# Pointwise regularity of solutions to nonlinear double obstacle problems 

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## 1. Introduction

In this paper we investigate the continuity at a given point $x_{0}$ of the solution $u$ of a nonlinear elliptic variational inequality with a double obstacle constraint of the form $\psi_{1} \leqq u \leqq \psi_{2}$. The partial differential operators associated with our obstacle problem are quasi-linear and include operators with only $L^{\infty}$ coefficients. The obstacles in this context are to be regarded as quite general and irregular. In particular, they may be discontinuous. Since the partial differential operators associated with our problem may have only $L^{\infty}$ coefficients, we can expect at most Hölder continuity for the regularity of our solution. Indeed, we show that if both obstacles are locally Hölder continuous, then the solution is also locally Hölder continuous. We also show that if the obstacles are not continuous, but satisfy a Wiener-type regularity condition, the solution is still continuous. This work extends that of [MZ1] in which a similar investigation was undertaken for the case of a single obstacle. This work also extends the recent paper of [DMV] which is devoted to the double obstacle problem for linear operators with bounded measurable coefficients. Because their work involves linear operators, they are able to employ potential theoretic techniques to obtain many of their estimates. These techniques are not available for us in our context of nonlinear operators.

Since the one-obstacle problem is a special case of the two-obstacle problem, one cannot expect better results in this latter case. On the other hand, the twoobstacle problem is so similar to the unilateral case, that one would anticipate virtually identical results. However, we have not been able to achieve this with the general structure we consider. One significant difference between the single and

[^0]double-obstacle problem is that in the former case, the solution to the obstacle problem turns out to be a supersolution of the corresponding differential equation whereas this is no longer true in the double obstacle situation. This does not allow the use of the weak Harnack inequality for supersolutions, which was a critical tool in the analysis of the unilateral problem. If we impose more conditions on the structure of our nonlinear problem, then we show that the solution does become a supersolution of the corresponding equation and in this case we obtain results that run parallel to the single obstacle problem. The conditions imposed on the structure are general enough to include the case of linear operators.

Let $A$ and $B$ denote, respectively, vector and scalar valued Borel functions defined on $\Omega \times R^{1} \times R^{n}$ satisfying the following structure conditions for fixed $n>$ $p>1, \mu>0$, and $v \geqq 0$ :

$$
\begin{gather*}
|A(x, z, h)| \leqq \mu|h|^{p-1}+\mu|z|^{p-1}+v \\
h \cdot A(x, z, h) \geqq|h|^{p}-\mu|z|^{p}-v  \tag{1.1}\\
|B(x, z, h)| \leqq \mu|h|^{p-1}+\mu|z|^{p-1}+v
\end{gather*}
$$

for a.e. $x \in \Omega$ and all $z \in R^{1}, h \in R^{n}$.
The obstacles $\psi_{1}$ and $\psi_{2}$ are defined to be real-valued functions defined on $\Omega$ with $\psi_{1} \leqq \psi_{2}$. In addition, we assume that $\psi_{1}$ is bounded above and that $\psi_{2}$ is bounded below. A bounded function $u \in W^{1, p}(\Omega)$ is said to be a solution of the double obstacle problem if $\psi_{1} \leqq u \leqq \psi_{2}$ and

$$
\begin{equation*}
\int_{\Omega} A(x, u, \nabla u) \cdot \nabla \varphi d x+\int_{\Omega} B(x, u, \nabla u) \varphi d x \geqq 0 \tag{1.2}
\end{equation*}
$$

for all $\varphi \in W_{0}^{1, p}(\Omega)$ such that $\psi_{1} \leqq u+\varphi \leqq \psi_{2}$.
Although we have made the assumption that $1<p<n$, it can be shown without too much difficulty that virtually all of the results below are valid in case $p=n$.

## 2. Preliminaries

Throughout the paper, we will utilize properties of the following classes of functions introduced in [MZ1] which are similar but not quite as general as the De Giorgi classes discussed in [LU]. If $G \subset R^{n}$ is an open set, $C>0$ and $\lambda \geqq 0$, we let

$$
S^{e}(G, p, C, \lambda)
$$

denote the set of all nonnegative functions $v \in W^{1, p}(G)$ such that

$$
\begin{gather*}
\int_{G}\left|\nabla(v(x)-k)^{f}\right| \eta(x)^{p} d x \leqq C \int_{G}\left[(v-k)^{)^{p}\right]^{p}\left[\eta(x)^{p}+|\nabla \eta(x)|^{p}\right] d x}+\quad+C\left(k^{p}+\lambda\right) \int_{((\rho(x)-k) \leq>0\}} \eta(x)^{p} d x\right. \tag{2.1}
\end{gather*}
$$

for all $k \geqq 0$ and all nonnegative Lipschitz functions $\eta$ on $R^{n}$ with compact support in $G$. Here $\varepsilon$ stands for either + or - .

The next two lemmas reveal the relationship between the solution $u$ of the double obstacle problem and the classes $S^{e}(G, p, C, \lambda)$.

Lemma 2.1. Let $u$ be a solution of (1.2) and let $d_{0} \geqq 0$. If $|d| \leqq d_{0}$, then there exist constants $C>0$ and $\lambda>0$ such that
(i) If $d \geqq \psi_{1}$ in $G \subset \Omega$, then $(u-d)^{+} \in S^{+}(G, p, C, \lambda)$
(ii) If $d \leqq \psi_{2}$ in $G \subset \Omega$, then $(u-d)^{-} \in S^{+}(G, p, C, \lambda)$.

