the $T_{q}^{p}$ spaces in [13]. Alternatively one could deduce this part of the theory using similar arguments to those given in [13] (cf. also [15]).

Our attention in this paper will be devoted to the study of the extremal spaces $T_{1}^{\infty}, T_{\infty}^{p}, T_{\infty}^{p, 1}$ and their duals. Let us start by making explicit the remark that followed (2.8).
(2.9) Definition. The $A_{1}, C_{1}$ functionals for measures are given by

$$
\begin{gather*}
A_{1}(w)(x)=\int_{\Gamma(x)} t^{-n} d|w|(y, t)  \tag{2.10}\\
C_{1}(w)(x)=\sup _{B \ni x} \frac{1}{|B|} \int_{T(B)} d|w|(y, t) . \tag{2.11}
\end{gather*}
$$

For $0<p<\infty, 0<q \leqq \infty$, let $\tau_{1}^{p, q}=\left\{w \mid A_{1}(w) \in L^{p, q}\right\}, \tau_{1}^{\infty}=\left\{w \mid C_{1}(w) \in L^{\infty}\right\}$ where $w$ is a measure on $\overline{\mathbf{R}}_{+}^{n+1}$. One easily cheks that $\tau_{1}^{\infty}=V^{1}$, on the other hand, if $1<p<\infty$, $\tau_{1}^{p, p}=\tau_{1}^{p}$ coincides with the space $W^{1 / p^{\prime}}$ introduced by Amar and Bonami in [2].

The $A_{1}$ and $C_{1}$ functionals are related by

$$
C_{1}(w)(x) \leqq c M\left(A_{1}(w)\right)(x)
$$

In particular,

$$
\begin{equation*}
\left|\left\{x \mid C_{1}(w)(x)>\lambda\right\}\right| \leqq\left|\left\{x \mid c M\left(A_{1}(w)\right)(x)>\lambda\right\}\right| \tag{2.12}
\end{equation*}
$$

where $M$ is the Hardy-Littlewood maximal operator. This follows readily from the definitions (cf. [3]).

Finally, in our estimates of $K$ functionals we need the well-known concept of nonincreasing rearrangment (cf. [17]): for a given function $f$ the nonincreasing rearrangement $f^{*}$ is the generalized inverse of the distribution function $|\{x||f(x)|>t\} \mid$ therefore $\left|\left\{x\left||f(x)|>f^{*}(t)\right\} \mid \leqq t\right.\right.$. The double star function $f^{* *}$ is simply the average of $f^{*}$ at time $t ; f^{* *}(t)=t^{-1} \int_{0}^{t} f^{*}(s) d s, t>0$, and is also decreasing.

## 3. Interpolation of $T_{\infty}^{p}$ spaces

In this section we study the interpolation properties of the $T_{\infty}^{p}$ spaces. Our results complement earlier work by Coifman, Meyer and Stein [13] and Amar and Bonami [2].

We shall consider the real method of interpolation using Peetre's $K$ functionals.

We review briefly some pertinent facts about real interpolation. For more information on interpolation theory we refer the reader to [4].

Let $B_{0}, B_{1}$ be quasi normed spaces embedded in a suitable topological Hausdorff vector space $V$, and define for $f \in B_{0}+B_{1}, t>0$,

For $0<\theta<1,0<q \leqq \infty$, let

$$
K\left(t, f, B_{0}, B_{1}\right)=\inf _{\substack{f=f_{0}+f_{1} \\ f_{i} \in B_{i}}}\left\{\left\|f_{0}\right\|_{B_{0}}+t\left\|f_{1}\right\|_{B_{1}}\right\}
$$

$$
\left(B_{0}, B_{1}\right)_{\theta, q}=\left\{f \in B_{0}+B_{1} /\|f\|_{\left(B_{0}, B_{1}\right)_{\theta, q}}=\left\{\int_{0}^{\infty}\left[t^{-\theta} K\left(t, f, B_{0}, B_{1}\right)\right]^{q} \frac{d t}{t}\right\}^{1 / q}<\infty\right\}
$$

We shall also use, in §4, an important complement to the reiteration theorem obtained recently by T. Wolff [22]: Let $A_{i}, i=0, \ldots$, be quasi Banach spaces continuously embedded in a common topological Hausdorff vector space; let $0<\theta<\eta<1, \theta=$ $\lambda \eta, \eta=(1-\mu) \theta+\mu, 0<\lambda, \mu<1$, then if $A_{1}=\left(A_{0}, A_{2}\right)_{\lambda, p}, A_{2}=\left(A_{1}, A_{3}\right)_{\mu, q}$ we have $A_{1}=\left(A_{0}, A_{3}\right)_{\theta, p}, A_{2}=\left(A_{0}, A_{3}\right)_{\eta, q}$. In other words Wolff's theorem allows to treat each end point space separately.

When considering the $\left\{T_{\infty}^{p}\right\}_{p>0}$ scale it is natural to let $T_{\infty}^{\infty}=L^{\infty}$ since $A_{\infty}(f) \in$ $L^{\infty} \Leftrightarrow f \in L^{\infty}$. Our main result concerning this scale is
(3.1) Theorem. Let $0<p<\infty, 0<\theta<1,0<r \leqq \infty, \frac{1}{p_{\theta}}=\frac{1-\theta}{p}$, then
(i) $\exists c_{1}, c_{2}>0$ such that

$$
c_{1}\left\{\int_{0}^{t^{p}} A_{\infty}(f)^{* p}(s) d s\right\}^{1 / p} \leqq K\left(t, f, T_{\infty}^{p}, L^{\infty}\right) \leqq c_{2}\left\{\int_{0}^{t^{p}} A_{\infty}(f)^{* p}(s) d s\right\}^{1 / p}
$$

In other words,

$$
K\left(t, f, T_{\infty}^{p}, L^{\infty}\right) \approx K\left(t, A_{\infty}(f), L^{p}, L^{\infty}\right)
$$

(ii)

$$
\left(T_{\infty}^{p}, L^{\infty}\right)_{0, r}=T_{\infty}^{p_{0}, r} .
$$

Proof. The second half of the theorem is an easy consequence of the first half.
Consider the proof of (i). Let $f=f_{1}+f_{2}$ be a decomposition of $f$ with $f_{0} \in T_{\infty}^{p}$, $f_{1} \in L^{\infty}$, then

$$
\begin{aligned}
\left\{\int_{0}^{t p} A_{\infty}(f)^{* p}(s) d s\right\}^{1 / p} \leqq c & {\left[\left\{\int_{0}^{t^{p}} A_{\infty}\left(f_{0}\right)^{* p}(s) d s\right\}^{1 / p}+\left\{\int_{0}^{t p} A_{\infty}\left(f_{1}\right)^{* p}(s) d s\right\}^{1 / p}\right] } \\
& \leqq c\left[\left\|f_{0}\right\|_{r_{\infty}^{p}}+t\left\|f_{1}\right\|_{\infty}\right]
\end{aligned}
$$