Proof. (i) Let $0 \leqq \eta \leqq 1$ be a cut-off function and let $v=(u-d)^{+}$. For $k \geqq 0$, let $w=(v-k)^{+}$and $\varphi=-w \eta^{p}$. Then
and

$$
u+\varphi \leqq u \leqq \psi_{2}
$$

$$
u+\varphi \geqq u-(u-d)^{+}=\min (u, d) \geqq \varphi_{1}
$$

in $G$ and $\varphi$ is thus admissible in (1.2). Since

$$
\nabla \varphi=-\nabla w \eta^{p}-p w \eta^{p-1} \nabla \eta
$$

it follows that

$$
\begin{align*}
& \int_{\{u>d+k\}}\left(|\nabla u|^{p}-\mu|u|^{p}-v\right) \eta^{p} d x  \tag{2.2}\\
& \leqq p \int_{\{u>d+k\}}\left(\mu|\nabla u|^{p-1}+\mu|u|^{p-1}+v\right) w \eta^{p-1}|\nabla \eta| d x \\
&+ \int_{\{u>d+k\}}\left(\mu|\nabla u|^{p-1}+\mu|u|^{p-1}+v\right) w \eta^{p} d x
\end{align*}
$$

Applying Young's inequality, we obtain

$$
\begin{gather*}
\int_{\{u>d+k\}}|\nabla u|^{p} \eta^{p} d x \leqq C \int_{\{u>d+k\}}\left(\mu|u|^{p}+v\right) \eta^{p} d x  \tag{2.3}\\
+C \int_{\{u>d+k\}} w^{p}\left(\eta^{p}+|\nabla \eta|^{p}\right) d x+p \int_{(u>d+k\}}\left(\mu|u|^{p-1}+v\right) w\left(\eta^{p-1}|\nabla \eta|+\eta^{p}\right) d x
\end{gather*}
$$

Since $u>d+k$ implies $u=w+d+k$ we have

$$
\left(\mu|u|^{p}+v\right) \eta^{p} \leqq\left(\left(\mu w+\mu d_{0}+\mu k\right)^{p}+v\right) \eta^{p}
$$

and thus the first integral on the right side of (2.3) is bounded by

$$
\begin{equation*}
C_{1} \int_{\{w>0\}}\left(w^{p}+k^{p}+\lambda_{1}\right) \eta^{p} d x \quad \text { where } \quad \lambda_{1}=\lambda_{1}\left(d_{0}, v\right) \tag{2.4}
\end{equation*}
$$

Also, the inequality

$$
\begin{gathered}
\left(\mu|u|^{p-1}+v\right) w\left(\eta^{p}+\eta^{p-1}|\nabla \eta|\right) \\
=\mu|u|^{p-1} w \eta^{p-1}|\nabla \eta|+v w \eta^{p-1}|\nabla \eta|+\mu|u|^{p-1} w \eta^{p}+v w \eta^{p} \\
\leqq C\left(|u|^{p}+v^{p /(p-1)}+w^{p}\right) \eta^{p}+C\left(|u|^{p} \eta^{p}+w^{p}|\nabla \eta|^{p}\right)+C\left(\eta^{p}+w^{p}|\nabla \eta|^{p}\right)
\end{gathered}
$$

together with the fact

$$
\left(\mu|u|^{p}+v^{p /(p-1)}\right) \eta^{p} \leqq C_{2}\left(w^{p}+k^{p}+\lambda_{2}\right) \eta^{p} \quad \text { where } \quad \lambda_{2}=\lambda_{2}\left(d_{0}, v\right),
$$

implies the estimate

$$
\begin{equation*}
C_{2} \int_{\{w>0\}} w^{p}\left(\eta^{p}+|\nabla \eta|^{p}\right) d x+C_{2}\left(k^{p}+\lambda\right) \int_{\{w>0\}} \eta^{p} d x \tag{2.5}
\end{equation*}
$$

for the last integral of (2.3). The estimates (2.4) and (2.5) yield

$$
\int|\nabla w|^{p} \eta^{p} d x \leqq C \int w^{p}\left(\eta^{p}+|\nabla \eta|^{p}\right) d x+C\left(k^{p}+\lambda\right) \int_{\{w>0\}} \eta^{p} d x
$$

with $\lambda=\lambda\left(d_{0}, v\right)$. This proves (i).
The proof of (ii) is similar. Write $v=(u-d)^{-}$and $w=(v-k)^{+}$. Then

$$
\varphi=w \eta^{p}
$$

is admissible for the (1.2) Proceeding as in (i) we obtain the desired result.
Lemma 2.2. With the same hypotheses as Lemma 2.1 and with $v=(u-d)^{2}$, there exists a constant $C=C\left(\mu, v, p, n,\|u\|_{\infty ; G}\right)$ such that

$$
\int_{B\left(x_{0}, r\right)}|\nabla v|^{p} \eta^{p} d x \leqq C\left(\frac{1}{r} \int_{B\left(x_{0}, 2 r\right)}|\nabla v|^{p-1} \eta^{p-1} d x+r^{n-1}\right)
$$

for every $B\left(x_{0}, 2 r\right) \subset G$ and every nonnegative cut-off function whose support is contained in $B\left(x_{0}, 2 r\right)$.

Proof. Consider the proof of the case $v=(u-d)^{+}$. Let $\eta$ be a cut-off function such that $\eta=1$ on $B(x, r)$ and whose support is contained in $B(x, 2 r)$. The desired result follows from a standard application of Young's inequality in (2.2). The proof in case $v=(u-d)^{-}$is similar.

Lemma 2.3. Let $u$ be a solution of (1.2) and let $d_{0} \geqq 0$. If $d \leqq d_{0}$, then there exist constants $C>0$ and $\lambda \geqq 0$ such that if $M=\sup _{G}(u-d)^{2}$, then
(i) $d \geqq \psi_{1}$ in $G$ implies $M-(u-d)^{+} \in S^{-}(G, p, C, \lambda)$
and
(ii) $d \leqq \psi_{2}$ in $G$ implies $M-(u-d)^{-} \in S^{-}(G, p, C, \lambda)$.

Proof. It will be sufficient to establish (i), the proof of (ii) being similar. Let $v=M-(u-d)^{+}$and let $\eta$ be a cut-off function. Fix $k \geqq 0$ with $k \leqq M$. With
and

$$
w=(v-k)^{-}
$$

$$
\varphi=-w \eta^{p}
$$

we have

$$
\begin{aligned}
\psi_{1} \leqq \min (u, d) & =u-(u-d)^{+} \leqq u-(k-v) \\
& \leqq u+\varphi \leqq u \leqq \psi_{2}
\end{aligned}
$$

in $G$, whence $\varphi$ is an admissible test function for (1.2). Thus, proceeding as in Lemma 2.1, we obtain

$$
\begin{equation*}
\int|\nabla w|^{p} \eta^{p} d x \leqq C \int w^{p}\left(\eta^{p}+|\nabla \eta|^{p}\right) d x+C\left(k^{p}+\lambda\right) \int_{(w>0)} \eta^{p} d x \tag{2.6}
\end{equation*}
$$

where $\lambda=\lambda\left(d_{0}, v\right)$. This proves the necessary estimate for the case $k \leqq M$. If $k \geqq$ $\sup _{G} v=M$, then

$$
(v-k)^{-}=k-v=(k-M)+(v-M)^{-}
$$

and (2.6) holds with $w=(v-M)^{-}$and $k$ replaced by $M$. This implies that (2.6) holds with $w=(v-k)^{-}$since $k-M$ is nonnegative.