Thus, taking infimum over all decompositions gives the first inequality of (i).
Let us construct a nearly optimal decomposition to prove the main estimate of (i). Let $t>0, f \in T_{\infty}^{p}+L^{\infty}$, define $\Omega_{t}=\left\{x \mid A_{\infty}(f)(x)>A_{\infty}(f)^{*}\left(t^{p}\right)\right\}$, and consider the
decomposition of $f$ given by $f=f_{0}+f_{1}$, where $f_{0}=f \chi_{T\left(\Omega_{t}\right)}$. Then,

$$
\begin{equation*}
\left\|f_{0}\right\|_{T_{\infty}^{p}}=\left\|A_{\infty}\left(f_{0}\right)\right\|_{p} \leqq c\left\|\chi_{\Omega_{t}} A_{\infty}(f)\right\|_{p} \leqq c\left\{\int_{0}^{t^{p}} A_{\infty}(f)^{* p}(s) d s\right\}^{1 / p} \tag{3.2}
\end{equation*}
$$

Moreover,

$$
t\left\|f_{1}\right\|_{\infty} \leqq\left\{\int_{0}^{t p} A_{\infty}(f)^{* p}(s) d s\right\}^{1 / p}
$$

Combining these estimates we conclude that

$$
\begin{aligned}
K\left(t, f, T_{\infty}^{p}, L^{\infty}\right) & \leqq\left\|f_{0}\right\|_{T_{\infty}}+t\left\|f_{1}\right\|_{\infty} \\
& \leqq c\left\{\int_{0}^{t p} A_{\infty}(f)^{* p}(s) d s\right\}^{1 / p}
\end{aligned}
$$

## 4. Interpolation of $\tau_{1}^{p}$ spaces

In this section we extend the results of [13] (cf. also [2]) concerning the interpolation of the $\left\{\tau_{1}^{p}\right\}$ scale.
(4.1) Theorem. Let $0<p<\infty, 0<\theta<1,0<r \leqq \infty, \frac{1}{p_{\theta}}=\frac{1-\theta}{p}$, then

$$
\begin{equation*}
\left(\tau_{1}^{p}, \tau_{1}^{\infty}\right)_{\theta, r}=\tau_{1}^{p_{0}, r} \tag{4.2}
\end{equation*}
$$

The proof of (4.2) is based on the following estimates for the $K$ functionals for couples of $\tau_{1}^{p}$ spaces
(4.3) Theorem. (i) Let $1 \leqq p<\infty$, then

$$
\begin{equation*}
K\left(t, w, \tau_{1}^{p}, \tau_{1}^{\infty}\right) \leqq c\left\{\int_{0}^{t p} M\left(A_{1}(w)\right)^{* p}(s) d s\right\}^{1 / p} \tag{4.4}
\end{equation*}
$$

where $c$ is a positive constant independent of $w$.
(ii) Let $0<p_{0}<p_{1}<\infty$, then

$$
\begin{equation*}
K\left(t, w, \tau_{1}^{p_{0}}, \tau_{1}^{p_{1}}\right) \approx K\left(t, A_{1}(w), L^{p_{0}}, L^{p_{1}}\right) \tag{4.5}
\end{equation*}
$$

Using (i) and (ii) we can now provide the
Proof of (4.1). Observe that for $1<p<\infty$, we have (cf. [13])

$$
\begin{equation*}
\left\|A_{1}(w)\right\|_{p} \sim\left\|C_{1}(w)\right\|_{p} \tag{4.6}
\end{equation*}
$$

Consider the operator $C_{1}$, then by (4.6) we have

$$
\begin{equation*}
C_{1}: \tau_{1}^{p} \rightarrow L^{p}, \quad C_{1}: \tau_{1}^{\infty} \rightarrow L^{\infty}, \quad p>1 \tag{4.7}
\end{equation*}
$$

Interpolating the estimates in (4.7) and taking into account (4.6) once again gives

$$
\begin{equation*}
\left(\tau_{1}^{p}, \tau_{1}^{\infty}\right)_{\theta, r} \subseteq \tau_{1}^{p_{\theta}, r}, \frac{1}{p_{\theta}}=\frac{1-\theta}{p}, \quad \text { if } p>1 \tag{4.8}
\end{equation*}
$$

The reverse inclusion can now be obtained from (4.4) and Hardy's inequality. Therefore we have obtained

$$
\begin{equation*}
\left(\tau_{1}^{p}, \tau_{1}^{\infty}\right)_{\theta, r}=\tau_{1}^{p_{0}, r}, \frac{1}{p_{\theta}}=\frac{1-\theta}{p}, \quad 1<p<\infty . \tag{4.9}
\end{equation*}
$$

Using (4.5) we readily derive

$$
\begin{equation*}
\left(\tau_{1}^{p_{0}}, \tau_{1}^{p_{1}}\right)_{\theta, r}=\tau_{1}^{p_{\theta}, r}, \quad \frac{1}{p_{\theta}}=\frac{1-\theta}{p_{0}}+\frac{\theta}{p_{1}}, \quad 0<p_{0}<p_{1}<\infty . \tag{4.10}
\end{equation*}
$$

We can now invoke Wolff's theorem to combine (4.9) and (4.10) to obtain (4.2).
We are now ready to provide the
Proof of (4.3). We prove only (4.4) (the assertion (4.5) being similar). Moreover, we only need to consider the case $p=1$.

The proof of (4.4) consists in exhibiting a nearly optimal decomposition. Let $t>0, w \in \tau_{1}^{1}+\tau_{1}^{\infty}$, define $\Omega_{t}=\left\{x \mid C_{1}(w)(x)>c M\left(A_{1}(w)\right)^{*}(t)\right\}$, and consider the decomposition $w=\mu_{0}+\mu_{1}$, where $\mu_{0}=w \chi_{T\left(\Omega_{t}\right)}$. We have,

$$
\left\|\mu_{0}\right\|_{\tau_{1}^{1}}=\left\|A_{1}\left(\mu_{0}\right)\right\|_{1} \leqq\left\|A_{1}(w) \chi_{\Omega_{t}}\right\|_{1} \leqq \int_{0}^{\left|\Omega_{t}\right|} M\left(A_{1}(w)\right)^{*}(s) d s
$$

Now, recall that by (2.12)

$$
\left|\Omega_{t}\right| \leqq c\left|\left\{x \mid M\left(A_{1}(w)\right)(x)>M\left(A_{1}(w)\right)^{*}(t)\right\}\right| \leqq c t, c>1
$$

Therefore, using the fact that $f^{* *}(t) \downarrow$, we obtain

$$
\begin{equation*}
\left\|\mu_{0}\right\|_{\tau_{1}^{1}} \leqq c \int_{0}^{t} M\left(A_{1}(w)\right)^{*}(s) d s \tag{4.11}
\end{equation*}
$$

Consider now $C_{1}\left(\mu_{1}\right)(x)$; clearly if $x \in \Omega_{t}^{c}$ we get

$$
\begin{aligned}
t C_{1}\left(\mu_{1}\right)(x) & \leqq t C_{1}(w)(x) \leqq t M\left(A_{1}(w)\right)^{*}(t) \\
& \leqq \int_{0}^{t} M\left(A_{1}(w)\right)^{*}(s) d s
\end{aligned}
$$

Suppose now that $x \in \Omega_{t}$. In order to estimate $C_{1}\left(\mu_{1}\right)(x)$ we use a Whitney decomposition of the open set $\Omega_{t}$. (See [21], p. 167.) Therefore, $\Omega_{t}=\bigcup_{k=1}^{\infty} Q_{k}$, where the $Q_{k}$ 's are disjoint cubes such that dist $\left(Q_{k}, \Omega_{t}^{c}\right) \approx \operatorname{diam}\left(Q_{k}\right)$. More precisely, for each cube $Q_{k}$ let $x_{k}^{0}$ be its center and $2 r_{k}$ its diameter, we also let $\bar{Q}_{k}$ be the cube concentric with $Q_{k}$ and such that $\operatorname{diam} \bar{Q}_{k}=10 r_{k}$; then, $\frac{3}{4}\left(2 r_{k}\right)<\operatorname{dist}\left(\bar{Q}_{j}, \Omega_{t}^{c}\right)$, and dist $\left(Q_{k}, \Omega_{t}^{c}\right) \leqq$ $\leqq 4\left(2 r_{k}\right)$.