The following results established in [MZ1] will be needed.
Theorem 2.4. Let $r \in(0,1], x_{0} \in R^{n}$, and $v \in S^{\varepsilon}\left(B\left(x_{0}, r\right), p, C, \lambda\right)$. Define

$$
w(y)=r^{-1} v\left(r y+x_{0}\right)
$$

for all $y \in B(0,1)$. Then

$$
w \in S^{e}(B(0,1), p, C, \lambda)
$$

This is proved by a straightforward change of variables.
Theorem 2.5. If $v \in S^{t}(B(0,1), p, C, \lambda)$, then

$$
v+2 \lambda^{1 / p} \in S^{z}(B(0,1), p, 2 C, 0)
$$

From these two results, one easily concludes the following.
Corollary 2.6. If $v \in S^{\varepsilon}\left(B\left(x_{0}, r\right), p, C, \lambda\right)$ then

$$
v+2 \lambda^{1 / p} r \in S^{z}\left(B\left(x_{0}, r\right), p, 2 C, 0\right)
$$

Finally we also recall the weak Harnack inequalities established in [MZ1] for functions in the classes $S^{+}$and $S^{-}$. They play a central role in this paper.

Theorem 2.7. Let $p \in(1, n), C>0$ and $\gamma \in(0, p]$. There exists a constant $C^{\prime}=$ $C^{\prime}(n, p, C, \gamma)$ such that for $x_{0} \in R^{n}, \sigma \in(0,1), r \in(0,1), \lambda \geqq 0$, and
we have

$$
v \in S^{+}\left(B\left(x_{0}, r\right), p, C, \lambda\right)
$$

$$
\sup _{B\left(x_{0}, \sigma\right)} v \leqq C^{\prime}(1-\sigma)^{-\xi}\left[\left\{f_{B\left(x_{0}, r\right)} v(x)^{\gamma} d x\right\}^{1 / y}+\lambda^{1 / p_{r}}\right]
$$

where $\xi=n / \gamma$.

Theorem 2.8. Let $p \in(1, n), \sigma \in(0,1)$ and $C>0$. There exist positive constants $C^{\prime}, C^{\prime \prime}$, and $\gamma \in(0,1)$ depending only on $n, p, \sigma$ and $C$ such that for $x_{0} \in R^{n}, r \in(0,1)$, $\lambda \geqq 0$, and

$$
v \in S^{-}\left(B\left(x_{0}, r\right), p, C, \lambda\right)
$$

with $v \geqq 0$, we have

$$
\inf _{B\left(x_{0}, r\right)} v \geqq C^{\prime}\left\{f_{B\left(x_{0}, r\right)} v(x)^{\gamma} d x\right\}^{1 / \gamma}-C^{\prime \prime} \lambda^{1 / p} r
$$

Corollary 2.9. Let $u$ be a bounded solution of (1.2).
(i) If $d \geqq \psi_{1}$ in $G \subset \Omega$, then $(u-d)^{+}$is upper semicontinuous in $G$.
(ii) If $d \leqq \psi_{2}$ in $G \subset \Omega$, then $(u-d)^{-}$is upper semicontinuous in $G$.

Proof. This follows immediately from Lemmas 2.3 and 2.8. Indeed, for the first part of the corollary, we have $\mu(r)-(u-d)^{+} \in S^{-}(G, p, C, \lambda)$ where $\mu(r)=$ $\sup _{B\left(x_{0}, r\right)}(u-d)^{+}, x_{0} \in G$, and $r$ sufficiently small. Hence,

$$
\mu(r)-\mu(r / 2) \geqq C^{\prime}\left\{f_{B\left(x_{0}, r\right)}\left[\mu(r)-(u-d)^{+}\right]^{y} d x\right\}^{1 / v}-C^{\prime \prime} \lambda^{1 / p} r .
$$

Since $u$ is bounded, this implies that

$$
f_{B\left(x_{0} r\right)}\left[\mu(r)-(u-d)^{+}\right] d x \rightarrow 0 \quad \text { as } \quad r \rightarrow 0
$$

Thus, it follows that $(u-d)^{+}$has a Lebesgue point at each point of $G$ and that the value of $(u-d)^{+}$is equal to its upper limit at each point of $G$, which implies upper semicontinuity. A similar argument holds for the second part.

Remark 2.10. Theorems 2.7 and 2.8 are counterparts to the well-known weak Harnack inequalities for sub and super solutions of equations in divergence form [GT, Chapter 8]. While sub and super solutions are elements of the classes $S^{+}$ and $S^{-}$, the greater generality of these classes do not yield results with the same precision as with sub and super solutions. For example, the exponent $\gamma$ in Theorem 2.8 can be taken as any positive number less than $n(p-1) /(n-p)$ in the case of a weak supersolution (see [T]) whereas no such bound is available to functions in the class $S^{-}$. A similar phenomenon was encountered in [DT] where weak Harnack inequalities were established for quasiminima, but where the exponent in the weak Harnack inequality for super quasiminima is not quite as strong as for supersolution of elliptic equations. This slight difference will be a factor in some of our results below.

## 3. Regularity properties of the solution

In this section we will initiate a study of the regularity of the solution to (1.2). For this we will introduce the following notation. We let $|E|$ denote the Lebesgue measure of a set $E \subset R^{n}$. We will also need a more refined method of measuring sets.

Definition 3.1. For $1<p<n$, the $p$-capacity of a compact set $K \subset R^{n}$ is defined as

$$
\gamma_{p}(K)=\inf \left\{\int_{R^{n}}|\nabla v|^{p} d x\right\}
$$

where the infimum is taken over all functions $v \in C_{0}^{\infty}\left(R^{n}\right)$ such that $v \geqq 1$ on $K$. The definition of $\gamma_{p}$ can be extended to all sets by standard methods.