Let $B=B(x, r)$ be a ball centered at $x$ with radius $r$. If $T(B) \subset T\left(\Omega_{t}\right)$, then we clearly have $1 /|B| \int_{T(B)} d\left|\mu_{1}\right|=0$. Thus, from now on we suppose that $T(B) \cap$ $\left(T\left(\Omega_{t}\right)\right)^{c} \neq \emptyset$. Therefore, there exists $(y, s) \in T(B)$ such that $(y, s) \notin T\left(\Omega_{t}\right)$. If $(y, s) \notin$ $T\left(\Omega_{t}\right)$ then $B(y, s) \nsubseteq T\left(\Omega_{t}\right)$. Two possibilities arise: a) $y \notin \Omega_{t}$ or b) $y \in \Omega_{t}$ but $\operatorname{dist}\left(y, \Omega_{t}^{c}\right)<s$.

We consider the case b) first. By our assumption $y \in Q_{k}$ for some $k \in \mathbf{N}$. Therefore,

$$
\left(\frac{3}{4}\right)\left(2 r_{k}\right)<\operatorname{dist}\left(y, \Omega_{t}^{c}\right)<s<r
$$

where the last inequality follows from the fact that $(y, s) \in T(B)$.
Let $x_{k} \in \Omega_{t}^{c}$ be such that dist $\left(x_{k}, Q_{k}\right) \leqq 8 r_{k}$. It follows that $x_{k} \in \bar{B}=B(x, 10 r)$. In fact,

$$
\begin{aligned}
\left|x-x_{k}\right| & \leqq|x-y|+\left|y-x_{k}^{0}\right|+\left|x_{k}^{0}-x_{k}\right| \\
& \leqq r+r_{k}+10 r_{k}<9 r .
\end{aligned}
$$

Consequently,

$$
\begin{aligned}
\frac{1}{|B|} \int_{T(B)} d\left|\mu_{1}\right| & \leqq \frac{c}{|\bar{B}|} \int_{T(B)} d\left|\mu_{1}\right| \\
& \leqq c C_{1}\left(\mu_{1}\right)\left(x_{k}\right) \leqq c M\left(A_{1}(w)\right)^{*}(t)
\end{aligned}
$$

Finally if $y \notin \Omega_{t}$, then

$$
\frac{1}{|B|} \int_{T(B)} d\left|\mu_{1}\right| \leqq C_{1}\left(\mu_{1}\right)(y) \leqq M\left(A_{1}(w)\right)^{*}(t)
$$

Our analysis shows that if $x \in \Omega_{t}$,

$$
t C_{1}\left(\mu_{1}\right)(x) \leqq c \int_{0}^{t} M\left(A_{1}(w)\right)^{*}(s) d s
$$

Therefore

$$
\begin{equation*}
t\left\|\mu_{1}\right\|_{\tau_{1}^{\infty}} \leqq c \int_{0}^{t} M\left(A_{1}(w)\right)^{*}(s) d s \tag{4.12}
\end{equation*}
$$

Combining (4.11) and (4.12) we obtain

$$
K\left(t, w, \tau_{1}^{1}, \tau_{1}^{\infty}\right) \leqq c \int_{0}^{t} M\left(A_{1}(w)\right)^{*}(s) d s
$$

as desired.
In particular, we have the results of [2] for the ( $V^{\mathbf{0}}, V^{\mathbf{1}}$ ) couple. In fact since $\tau_{1}^{1}=V^{0}, \tau_{1}^{\infty}=V^{1}$, we get from (4.2)

$$
\begin{equation*}
\left(V^{0}, V^{1}\right)_{\theta, \infty}=\tau_{1}^{p_{\theta}, \infty}=V^{1 / p_{\theta}}, \quad \frac{1}{p_{\theta}}=1-\theta \tag{4.13}
\end{equation*}
$$

Moreover, we also have (cf. §2)

$$
\begin{equation*}
\left(V^{0}, V^{1}\right)_{\theta, p_{\theta}}=W^{1 / p_{\theta}}, \quad \frac{1}{p_{\theta}}=1-\theta \tag{4.14}
\end{equation*}
$$

## 5. Spaces of Carleson measures and duality of $T_{\infty}^{p, q}$ spaces

We take up the study of the duals of $T_{\infty}^{p}$ spaces. Recall the result of Coifman, Meyer and Stein [13]: $\left(T_{\infty}^{1}\right)^{*}=\tau_{1}^{\infty}=V^{1}$. In this section we consider more generally, the duals of $T_{\infty}^{p, q}$ spaces.

We begin our analysis with the case $p \geqq 1$.
(5.1) Theorem.
(i) $\left(T_{\infty}^{p, 1}\right)^{*}=V^{1 / p}, \quad 1 \leqq p<\infty$
(ii) $\left(T_{\infty}^{p}\right)^{*}=\tau_{1}^{p^{\prime}}, \quad 1<p<\infty$.

Proof. (i) The basic inequality is a reformulation of (2.5) for $\alpha=1 / p$, namely

$$
\left|\int f(x, t) d w(x, t)\right| \leqq c\|w\|_{V^{1^{\prime} p}}\left\|A_{\infty}(f)\right\|_{p, 1}
$$

It follows that $V^{1 / p} \subset\left(T_{\infty}^{p, 1}\right)^{*}$.
To prove the reverse inclusion we argue as in [13]. Let $l \in\left(T_{\infty}^{p, 1}\right)^{*}$, let $K$ be a compact set of $\overline{\mathbf{R}}_{+}^{n+1}$ and consider the restriction $l_{K}$ of $l$ to $C(K)$; then there exists a measure on $K$ representing $l_{K}$, i.e. $l_{K}(f)=\int f(x, t) d \mu_{K}(x, t) \forall f \in C(K)$; by letting $\left\{K_{n}\right\}$ be an increasing sequence of compact sets covering $\overline{\mathbf{R}}_{+}^{n+1}$ we obtain a measure $\mu$ on $\overline{\mathbf{R}}_{+}^{n+1}$.