The p-capacity is a capacity in the sense of Brelot-Choquet and enjoys the following properties that will be needed later, cf. [AH]. If $q<p$ then

$$
\begin{gather*}
\gamma_{q}(E) \leqq C\left[\gamma_{p}(E)\right]^{(n-q) /(n-p)} \quad \text { for every } \quad E \subset B(0,1)  \tag{3.1}\\
\gamma_{p}[B(x, r)]=C r^{n-p} \quad \text { for every ball } B(x, r) .
\end{gather*}
$$

Since we are concerned with point-wise regularity, we will consider a fixed point $x_{0} \in \Omega$ throughout the remainder of the paper. Let

$$
\begin{array}{cl}
\Psi_{1}(r)=p_{B\left(x_{0}, r\right)}^{p-\sup } \psi_{1} & \underline{\psi}_{1}(r)=p-\inf _{B\left(x_{0}, r\right)} \psi_{1} \\
\bar{\Psi}_{1}\left(x_{0}\right)=\lim _{r \rightarrow 0} \bar{\psi}_{1}(r) & \underline{\psi}_{1}\left(x_{0}\right)=\lim _{r \rightarrow 0} \psi_{1}(r)  \tag{3.2}\\
\underline{\psi}_{1}\left(E_{1}, r\right)=\inf _{E_{1} \cap B\left(x_{0}, r\right)} \psi_{1} & \Psi_{1}\left(E_{1}, r\right)=\sup _{E_{1} \cap B\left(x_{0}, r\right)} \psi_{1} .
\end{array}
$$

Here $p$-sup and $p$-inf are the essential supremum and infimum in the sense of $p$ capacity. Similar notation will be used for $\psi_{2}$ and the solution $u$.

Theorem 3.2. Let $u$ be a solution of (1.2) and assume $\Psi_{1}\left(x_{0}\right) \leqq \underline{\psi}_{2}\left(x_{0}\right)$ where $x_{0} \in \Omega$. Then $u$ possesses a Lebesgue point at $x_{0}$.

Proof. Let

$$
\begin{gathered}
\mu(r)=\sup _{B\left(x_{0}, r\right)}\left[u-\bar{\psi}_{1}(r)\right], \\
v(x)=\mu(r)-\left[u(x)-\bar{\psi}_{1}(r)\right]^{+} .
\end{gathered}
$$

Lemma 2.3 implies that $v \in S^{-}\left(B\left(x_{0}, r\right), p, C, \lambda\right)$ and therefore we conclude from Theorem 2.8 that

$$
\begin{aligned}
\left\{f_{B\left(x_{0}, r\right)}\right. & {\left.\left[\mu(r)-\left(u(x)-\bar{\psi}_{1}(r)\right)^{+}\right]^{y} d x\right\}^{1 / \gamma}-C^{\prime} \lambda^{1 / p} r \leqq C \inf _{B\left(x_{0}, r / 2\right)} v } \\
& \leqq C\left[\mu(r)-\bar{u}(r / 2)+\bar{\psi}_{1}(r)\right]=C[\bar{u}(r)-\bar{u}(r / 2)]
\end{aligned}
$$

Observe that the right-side tends to 0 as $r \rightarrow 0$. Thus,

$$
\lim _{r \rightarrow 0}\left\{f_{B\left(x_{0}, r\right)}\left[\mu(r)-\left(u(x)-\bar{\psi}_{1}(r)\right)^{+}\right]^{\gamma} d x\right\}^{1 / y}=0
$$

and since $u$ is bounded, we obtain the same conclusion with $\gamma=1$. Now let $\mu\left(x_{0}\right)=$ $\lim _{r \rightarrow 0} \mu(r)$ and assume that $\mu\left(x_{0}\right)>0$. Then we have

$$
\lim _{r \rightarrow 0} f_{B\left(x_{0} r\right)}\left[\mu\left(x_{0}\right)-\left(u(x)-\Psi_{1}\left(x_{0}\right)\right)^{+}\right] d x=0
$$

Now write

$$
\begin{gathered}
f_{B\left(x_{0}, r\right)}\left[\mu\left(x_{0}\right)-\left(u(x)-\Psi_{1}\left(x_{0}\right)\right)^{+}\right] d x \\
=\left|B\left(x_{0}, r\right)\right|^{-1} \int_{B\left(x_{0}, r\right) \cap\left\{u \geq \bar{W}_{1}\left(x_{0}\right)\right\}}\left|\mu\left(x_{0}\right)+\Psi_{1}\left(x_{0}\right)-u(x)\right| d x \\
+\mu\left(x_{0}\right) \frac{\left|B\left(x_{0}, r\right) \cap\left\{u<\Psi_{1}\left(x_{0}\right)\right\}\right|}{\left|B\left(x_{0}, r\right)\right|}
\end{gathered}
$$

Notice that both terms on the right tend to 0 as $r \rightarrow 0$ and since we are assuming $\mu\left(x_{0}\right) \neq 0$, we have

$$
\lim _{r \rightarrow 0} \frac{\left|B\left(x_{0}, r\right) \cap\left\{u<\Psi_{1}\left(x_{0}\right)\right\}\right|}{\left|B\left(x_{0}, r\right)\right|}=0 .
$$

Since $u$ is bounded, it now follows that

$$
\lim _{r \rightarrow 0} f_{B\left(x_{0}, r\right)}\left|\mu\left(x_{0}\right)+\bar{\psi}_{1}\left(x_{0}\right)-u(x)\right| d x=0
$$

if $\mu\left(x_{0}\right) \neq 0$, our desired conclusion. In case $\mu\left(x_{0}\right)=0$, then we have

$$
\text { ess } \limsup _{x \rightarrow x_{0}} u(x)=\bar{\psi}_{1}\left(x_{0}\right) .
$$

In this situation, let

$$
v(x)=\lambda(r)-\left(u(x)-\psi_{2}(r)\right)^{-}
$$

where $\lambda(r)=\sup _{B\left(x_{0}, r\right)}\left(u-\underline{\psi}_{2}(r)\right)^{-}$. Then as above, we find that $u$ possesses a Lebesgue point at $x_{0}$ if $\lim _{r \rightarrow 0} \lambda(r)=\lambda\left(x_{0}\right)>0$. If $\lambda\left(x_{0}\right)=0$, we have

$$
\text { ess } \liminf _{x \rightarrow x_{0}} u(x)=\underline{\psi}_{2}\left(x_{0}\right)
$$

This, along with the case $\mu\left(x_{0}\right)=0$ and the assumption $\bar{\psi}_{1}\left(x_{0}\right) \leqq \psi_{2}\left(x_{0}\right)$ implies that the essential limit of $u$ at $x_{0}$ exists and therefore $u$ possesses a Lebesgue point there.

This result immediately leads to our first conclusion on regularity.
Corollary 3.3. If $u$ is a solution of (1.2) and $\Psi_{1}\left(x_{0}\right) \leqq u\left(x_{0}\right) \leqq \underline{\psi}_{2}\left(x_{0}\right)$, then $u$ is continuous at $x_{0}$.

Proof. By the previous result, note that $u$ has a Lebesgue point at $x_{0}$ and therefore the theorem has meaning. If $d>u\left(x_{0}\right)$, then $d>\psi_{1}$ in a neighborhood of $x_{0}$. Moreover, by Corollary 2.9, v=(u-d)+ is upper semicontinuous at $x_{0}$ and $v\left(x_{0}\right)=0$. It follows that

$$
\limsup _{x \rightarrow x_{0}} u(x) \leqq \limsup _{x \rightarrow x_{0}} v(x)+d=d
$$

Now let $d \rightarrow u\left(x_{0}\right)$ to obtain that $u$ is upper semicontinuous at $x_{0}$. A dual argument shows that $u$ is also lower semicontinuous at $x_{0}$, thus establishing the desired conclusion.

Now we have shown that $u$ has a Lebesgue point under the assumption $\bar{\psi}_{1}\left(x_{0}\right) \leqq$ $\underline{\psi}_{2}\left(x_{0}\right)$, we wish to use this result to prove more, namely, that $u$ is finely continuous at $x_{0}$. For this purpose, we will need the following result.

Theorem 3.4. Let $p>1$ and $C>0$. There exist constants $C^{\prime}=C^{\prime}(n, p, C)$ and $\alpha<p$ such that if $v \in S^{-}(G, p, C, 0)$ with $v(x)>0$ for almost all $x \in G$, then

$$
\int_{G} v(x)^{-\beta}|\nabla v(x)|^{p} \eta(x)^{p} d x \leqq C^{\prime} \int_{G} v(x)^{p-\beta}\left[\eta(x)^{p}+|\nabla \eta(x)|^{p}\right] d x
$$

for every $\beta>\alpha$ and every nonnegative Lipschitz function $\eta$ with compact support in $G$.
Proof. Since $v+t \in S^{-}(G, p, C, 0)$ whenever $t>0$, we may assume that $v$ is bounded away from 0 on $G$. Observe also that (2.1) holds whenever $\eta \in W_{0}^{1, p}(G)$. Let $T>0, \gamma>0$ and set

$$
\begin{aligned}
& v_{T}(x)=\inf \{T, v(x)\} \\
& \eta_{T}(x)=v_{T}(x)^{\gamma} \eta(x)
\end{aligned}
$$

Now replace $\eta$ by $\eta_{T}$ in (2.1). This yields

$$
\begin{aligned}
\int_{G}\left|\nabla(v-k)^{-}\right|^{p} v_{T}^{p \gamma} \eta^{p} d x & \leqq C_{1} \int_{G}\left[(v-k)^{-}\right]^{p}\left[v_{T}^{p \gamma} \eta^{p}+\gamma^{p} v_{T}^{p \gamma-p}\left|\nabla v_{T}\right|^{p} \eta^{p}\right. \\
& \left.+v_{T}^{p \gamma}|\nabla \eta|^{p}\right] d x+C_{1} k^{p} \int_{\left\{(v-k)^{->0}\right.} v_{T}^{p \gamma} \eta^{p} d x
\end{aligned}
$$

where $C_{1}=C_{1}(C, p)$. Let $\beta \geqq p$ and multiply both sides of the previous inequality by $k^{-p y-\beta-1}$ and integrate with respect to $k$ over $(0, \infty)$ to obtain

$$
\begin{gathered}
\frac{1}{p \gamma+\beta} \int_{G} v^{-p \gamma-\beta} v_{T}^{p \gamma}|\nabla v|^{p} \eta^{p} d x \\
\leqq \frac{2 C_{1}}{-p+p \gamma+\beta} \int_{G} v^{-p-p y-\beta}\left[v_{T}^{p \gamma} \eta^{p}+\gamma^{p} v_{T}^{-p+p \gamma}\left|\nabla v_{T}\right|^{p} \eta^{p}+v_{T}^{p \gamma}|\nabla \eta|^{p}\right] d x .
\end{gathered}
$$

By letting $T \rightarrow \infty$ and observing that

$$
\frac{p \gamma+\beta}{-p+p \gamma+\beta}=1+\frac{p}{-p+p \gamma+\beta} \leqq 1+\frac{1}{\gamma}
$$

we obtain

$$
\begin{aligned}
\int_{G} v^{-\beta}|\nabla v|^{p} \eta^{p} d x & \leqq 2 C_{1}\left(1+\frac{1}{\gamma}\right) \int_{G} v^{p-\beta}\left[\eta^{p}+|\nabla \eta|^{p}\right] d x \\
& +2 C_{1}\left(1+\frac{1}{\gamma}\right) \gamma^{p} \int_{G} v^{-\beta}|\nabla v|^{p} \eta^{p} d x
\end{aligned}
$$

A suitable choice of $\gamma$ now yields the required inequality under our assumption of $\beta \geqq p$. However, it is possible to make a suitable choice of $\gamma$ under a slightly weaker constraint on $\beta$. For a certain range of $\beta$ less than $p$, we need $\gamma>0$ such that

$$
\begin{equation*}
0<\frac{(\gamma)^{p}(p \gamma+\beta)}{p \gamma+\beta-p}<\varepsilon \tag{3.3}
\end{equation*}
$$

for arbitrary $\varepsilon>0$. Clearly, such a $\gamma$ exists, say $\gamma_{1}$, in case $\beta \geqq p$. Under the assumption that $\beta<p$, there exists $\gamma_{2}$ such that

$$
\begin{equation*}
\left(\frac{p \gamma_{2}+\beta}{p \gamma_{2}}\right)\left(\gamma_{2}\right)^{p} \leqq\left(\frac{p \gamma_{2}+p}{p \gamma_{2}}\right)\left(\gamma_{2}\right)^{p} \leqq \frac{\varepsilon}{2} . \tag{3.4}
\end{equation*}
$$

Now let $\gamma_{0}=\min \left(\gamma_{1}, \gamma_{2}\right)$ and place an additional constraint on $\beta$ by requiring

$$
\frac{\beta-p}{p \gamma_{0}} \geqq-\frac{1}{2} .
$$

Then, $p \gamma_{0}+\beta-p>0$ and because of (3.4), it is easy to see that (3.3) is satisfied.
We now proceed to prove continuity of $u$ at points $x_{0}$ where the obstacles $\psi_{1}$ and $\psi_{2}$ possess some regularity. The amount of regularity required is given by the following definition.