We claim that $\mu \in V^{1 / p}$. We may suppose that $\mu \geqq 0$. Then given any open bounded set $\Omega$, let $\left\{f_{n}\right\}$ be a sequence of continuous bounded functions with compact support such that $f_{n} \uparrow \chi_{T(\Omega)}$; then $\left|\int f_{n}(x, t) d \mu\right| \leqq\|I\|\left\|f_{n}\right\|_{T^{p, 1}}$, and therefore

$$
\left|\int_{T(\Omega)} d \mu\right| \leqq\|l\|\left\|A_{\infty}\left(\chi_{T(\Omega)}\right)\right\|_{p, 1}=p\|l\||\Omega|^{1 / p}
$$

(ii) The basic inequality here is the elementary inequality

$$
\begin{equation*}
\left|\int_{\mathbf{R}_{+}^{n+1}} f(x, t) d \mu(x, t)\right| \leqq c \int_{\mathbf{R}^{n}} A_{\infty}(f)(x) A_{1}(\mu)(x) d x \tag{5.2}
\end{equation*}
$$

The proof of (5.2) runs as follows:

$$
\begin{gathered}
\int A_{\infty}(f)(x) A_{1}(\mu)(x) d x=\int_{\mathbf{R}^{n}} \sup _{\Gamma(x)}|f(y, t)| \int_{\Gamma(x)} \frac{d|\mu|(y, t)}{t^{n}} d x \\
\geqq \int_{\mathbf{R}^{n}} \int_{\Gamma(x)} \frac{|f(y, t)|}{t^{n}} d|\mu|(y, t) d x=c \int_{\mathbf{R}^{n}} \int_{0}^{\infty} \int_{|x-y|<t} \frac{|f(y, t)|}{t^{n}} d|\mu|(y, t) d x \\
=c \int_{\mathbf{R}_{+}^{n+1}}|f(y, t)| d|\mu|(y, t)
\end{gathered}
$$

It follows that $\tau_{1}^{p^{\prime}} \subset\left(T_{\infty}^{p}\right)^{*}$. Let $l \in\left(T_{\infty}^{p}\right)^{*}$ and argue as in the proof of (i) to obtain a measure $\mu$ representing $l$. We shall show that $A_{1}(\mu) \in L^{p^{\prime}}$. Assume that $\mu$ is positive.

Let us observe that $A_{1}$ is the adjoint of the operator $P_{t}^{0}$ defined in §2: $P_{t}^{0} f(x)=$ $\frac{1}{t^{n}} \int_{|y| \xi_{t}} f(x-y) d y$. Therefore,

$$
\left\|A_{1}(\mu)\right\|_{p^{\prime}}=\sup _{\|g\|_{p} \leqq 1} \int A_{1}(\mu)(x) g(x) d x=\sup _{\|g\|_{p} \leqq 1} \int P_{t}^{0} g(x) d \mu(x, t)
$$

At this point we use the fact that $\mu$ represents the functional $l$, and obtain

$$
\left\|A_{1}(\mu)\right\|_{p^{\prime}} \leqq \sup _{\|g\|_{p} \leqq 1}\left\|P_{t}^{0} g\right\|_{T_{\infty}^{p}}\|l\| \leqq\|l\| \sup _{\|g\|_{p} \leqq 1}\|M g\|_{p} \leqq c\|l\|
$$

where in the last inequality we use the Hardy-Littlewood maximal theorem.
To deal with the case $0<p<1$ we use the atomic theory of $T_{\infty}^{p}$ spaces as developed in [13]. A $p$-atom is a function $a(x, t)$ such that there exists a ball $B$ with supp $a \subset T(B)$ and $\|a\|_{\infty} \leqq|B|^{-1 / p}$. Observe that if a is a $p$-atom then $\|a\|_{T_{\infty}^{p}} \leqq 1$. Moreover, notice that these atoms do not have zero mean.
(5.3) Theorem (cf. [13]). Suppose that $f \in T_{\infty}^{p}, \quad 0<p \leqq 1$, then $f=\sum_{j=1}^{\infty} \lambda_{j} a_{j}$ in the sense of $\mathscr{S}^{\prime}$, where the $a_{j}^{\prime}$ 's are p-atoms, $\lambda_{j} \in \mathbf{C}$, and $\left(\sum_{j=1}^{\infty}\left|\lambda_{j}\right|^{p}\right)^{1 / p} \leqq c\|f\|_{T_{\infty}^{p}}$. In fact, $\|f\|_{T_{\infty}^{p}} \approx \inf \left\{\left(\sum_{j=1}^{\infty}\left|\lambda_{j}\right|^{p}\right)^{1 / p}: f=\sum_{j=1}^{\infty} \lambda_{j} a_{j}, a_{j}\right.$ p-atom $\}$.

Finally one more observation is needed before stating the duality theorem. It is easy to check, using for example Whitney's decompositions, that $w \in V^{r}$ with $r \geqq 1$ if and only if for some $c>0$ and all balls $B \subset \mathbf{R}^{m}$,

$$
\begin{equation*}
|w|(T(B)) \leqq c|B|^{r} \tag{5.4}
\end{equation*}
$$

(5.5) Theorem. Suppose $0<p<1$, then

$$
\begin{equation*}
\left(T_{\infty}^{p}\right)^{*}=V^{1 / p} \tag{5.6}
\end{equation*}
$$

Proof. Let a be a $p$-atom, supp $a \subset T(B)$, for some ball $B \subset R^{n}$, then if $\mu \in V^{1 / p}$

$$
\left|\int_{T(B)} a(x, t) d \mu(x, t)\right| \leqq|B|^{-1 / p} \int_{T(B)} d|\mu|(x, t) \leqq\|\mu\|_{V^{1 / p}}
$$

Therefore, by (5.3), we see that $V^{1 / p} \subset\left(T_{\infty}^{p}\right)^{*}$. The reverse inclusion is obtained as in the proof of the case $p \geqq 1$ above using the remark (5.4).

## 6. Applications

We consider applications of our results to some problems in harmonic analysis.

## A. Multiplier operators and maximal operators on $V^{\alpha}, \tau_{1}^{p, q}$ and $T_{\infty}^{p}$ spaces

Let $X$ denote any of the spaces $V^{\alpha}, \tau_{1}^{p, q}$ or $T_{\infty}^{p}$. A natural problem that arises is the study of multipliers of $X$ : i.e., let $\mu(x, t)$ be a function on $\mathbf{R}_{+}^{n+1}$, then under what conditions is the map $w \rightarrow \mu w$ well defined and continuous from $X$ into itself?

Observe that when $X=T_{\infty}^{p}$ a necessary condition is the continuity of $\mu$. In other cases the necessity of this condition is not so explicit but is natural in view of our next.
(6.1) Lemma (cf. [18]). Let $P=\left(x_{0}, t_{0}\right) \in \mathbf{R}_{+}^{n+1}$ be fixed and let $\delta_{P}$ be the Dirac measure concentrated in $P$. Then, $\delta_{P} \in V^{\alpha}, \delta_{P} \in \tau_{1}^{p, q}$. Moreover,
(i) $\left\|\delta_{P}\right\|_{V^{\alpha}}=\frac{c}{t_{0}^{n \alpha}}, \quad c=c(n, \alpha)$
(ii) $\left\|\delta_{P}\right\|_{\tau_{1}^{p, g}}=c t_{0}^{n((1 / p)-1)}, \quad c=c(n, p, q)$.

Proof. Let $\Omega \subset \mathbf{R}^{n}$ be a bounded open set, then $\delta_{P}(T(\Omega))$ vanishes unless $\left(x_{0}, t_{0}\right) \in T(\Omega)$. Thus, if $\delta_{P}(T(\Omega)) \neq 0$ we see that $c t_{0}^{n} \leqq|\Omega|$, and consequently, for $\alpha>0$

$$
\delta_{P}(T(\Omega)) \leqq \frac{1}{\left(c t_{0}^{n}\right)^{\alpha}}|\Omega|^{\alpha}
$$

On the other hand $\delta_{P}\left(T\left(B\left(x_{0}, t_{0}\right)\right)\right)=1=\left|B\left(x_{0}, t_{0}\right)\right|^{\alpha} /\left(c t_{0}^{n}\right)^{\alpha}$. These two observations prove (i).