Definition 3.5. A point $x_{0}$ is said to be an upper ( $p, \eta$ )-Wiener point for $\psi_{1}$ if there exists a set $E_{1}$ such that

$$
\begin{equation*}
\int_{0}^{1}\left(\frac{\gamma_{p}\left[E_{1} \cap B\left(x_{0}, r\right)\right]}{r^{n}-p}\right)^{1 / n} \frac{d r}{r}=\infty \tag{3.5}
\end{equation*}
$$

and

$$
\liminf _{\substack{x \rightarrow x_{0} \\ x \in E_{1}}} \psi_{1}(x) \geqq \bar{\psi}_{1}\left(x_{0}\right)
$$

Similarly, $x_{0}$ is said to be a lower $(p, \eta)$-Wiener point for $\psi_{2}$ if there exists a set $E_{2}$ satisfying the same condition as (3.5) and such that

$$
\limsup _{\substack{x \rightarrow x_{0} \\ x \in E_{2}}} \psi_{2}(x) \leqq \psi_{2}\left(x_{0}\right) .
$$

A set $E_{1}$ satisfying (3.5) is said to be not thin at $x_{0}$.

Remark 3.6. In the usual definition of the Wiener integral for nonlinear problems, the exponent $1 / \eta$ is replaced by $1 /(p-1)$, cf. [MA], [GZ]. As suggested in Remark 2.10, we are forced to work with a slightly weaker Wiener condition because of the greater generality of the classes $S^{+}$and $S^{-}$. This will manifest itself in our main result below where a crucial role is played by the weak Harnack inequality, Theorem 2.8.

Theorem 3.7. Let $x_{0} \in \Omega$ be both an upper ( $p, \gamma$ )-Wiener point for $\psi_{1}$ and a lower $(p, \gamma)$-Wiener point for $\psi_{2}$ where $1 / \gamma$ is the exponent that appears in Theorem 2.8. If $\Psi_{1}\left(x_{0}\right) \leqq \psi_{2}\left(x_{0}\right)$, then a weak solution $u$ of $(1.2)$ is continuous at $x_{0}$.

Proof. Recall Theorem 3.2 which states that $u$ of (1.2) is defined at $x_{0}$. If

$$
\begin{equation*}
\bar{\psi}_{1}\left(x_{0}\right) \leqq u\left(x_{0}\right) \leqq \psi_{2}\left(x_{0}\right) \tag{3.6}
\end{equation*}
$$

then $u$ is continuous at $x_{0}$ by Corollary 3.3. Thus, it will be sufficient to prove (3.6). In fact, it is sufficient to prove

$$
\begin{equation*}
\bar{\psi}_{1}\left(x_{0}\right) \leqq u\left(x_{0}\right) . \tag{3.7}
\end{equation*}
$$

Since the other inequality in (3.6) will follow by a dual argument. Let

$$
v=(u-d)^{-} \quad \text { where } \quad d<\bar{\psi}_{1}\left(x_{0}\right)=u\left(x_{0}\right) .
$$

Since $d<\Psi_{2}(2 r)$ for all small $r>0$, note that

Also, let

$$
v \in S^{+}\left(B\left(x_{0}, 2 r\right), p, C, \lambda\right)
$$

$$
\begin{aligned}
m(r) & =\sup _{B\left(x_{0}, r\right)}(u-d)^{-} \\
w & =m(2 r)-v \in S^{-}\left(B\left(x_{0}, 2 r\right), p, C, \lambda\right) \\
A & =\{u>d\} .
\end{aligned}
$$

If $E_{1}$ is a set which is not thin at $x_{0}$ and has the property that

$$
\liminf _{\substack{x \rightarrow x_{0} \\ x \in E_{1}}} \psi_{1}(x) \geqq \bar{\psi}_{1}\left(x_{0}\right),
$$

then since $u \geqq \psi_{1}$ q.e. we have

$$
\begin{equation*}
E_{1} \cap B\left(x_{0}, r\right) \subset A \cap B\left(x_{0}, r\right) \tag{3.8}
\end{equation*}
$$

for all small $r>0$. Let $\eta$ be a cut-off function supported in $B\left(x_{0}, 2 r\right)$ such that $\eta=1$ on $B\left(x_{0}, r\right)$. Then,

$$
\begin{aligned}
m(2 r)^{p} \gamma_{p}\left[A \cap B\left(x_{0}, r\right)\right] & \leqq \int_{B\left(x_{0}, 2 r\right)}|\nabla(\eta w)|^{p} d x \\
& \leqq C\left(\int_{B\left(x_{0}, 2 r\right)}|\nabla w|^{p} \eta^{p}+w^{p}|\nabla \eta|^{p} d x\right) .
\end{aligned}
$$

Now let $w^{\prime}=w+2 \lambda^{1 / p} r$ and refer to Corollary 2.6 to find that

$$
w^{\prime} \in S^{-}\left(B\left(x_{0}, 2 r\right), p, C, 0\right)
$$

Thus, we obtain

$$
\begin{aligned}
& \left(\int_{B\left(x_{0}, 2 r\right)}|\nabla w|^{p} \eta^{p}+w^{p}|\nabla \eta|^{p} d x\right)=\left(\int_{B\left(x_{0}, 2 r\right)}\left|\nabla w^{\prime}\right|^{p} \eta^{p}+w^{p}|\nabla \eta|^{p} d x\right) \\
& \quad \leqq C\left(\frac{1}{r} \int_{B\left(x_{0}, 2 r\right)}\left|\nabla w^{\prime}\right|^{p-1} \eta^{p-1} d x+r^{n-p} f_{B\left(x_{0}, 2 r\right)} w^{p} d x+r^{n-1}\right)
\end{aligned}
$$

(by Lemma 2.2)

$$
\leqq C\left(\frac { 1 } { r } \int _ { B ( x _ { 0 } , 2 r ) } \left(\eta\left(w^{\prime}\right)^{\left.\left.(-\beta) / p\left|\nabla w^{\prime}\right|\right)^{p-1}\left(\left(w^{\prime}\right)^{(\beta(p-1)) / p}\right) d x+I(r)\right), ~(x)}\right.\right.
$$

where $\beta<p$ is obtained from Theorem 3.4 and

$$
\begin{gathered}
I(r)=r^{n-p} f_{B\left(x_{0}, 2 r\right)} w^{p} d x+r^{n-1} \\
\leqq C\left(\frac{1}{r}\left(\int_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{-\beta}\left|\nabla w^{\prime}\right|^{p} \eta^{p} d x\right)^{(p-1) / p}\left(\int_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{\beta(p-1)} d x\right)^{1 / p}+I(r)\right) \\
\leqq C\left(\frac{1}{r}\left(\int_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{p-\beta}\left[\eta^{p}+|\nabla \eta|^{p}\right] d x\right)^{(p-1) / p}\left(\int_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{\beta(p-1)} d x\right)^{1 / p}+I(r)\right)
\end{gathered}
$$

(by Theorem 3.4)

$$
\leqq C\left(\frac{1}{r} r^{n(p-1) / p}\left(f_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{p-\beta} d x\right)^{(p-1) / p} \cdot r^{n / p}\left(f_{B\left(x_{j}, 2 r\right)}\left(w^{\prime}\right)^{\beta(p-1)} d x\right)^{1 / p}+I(r)\right)
$$

$$
\leqq C\left(r^{n-p}\left(f_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{p-\beta} d x\right)^{(p-1) / p}\left(f_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{\beta(p-1)} d x\right)^{1 / p}+I(r)\right)
$$

$$
\leqq C r^{n-p}\left(\left(f_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{p-\beta} d x\right)^{(p-1) / p}\left(f_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{y} d x\right)^{1 / p}+f_{B\left(x_{0}, 2 r\right)} w^{p} d x+r^{p-1}\right)
$$

(with $\beta(p-1)>\gamma$ and from the fact that $w^{\prime}$ is bounded)
$\leqq C r^{n-p}\left(f_{B\left(x_{0} \cdot 2 r\right)}\left(w^{\prime}\right)^{\gamma} d x+f_{B\left(x_{0}, 2 r\right)} w^{p} d x+r^{p-1}\right)$

$$
\leqq C r^{n-p}\left(\left(\inf _{B\left(x_{0}, r\right)} w^{\prime}\right)^{\gamma}+\left(\inf _{B\left(x_{0}, r\right)} w\right)^{\gamma}+r^{p-1}\right)
$$

(using the fact that $w$ is bounded and from Theorem 2.8)

$$
\leqq C r^{n-p}\left[\left(m(2 r)-m(r)+2 \lambda^{1 / p} r\right)^{y}+r^{p-1}\right]
$$

$$
\text { (since } \inf _{B\left(x_{0}, r\right)} w \text { is }<1 \text { for small } r \text { ). }
$$

That is,
(3.9) $m(2 r)^{p}\left(\frac{\gamma_{p}\left[A \cap B\left(x_{0}, r\right)\right]}{r^{n-p}}\right) \leqq C\left[m(2 r)-m(r)+2 \lambda^{1 / p} r\right]^{\gamma}$ for $0<r<1$,
which implies

$$
\int_{0}^{1} m(2 r)^{p / v}\left(\frac{\gamma_{p}\left[A \cap B\left(x_{0}, r\right)\right]}{r^{n-p}}\right)^{1 / v} \frac{d r}{r}<\infty .
$$

Referring to (3.8), we obtain

$$
\infty=\int_{0}^{1}\left(\frac{\gamma_{p}\left[E_{1} \cap B\left(x_{0}, r\right)\right]}{r^{n-p}}\right)^{1 / \gamma} \frac{d r}{r} \leqq \int_{0}^{1}\left(\frac{\gamma_{p}\left[A \cap B\left(x_{0}, r\right)\right]}{r^{n-p}}\right)^{1 / y} \frac{d r}{r}
$$

it follows that $m(r) \rightarrow 0$ as $r \rightarrow 0$. This implies that $\lim _{r \rightarrow 0} \underline{u}(r) \geqq d$; that is, $\underline{u}\left(x_{0}\right) \geqq d$.
Recall that this is proved for any $d<\bar{\psi}_{1}\left(x_{0}\right)$. Thus, we have that

$$
u\left(x_{0}\right)=\lim _{r \rightarrow 0} f_{B\left(x_{0} r\right)} u(x) d x \geqq \underline{u}\left(x_{0}\right) \geqq \bar{\psi}_{1}\left(x_{0}\right) .
$$

Remark 3.8. It is easily seen that the hypothesis $\bar{\psi}_{1}\left(x_{0}\right) \leqq \psi_{2}\left(x_{0}\right)$ is necessary for a solution to be continuous at $x_{0}$. Indeed, since $u \geqq \psi_{1}$ q.e. and $u \leqq \psi_{2}$ q.e., it follows

$$
\begin{aligned}
\Psi_{1}\left(x_{0}\right) & =\inf _{r>0} p-\sup \psi_{B\left(x_{0}, r\right)} \psi_{1} \leqq \inf _{r>0} p-\sup u=u\left(x_{0}\right) \\
& =\sup _{r>0} p-\inf u \leqq \sup _{B\left(x_{0}, r\right)} u-\text {-inf } \psi_{r>0} \psi_{B\left(x_{0}, r\right)}=\psi_{2}\left(x_{0}\right) .
\end{aligned}
$$

## 4. Modulus of continuity

Inequality (3.9) is fundamental in establishing the continuity of the solution. However, it is not strong enough to establish a meaningful bound on the modulus of continuity. For this purpose, one might employ a substitute for $\gamma_{p}$. The capacity needed is $\gamma_{q}$ where

$$
q=n-\frac{p(n-p)}{\gamma}
$$

where $\gamma$ is the number that appears in Theorem 2.8. Referring to (3.1) it is easily seen that (3.9) becomes

$$
\begin{equation*}
m(2 r)^{\gamma}\left(\frac{\gamma_{q}\left[A \cap B\left(x_{0}, r\right)\right]}{r^{n-q}}\right)^{\gamma / p} \leqq C\left[m(2 r)-m(r)+2 \lambda^{1 / p} r\right]^{\gamma} \text { for } 0<r<1 \tag{4.1}
\end{equation*}
$$

That is,

$$
\begin{equation*}
m(2 r)\left(\frac{\gamma_{q}\left[A \cap B\left(x_{0}, r\right)\right]}{r^{n-q}}\right)^{1 / p} \leqq C\left[m(2 r)-m(r)+2 \lambda^{1 / p} r\right] \text { for } 0<r<1 \tag{4.2}
\end{equation*}
$$

Now that we have (4.2), it follows immediately from [GZ, Theorem 2.7] that $m(r)$
satisfies the following growth condition:

$$
m(r) \leqq C_{1} \exp \left(-C_{2} \int_{2 s}^{r} A(t) \frac{d t}{t}\right) \quad \text { for every } \quad s \leqq r / 2
$$

where

$$
A(t)=\left(\frac{\gamma_{q}\left[E_{1} \cap B\left(x_{0}, t\right)\right]}{t^{n-q}}\right)^{1 / p}
$$

From this it would be possible to obtain a modulus of continuity. The difficulty with this argument is that in the definition of the number $q$, we have no estimate for the number $\gamma$ and therefore $q$ might be negative.