The second half of the lemma follows readily from the observation that $A_{1}\left(\delta_{P}\right)(x)=t_{0}^{-n} \chi_{B\left(x_{0}, t_{0}\right)}(x)$.

Let us denote by $M_{X}$ the space of multipliers of $X$ as a subset of $C\left(\mathbf{R}_{+}^{n+1}\right)$. Let $B\left(\mathbf{R}_{+}^{n+1}\right)$ be the space of bounded continuous functions on $\mathbf{R}_{+}^{n+1}$. We then have
(6.2) Theorem. $M_{X}=B\left(\mathbf{R}_{+}^{n+1}\right)$.

Proof. It is clear that $B\left(\mathbf{R}_{+}^{n+1}\right) \subseteq M_{X}$. Now, if $\mu$ is a multiplier of $V^{\alpha}$, then according to (6.1),

$$
\mid \mu\left(x_{0}, t_{0}\right)\| \| \delta_{P}\left\|_{V^{x}}=\right\| \mu \delta_{P}\left\|_{V^{z}} \leqq C_{\mu}\right\| \delta_{P} \|_{V^{x}}
$$

Thus, $\mu \in B\left(\mathbf{R}_{+}^{n+1}\right)$.
The same argument can be used to show that any continuous multiplier of $\tau_{1}^{p, q}$ must be bounded. We also readily see that $M_{T_{\infty}^{p}}=B\left(\mathbf{R}_{+}^{n+1}\right)$.

We introduce now two weighted maximal functions, which are variants of the one considered in [13].

Let $\mu$ be a function defined on $\mathbf{R}_{+}^{n+1}$. Given $f \in T_{\infty}^{p}, 0<p \leqq 1$, let

$$
N_{\mu}(f)(x)=\sup _{t>0}|\mu(x, t) f(x, t)| .
$$

If $\varphi$ is a function in the Schwartz class $S\left(\mathbf{R}^{n}\right)$ such that $\int \varphi(x) d x=1$, let $\varphi_{t}(x)=$ $t^{-n} \varphi(x / t)$. Given $f \in H^{p}\left(\mathbf{R}^{n}\right), 0<p \leqq 1$, define

$$
N_{\mu}^{\varphi}(f)(x)=\sup _{t \geqq 0}\left|\mu(x, t)\left(\varphi_{t} * f\right)(x)\right| .
$$

(6.3) Theorem. Let $0<p \leqq 1$. The following conditions are equivalent
(i) $N_{\mu}: T_{\infty}^{p} \rightarrow L^{p}$, boundedly.
(ii) $\sup 1 /|B| \int_{B}\left(\sup _{t<r_{B}}|\mu(x, t)|\right)^{p} d x<\infty$, where the sup is taken over all balls and $r_{B}=$ radius of $B$.

Proof. Suppose that (i) holds. Given a ball $B$, let $f$ be a continuous function such that $\chi_{T(B)} \leqq f \leqq \chi_{T(2 B)}$ where $2 B$ denotes the ball with the same center as $B$ and radius $=$ $2 r_{B}$. We have,

$$
\chi_{B}(x) \sup _{t<r_{B}}|\mu(x, t)|=\sup _{t>0}\left|\mu(x, t) \chi_{T(B)}(x, t)\right| \leqq N_{\mu}(f)(x) .
$$

Thus,

$$
\int_{B}\left(\sup _{t<r_{B}}|\mu(x, t)|\right)^{p} d x \leqq c_{\mu}\left\|\chi_{T(2 B)}\right\|_{T_{\infty}^{p}}^{p}=c_{\mu}|B|
$$

which is (ii).
Conversely, suppose that $\mu$ satisfies (ii). Let $f$ be a $p$-atom supported in $T(B)$, for some ball $B$. Then,

$$
\begin{aligned}
N_{\mu}(f)(x) & \leqq \frac{1}{|B|^{1 / p}} \sup _{t>0}\left|\mu(x, t) \chi_{T(B)}(x, t)\right| \\
& \leqq \frac{\chi_{B}(x)}{|B|^{1 / p}} \sup _{t<r_{B}}|\mu(x, t)| .
\end{aligned}
$$

Consequently,

$$
\left\|N_{\mu}(f)\right\|_{p}^{p} \leqq \sup _{B} \frac{1}{|B|} \int_{B}\left(\sup _{t<r_{B}}|\mu(x, t)|\right)^{p} d x
$$

and (i) follows.
(6.4) Corollary. Suppose that $\mu$ satisfies condition (ii) of (6.9). Then, $N_{\mu}^{\varphi}$ defines a bounded operator, $N_{\mu}^{\varphi}: H^{p} \rightarrow L^{p}, 0<p \leqq 1$.

Proof. Indeed, if $f \in H^{p}$, then $\varphi_{t} * f \in T_{\infty}^{p}$ (see [16], p. 183). Consequently, $N_{\mu}^{\varphi}(f)=$ $N_{\mu}\left(\varphi_{t} * f\right)$ and (6.4) follows from (6.3).

## B. Interpolation of $H^{p}$ spaces and Lipschitz spaces

As pointed out in [13] the theory of tent spaces can be used to derive the interpolation theory (real and complex) of $H^{p}$ spaces. In fact, using a "trick" of A. P. Calderón (cf. [6]) we can "invert" the $\pi_{\varphi}$ operators (cf. [13]). Let $\varphi \in C_{0}^{\infty}\left(\mathbf{R}^{n}\right)$ be radial satisfying all the conditions of [13, p. 328]. Moreover, suppose $\varphi$ satisfies $\int \varphi(\xi t)|\xi| e^{-2 \pi|\xi| t} d t=-\frac{1}{2 \pi}, \forall \xi \neq 0$. Then for sufficiently regular $f, \pi_{\varphi}\left(t \frac{d}{d t} P_{t} f\right)=f$, where $P_{t} f=$ Poisson integral of $f$. (This can be easily checked using Fourier trans-
forms). Now, one observes that $f \in H^{p_{\leftrightarrow}} t \frac{d}{d t}\left(P_{t} f\right) P_{t} f \in T_{2}^{p}$ to obtain the interpolation theorems for $H^{p}$ spaces from the ones for $T_{2}^{p}$ spaces.

The complex interpolation theory for the $T_{\infty}^{p}$ and $\tau_{\infty}^{p}$ spaces can be readily derived from the results of $\S 3, \S 4$ and the methods of [20].
(6.5) Theorem. Let $1<p<p_{0}<\infty, 0<\theta<1$, then
(i) $\left[T_{\infty}^{p}, T_{\infty}^{p_{0}}\right]_{\theta}=T_{\infty}^{p_{\theta}}, \frac{1}{p_{\theta}}=\frac{1-\theta}{p}+\frac{\theta}{p_{0}}$.
(ii) $\left[\tau_{1}^{p}, \tau_{1}^{p_{0}}\right]_{\theta}=\tau_{1}^{p_{\theta}}, \frac{1}{p_{\theta}}=\frac{1-\theta}{p}+\frac{\theta}{p_{0}}$.

Proof. We shall only provide the details of the proof of (ii). (The proof of (i) is similar.)

Using (i), and our duality theorem (5.1) (ii) we have

$$
\tau_{1}^{p_{\theta}}=\left[T_{\infty}^{p^{\prime}}, T_{\infty}^{p_{0}^{\prime}}\right]_{\theta}^{*}=\left[\tau_{1}^{p}, \tau_{1}^{P}\right]_{\theta}
$$

## C. Balayages

Let $P_{t}$ be a kernel on $\mathbf{R}^{n} \times \mathbf{R}^{n}$ (see §2). The balayage $P * w$ is defined formally by (cf. [2])

$$
\begin{equation*}
(P * w)(y)=\int_{\mathbf{R}_{+}^{n+1}} P_{t}(x, y) d w(x, t) \tag{6.6}
\end{equation*}
$$

where $w$ is a measure on $\overline{\mathbf{R}}_{+}^{n+1}$.
Therefore, we have the duality

$$
\begin{equation*}
\int_{\mathbf{R}_{+}^{n+1}} P_{t} f(x) d w(x, t)=\int_{\mathbf{R}^{n}} f(y)(P * w)(y) d y \tag{6.7}
\end{equation*}
$$

In particular, if $P_{t} f$ is the Poisson integral of $f$, we have

$$
\begin{equation*}
P_{t}: H^{p} \rightarrow T_{\infty}^{p}, \quad 0<p \leqq 1 \tag{6.8}
\end{equation*}
$$

and by duality,

$$
\begin{equation*}
P_{*}: V^{1 / p} \rightarrow \operatorname{Lip} n\left(\frac{1}{p}-1\right), \quad 0<p<1 \tag{6.9}
\end{equation*}
$$

where the Lip space is modulo polynomials of degree $\left[n\left(\frac{1}{p}-1\right)\right]$. More precisely, in order to give a sense to (6.6) it is necessary in general to suppose that the measure $w$ is finite.

Let us observe that (6.9) also holds for any kernel $P_{t}$ satisfying (6.8). In particular, we will give a direct proof of the following result.
(6.10) Theorem. Let $P_{t}$ be a kernel on $\mathbf{R}^{n} \times \mathbf{R}^{n}$ satisfying for some $N=1,2,3, \ldots$, the following conditions:
(i) $\quad\left|D_{y}^{\beta} P_{t}(x, y)\right| \leqq c t(t+|x-y|)^{-n-1-|\beta|}, \quad|\beta| \leqq N-1$
(ii) $\left|D_{y}^{\beta} P_{t}(x, y)-\sum_{|\gamma+\beta|<N} D_{y}^{\gamma+\beta} P_{t}\left(x, y_{0}\right) \frac{\left(y-y_{0}\right)^{\gamma}}{\gamma!}\right| \leqq c \frac{\left|y-y_{0}\right|^{N}}{\left(t+\left|x-y_{0}\right|\right)^{n+N+|\beta|}}$

$$
\text { if } t+\left|x-y_{0}\right|>2\left|y-y_{0}\right|,|\beta| \leqq N-1
$$

Let $w \in V^{\alpha}, \alpha \geqq 1$ be a finite measure. Then, if $1+\frac{N-1}{n}<\alpha<1+\frac{N}{n}$, the $b a$ layage $P_{*} w$ of $w$ belongs to $\operatorname{Lip} n(\alpha-1)$.

Proof. Let us first show that $P_{*} w$ is well defined a.e. In fact, condition (6.10) (i) with $\beta=0$ implies that $\sup _{x, t} \int\left|P_{t}(x, y)\right| d y<\infty$. Since $w$ is a finite measure, we deduce that

$$
\int_{\mathbf{R}^{n}}\left|P^{*} w(y)\right| d y \leqq|w|\left(\mathbf{R}_{+}^{n+1}\right)
$$

which implies that $P * w$ is finite a.e. Now, given $|\beta|=N-1$, we will show that $D^{\beta} P * w$ is a bounded function satisfying a Lipschitz condition

$$
\left|D^{\beta}\left(P_{*} w\right)(y)-D^{\beta}\left(P_{*} w\right)\left(y_{0}\right)\right| \leqq c\left|y-y_{0}\right|^{n(\alpha-1)-N+1}
$$

If $B$ denotes $B(y, 1)$, we see that (see [2], p. 33),

$$
\left|\int_{T(B)} D_{y}^{\beta} P_{t}(x, y) d w(x, t)\right|+\left|\int_{T(B)} D_{y}^{\beta} P_{t}(x, y) d w(x, t)\right|<\infty
$$

According to (6.10) (i), since $\alpha n>n+N-1$, we can appeal to Lemma 2(1) in [2] to deduce that the first term is bounded. In $T(B)^{c}$ we have $t+|x-y|>1$; thus, the second integral can be majorized by $c|w|\left(\mathbf{R}_{+}^{n+1}\right)<\infty$.

Finally, let us verify the Lipschitz condition. Fix $y, y_{0}$ and let $B=B\left(y_{0}, 2\left|y-y_{0}\right|\right)$. Then, we estimate

$$
\begin{gathered}
\int_{T(B)}\left[D_{y}^{\beta} P_{t}(x, y)-D_{y}^{\beta} P_{t}\left(x, y_{0}\right)\right] d w(x, t)+\int_{T(B) c}\left[D_{y}^{\beta} P_{t}(x, y)-\right. \\
\left.-D_{y}^{\beta} P_{t}\left(x, y_{0}\right)\right] d w(x, t)=(I)+(J)
\end{gathered}
$$

The term ( $I$ ) can be majorized by the sum of two integrals over tents

$$
\begin{aligned}
(I) & \leqq \int_{T\left(B\left(y, 3\left|y-y_{0}\right|\right)\right)}\left|D_{y}^{\beta} P_{t}(x, y)\right| d|w|(x, t) \\
& +\int_{T(B)}\left|D_{y}^{\beta} P_{t}\left(x, y_{0}\right)\right| d|w|(x, t)
\end{aligned}
$$

and each of these integrals can be estimated in the same way using (6.10) (i) and Lemma 2 in [2]. In fact, let us consider the first one, say ( $I_{1}$ ).

$$
\left(I_{1}\right) \leqq c \int_{T\left(B\left(y, 3\left|y-y_{0}\right|\right)\right)}(t+|x-y|)^{-n-N+1} d|w|(x, t)
$$

Since $\alpha n>n+N-1$, we are again in the conditions of Lemma 2(2) in [2] and consequently the last integral is bounded by a constant times $\left|y-y_{0}\right|^{n(\alpha-1)-N+1}$.

To estimate $(J)$, we observe that over $T(B)^{c}$ we have $t+\left|x-y_{0}\right|>2\left|y-y_{0}\right|$ and therefore, we can use (6.10) (ii), to get

$$
(J) \leqq c\left|y-y_{0}\right|^{N} \int_{T(B)^{c}}\left(t+\left|x-y_{0}\right|\right)^{-n-2 N+1} d|w|(x, t)
$$

Since this time $\alpha n<n+N \leqq n+2 N-1$, we can apply Lemma 2(1) in [2], to obtain $(J) \leqq c\left|y-y_{0}\right|^{(\alpha-1) n-N+1}$. This completes the proof of the theorem.
(6.11) Remark. When $\frac{1}{p}$ in (6.9) has the critical values $\frac{N+n}{n}, N=1,2, \ldots$, one gets that $D^{\beta}(P * w),|\beta|=N-1$, belongs to the Zygmund space Lip. 1 However, when $N \geqq 2$, the same technique of Theorem (6.10) gives the improvement

$$
\left|D^{\beta}(P * w)(y)-D^{\beta}(P * w)\left(y_{0}\right)\right| \leqq c\left|y-y_{0}\right| .
$$

Let us observe that the hypothesis that $w$ is finite can be dropped in (6.10) by considering a regularization of the balayage (see [18]). More precisely, if $1+\frac{N-1}{n}<$ $\alpha<1+\frac{N}{n}$, we consider for $|\beta| \leqq N-1$,

$$
\int\left[D_{y}^{\beta} P_{t}(x, y)-\sum_{|\beta+y|<N} D_{y}^{\beta+\gamma} P_{t}\left(x, y_{0}\right) \frac{\left(y-y_{0}\right)^{\gamma}}{\gamma!} \chi_{r(\bar{B})^{c}}(x, t)\right] d w(x, t)
$$

where $B=B\left(y_{0}, r_{0}\right), \quad \bar{B}=B\left(y_{0}, 2 r_{0}\right)$. This integral can be seen to be finite for a.e. $y \in B$. In fact,

$$
\int_{T(B)}\left|D_{y}^{\beta} P_{t}(x, y)\right| d|w|(x, t) \leqq c \int_{T\left(B\left(y, 3 r_{0}\right)\right)}(t+|x-y|)^{-n-|\beta|} d|w|(x, t)
$$

Since $\alpha n>n+N-1 \geqq n+|\beta|$, we can apply Lemma 2(2) in [2] to estimate it by $c r_{0}^{(\alpha-1) n-|\beta|}$. On the other hand,

$$
\begin{gathered}
\left|\int_{T(B)^{c}}\left[D_{y}^{\beta} P_{t}(x, y)-\sum_{|\beta+y|<N} D_{y}^{\beta+\gamma} P_{t}\left(x, y_{0}\right) \frac{\left(y-y_{0}\right)^{\gamma}}{\gamma!}\right] d w(x, t)\right| \\
\quad \leqq c \int_{T(B) c}\left|y-y_{0}\right|^{N}\left(t+\left|x-y_{0}\right|\right)^{-n-N-|\beta|} d|w|(x, t)
\end{gathered}
$$

This time we have $\alpha n<n+N \equiv n+N+|\beta|$. Thus, applying Lemma 2(1) in [2], the last term is bounded by

$$
c\left|y-y_{0}\right|^{N} r_{0}^{(\alpha-1) n-N-|\beta|} .
$$

Taking $r_{0}=\left|y-y_{0}\right|$, these calculations also show that the increment can be majorized by $c\left|y-y_{0}\right|^{(\alpha-1) n-N+1}$, when $|\beta|=N-1$.

## 7. Carleson measures on product spaces

The theory developed in this paper can be easily adapted to the study of certain tent spaces and spaces of Carleson measures on product spaces. A theory of Carleson measures on product spaces has been developed by L. Carleson, S. Y. A. Chang, R. Fefferman (cf. [8], [9], [10], [11]).

In this section we shall outline an extension of some of our results to the two parameter setting. A more detailed theory of tent spaces on product spaces will be developed elsewhere,

We will work on $\mathbf{R}_{+}^{m} \times \mathbf{R}_{+}^{n}=\left\{(x, t) \mid x=\left(x_{1}, x_{2}\right) \in \mathbf{R}^{m} \times \mathbf{R}^{n}, t=\left(t_{1}, t_{2}\right) \in \mathbf{R}^{2}, t_{1}, t_{2}>0\right\}$. Given $(x, t) \in \mathbf{R}_{+}^{m} \times \mathbf{R}_{+}^{n}, B(x, t)$ stands for the product of balls $B\left(x_{1}, t_{1}\right) \times B\left(x_{2}, t_{2}\right) \subset$ $\mathbf{R}^{m} \times \mathbf{R}^{n}$. Moreover, given $x \in \mathbf{R}^{m} \times \mathbf{R}^{n}, \Gamma(x)$ will denote the product of cones $\Gamma\left(x_{1}\right) \times$ $\Gamma\left(x_{2}\right) \subset \mathbf{R}_{+}^{m} \times \mathbf{R}_{+}^{n}$.

Given a subset $\Omega \subset \mathbf{R}^{m} \times \mathbf{R}^{n}$, the tent over $\Omega, T(\Omega)$, is defined as

$$
\left\{(x, t) \in \mathbf{R}_{+}^{m} \times \mathbf{R}_{+}^{n} \mid B(x, t) \subset \Omega\right\} .
$$

(7.1) Definition. In accordance with the one parameter case, we define

$$
A_{\infty}(f)(x)=\sup _{r(x)}|f(y, t)|
$$

For $0<p<\infty, 0<r \leqq \infty, T_{\infty}^{p, r}$ denotes the class of functions $f$ which are continuous in $\mathbf{R}_{+}^{m} \times \mathbf{R}_{+}^{n}$, for which $A_{\infty}(f)(x) \in L(p, r)$ and $\left\|A_{\infty}\left(f_{\varepsilon}-f\right)\right\|_{p, r} \rightarrow 0$, as $\varepsilon \rightarrow 0$, where $f_{\varepsilon}(x, t)=f(x, t+\varepsilon), \varepsilon=\left(\varepsilon_{1}, \varepsilon_{2}\right), \varepsilon_{1}, \varepsilon_{2}>0$. These spaces are normed spaces for $1<p<\infty, 1 \leqq r \leqq \infty$.

As in the one parameter case, we have the inequality

$$
|f(x, t)| \leqq \inf _{y \in B(x, t)} A_{\infty}(f)(y) .
$$

By integration on $B(x, t)$, we get

$$
\begin{equation*}
|f(x, t)| \leqq C t_{1}^{-m / p} t_{2}^{-n / p}\left\|A_{\infty}(f)\right\|_{p} \tag{7.2}
\end{equation*}
$$

This shows that $T_{\infty}^{p}$ is a complete space.
The spaces $T_{\infty}^{p}, 0<p \leqq 1$, can be characterized through an atomic decomposition. First, a suitable notion of atom is needed.

A function $a(y, t)$ is a $p$-atom if there exists an open subset $\Omega \subset \mathbf{R}^{m} \times \mathbf{R}^{n}$ of finite measure such that
(i) $\operatorname{supp}(a) \subset T(\Omega)$,
(ii) $\|a\|_{\infty} \leqq|\Omega|^{-1 / p}$.

It is readily seen that a $p$-atom a belongs to $T_{\infty}^{p}$ and $\left\|A_{\infty}(a)\right\|_{p} \leqq 1$.
(7.3) Theorem. Given a function $f$ defined in $\mathbf{R}_{+}^{m} \times \mathbf{R}_{+}^{n}$, the following statements are equivalent:
a) $f \in T_{\infty}^{p}$,
b) $f=\sum \lambda_{k} a_{k}$ in the sense of $\mathscr{D}^{\prime}\left(\mathbf{R}_{+}^{m} \times \mathbf{R}_{+}^{n}\right)$, where $a_{k}$ are p-atoms, $\sum\left|\lambda_{k}\right|^{p} \leqq$ $c\left\|A_{\infty}(f)\right\|^{p}$.
Moreover, $\|f\|_{T_{\infty}^{p}}^{p} \approx \inf \left\{\Sigma\left|\lambda_{k}\right|^{p}: f=\sum \lambda_{k} a_{k}\right\}$, as in the one parameter case.
Proof. We only need to show a) $\Rightarrow$ b). First observe that given a function $f$ defined in $\mathbf{R}_{+}^{m} \times \mathbf{R}_{+}^{n}$, we have for each $\lambda>0$,

$$
\begin{equation*}
\{(x, t) \| f(x, t) \mid>\lambda\} \subset T\left\{x \mid A_{\infty}(f)(x)>\lambda\right\} . \tag{7.4}
\end{equation*}
$$

Now, given $f \in T_{\infty}^{p}$, let us consider for each $k \in \mathbf{Z}, O_{k}=\left\{x \mid A_{\infty}(f)>2^{k}\right\} . O_{k}$ is an open subset of $\mathbf{R}^{m} \times \mathbf{R}^{n}$ of finite measure and $O_{k+1} \subset O_{k}$. Let $\mathcal{O}=U_{k} T\left(O_{k}\right)$. From (7.4) it follows that $\operatorname{supp}(f) \subset \mathcal{O}$. Now, let $\mathcal{O}_{k}=T\left(O_{k}\right) \backslash T\left(O_{k+1}\right)$. The sets $\mathcal{O}_{k}$ are a disjoint covering of supp $(f)$. Moreover, on $\mathcal{O}_{k}$ we have $|f(x, t)| \leqq 2^{k+1}$. Thus, we can write $f=\sum \lambda_{k} a_{k}$, where $\lambda_{k}=2^{k+1}\left|O_{k}\right|^{1 / p}, a_{k}=2^{-k-1}\left|O_{k}\right|^{1 / p} \chi_{\mathcal{\theta}_{k}} f$. It is easy to see that $a_{k}$ is a $p$-atom supported on $\mathcal{O}_{k}$. Moreover, let $\Delta_{j}=O_{j} \backslash O_{j+1}$. Then, $O_{k}=\bigcup_{j=k}^{\infty} \Delta_{j}$ and

$$
\begin{aligned}
\sum_{k=-\infty}^{\infty} \lambda_{k}^{p} & =2^{p} \sum_{k=-\infty}^{\infty} 2^{p k}\left|O_{k}\right| \\
& =2^{p} \sum_{k=-\infty}^{\infty} 2^{p k} \sum_{j=k}^{\infty}\left|A_{j}\right| \\
& \leqq 2^{p} \sum_{k=-\infty}^{\infty} 2^{p k} \sum_{j=k}^{\infty} 2^{-p j} \int_{\Delta_{j}} A_{\infty}(f)(x)^{p} d x \\
& =2^{p} \sum_{j=-\infty}^{\infty} 2^{-p j} \sum_{k=-\infty}^{j} 2^{p k} \int_{\Delta_{j}} A_{\infty}(f)(x)^{p} d x \\
& =c \sum_{j=-\infty}^{\infty} \int_{\Delta} A_{\infty}(f)(x)^{p} d x \\
& =c\left\|A_{\infty}(f)\right\|_{L^{p}} .
\end{aligned}
$$

This completes the proof of the theorem.

We want now to identify the duals of the spaces $T_{\infty}^{p}, T_{\infty}^{p, 1}$. In analogy with the one parameter case, a suitable notion of Carleson measure in the product space $\overline{\mathbf{R}}_{+}^{m} \times \overline{\mathbf{R}}_{+}^{n}$ plays an important role.
(7.5) Definition. Given a measure $w$ in $\overline{\mathbf{R}}_{+}^{m} \times \overline{\mathbf{R}}_{+}^{n}$, the $A_{1}$ functional is defined as

$$
A_{1}(w)(x)=\int_{\Gamma(x)} \frac{d|w|(x, t)}{t_{1}^{m} t_{2}^{n}}
$$

For $\alpha \geqq 0$, we also define

$$
C_{1, \alpha}(w)(x)=\sup \frac{1}{|\Omega|^{\alpha}} \int_{T(\Omega)} d|w|(y, t)
$$

where the supremum is taken over all the open subsets of $\mathbf{R}^{m} \times \mathbf{R}^{n}$ of finite measure containing $x$. For $0<p, \alpha<\infty$, let $\tau_{1}^{p}=\left\{w \mid A_{1}(w) \in L^{p}\right\}, V^{\alpha}=\left\{w \mid C_{1, \alpha}(w) \in L^{\infty}\right\}$.

According to [8], $V^{1}$ is the space of Carleson measures. If $P_{t}$ denotes the tensor product of the usual Poisson kernels in $\mathbf{R}_{+}^{m}$ and $\mathbf{R}_{+}^{n}$, we have the following characterization of $V^{1}$. A measure $w \in V^{1}$ if and only if (see [19])

$$
\begin{equation*}
\int_{\mathbf{R}_{+}^{m} \times \mathbf{R}_{+}^{n}} \mid P_{t}(f)\left((y)^{p} d|w|(y, t) \leqq c_{p} \int_{\mathbf{R}^{m} \times \mathbf{R}^{n}}|f(y)|^{p} d y \quad 1 \leqq p<\infty\right. \tag{7.6}
\end{equation*}
$$

The functionals $A_{\infty}$ and $C_{1, \alpha}$ are related by the following inequality

$$
\begin{equation*}
\int|f(y, t)|^{1 / \alpha} d|w|(y, t) \leqq \int A_{\infty}(f)(x) C_{1, \alpha}(w)(x)^{1 / \alpha} d x \tag{7.7}
\end{equation*}
$$

In fact, the proof follows readily from (7.4).
In the spirit of (7.6) we also have a characterization of $V^{\alpha}, 0<\alpha<1$. In fact, a measure $w \in V^{\alpha}$ if and only if

$$
\begin{equation*}
\int\left|P_{t}(f)(y)\right| d|w|(y, t) \leqq c\|f\|_{1 / \alpha, 1} \tag{7.7}
\end{equation*}
$$

Finally, a measure $w \in \tau_{1}^{p}, 1<p<\infty$, if and only if

$$
\begin{equation*}
\int_{\mathbf{R}_{+}^{m} \times \mathbf{R}_{+}^{n}}\left|P_{t}(f)(y)\right| d|w|(y, t) \leqq c\left(\int_{\mathbf{R}^{m} \times \mathbf{R}^{n}}|f(x)|^{p^{\prime}} d x\right)^{1 / p^{\prime}} \tag{7.8}
\end{equation*}
$$

All the above assertions follow readily as in the one parameter case.

## (7.9) Theorem.

a) Let $0<p \leqq 1$, then $\left(T_{\infty}^{p}\right)^{*}=V^{1 / p}$.

Let $1<p<\infty$, then
b) $\left(T_{\infty}^{p}\right)^{*}=\tau_{1}^{p^{\prime}}$.
c) $\left(T_{\infty}^{p, 1}\right)^{*}=V^{1 / p}$.

Proof. The only change needed in the theory of $\S 5$ is to replace in the proof of (5.1) (ii) the maximal operator of Hardy Littlewood by the strong maximal operator.

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