In order to obtain a modulus of continuity under the general structure (1.1), we require the following assumptions on the sets $E_{1}$ and $E_{2}$ that appear in Definition 3.5:

$$
\begin{equation*}
\liminf _{r \rightarrow 0} \frac{\left|E_{1} \cap B\left(x_{0}, r\right)\right|}{\left|B\left(x_{0}, r\right)\right|}>0 \tag{4.3}
\end{equation*}
$$

and

$$
\begin{equation*}
\liminf _{r \rightarrow 0} \frac{\left|E_{2} \cap B\left(x_{0}, r\right)\right|}{\left|B\left(x_{0}, r\right)\right|}>0 \tag{4.4}
\end{equation*}
$$

Theorem 4.1. If $u$ is a solution of (1.2), $E_{1}, E_{2}$ respectively satisfy (4.3) and (4.4), and $\Psi_{1}\left(x_{0}\right) \leqq \Psi_{2}\left(x_{0}\right)$, then $u$ is continuous at $x_{0}$ and its modulus of continuity is estimated by

$$
\omega(r) \leqq C\left[r^{\alpha}+M_{1}(r)+M_{2}(r)\right]
$$

where

$$
\begin{aligned}
& M_{1}(r)=\sup \left(\bar{\psi}_{1}\left(x_{0}\right)-\underline{\psi}_{1}\left(E_{1}, r\right), \underline{\psi}_{2}\left(x_{0}\right)-\underline{\psi}_{2}(2 r)\right) \\
& M_{2}(r)=\sup \left(\Psi_{2}\left(E_{2}, r\right)-\underline{\psi}_{2}\left(x_{0}\right), \bar{\psi}_{1}(2 r)-\bar{\psi}_{1}\left(x_{0}\right)\right)
\end{aligned}
$$

Proof. We know from Theorem 3.2 that $u$ is defined at $x_{0}$ and that (4.3) implies $\Psi_{1}\left(x_{0}\right) \leqq u\left(x_{0}\right)$. Likewise, (4.4) implies that $u\left(x_{0}\right) \leqq \psi_{2}\left(x_{0}\right)$. Hence, by Theorem 3.7 it follows that $u$ is continuous at $x_{0}$. If

$$
\bar{\psi}_{1}\left(x_{0}\right)<u\left(x_{0}\right)<\underline{\psi}_{2}\left(x_{0}\right),
$$

then $u$ is a solution of the associated equation in a neighborhood of $x_{0}$ and is therefore Hölder continuous there. This establishes part of our conclusion.

Next we assume

$$
\begin{equation*}
\Psi_{1}\left(x_{0}\right)=u\left(x_{0}\right)<\psi_{2}\left(x_{0}\right) \tag{4.5}
\end{equation*}
$$

and consider the function $w$ defined in the proof of Theorem 3.7 with $d=d(r)=$
$\Psi_{1}\left(E_{1}, r\right)-r$ where $\psi_{1}\left(E_{1}, r\right)=\inf _{B\left(x_{0}, r\right) \cap E_{1}} \psi_{1}$. Because of assumption (4.5), note that $d(r)<\Psi_{2}(r)$ and that $E_{1} \cap B\left(x_{0}, r\right) \subset A \cap B\left(x_{0}, r\right)$ for all small $r>0$. Then, the functions $v$ and $w$ introduced in Theorem 3.7 belong to the classes $S^{+}$and $S^{-}$ respectively and by Theorem 2.8

$$
\begin{aligned}
& m(r)\left(\frac{\left|E_{1} \cap B\left(x_{0}, r\right)\right|}{\left|B\left(x_{0}, r\right)\right|}\right)^{1 / \gamma} \leqq\left(f_{B\left(x_{0}, r\right)}(w(x))^{\gamma} d x\right)^{1 / \gamma} \\
\leqq & C_{B\left(x_{0}, r / 2\right)} w+C^{\prime} \lambda^{1 / p} r \leqq C\left[m(r)-m(r / 2)+C^{\prime} \lambda^{1 / p} r\right] .
\end{aligned}
$$

Because of (4.3), this implies

$$
m(r) \leqq C^{\prime \prime}\left[m(r)-m(r / 2)+C^{\prime} \lambda^{1 / p} r\right]
$$

for all small $r>0$. This implies

$$
m(r) \leqq C r^{a}
$$

for some $0<\alpha<1$ and for all small $r>0$, of. [GT, Lemma 8.23]. Now

$$
u\left(x_{0}\right)-\underline{u}(r)=m(r)+\Psi_{1}\left(x_{0}\right)-\underline{\psi}_{1}\left(E_{1}, r\right)+r
$$

so that

$$
\begin{equation*}
\underline{\omega}(r) \equiv u\left(x_{0}\right)-\underline{u}(r) \leqq \psi_{1}\left(x_{0}\right)-\underline{\psi}_{1}\left(E_{1}, r\right)+C r^{\alpha} . \tag{4.6}
\end{equation*}
$$

This provides an estimate for the lower oscillation of $u$ under assumption (4.5). We now proceed to obtain an estimate for the lower oscillation in case

$$
\bar{\psi}_{1}\left(x_{0}\right)=u\left(x_{0}\right)=\psi_{2}\left(x_{0}\right) .
$$

The proof in this situation proceeds exactly as the one above except that now we take

$$
d(r)=\inf \left\{\underline{\psi}_{1}\left(E_{1}, r\right), \underline{\psi}_{2}(r)\right\}-r<\underline{\psi}_{2}(r)
$$

which will ensure that the functions $v$ and $w$ are in the appropriate classes and that $E_{1} \cap B\left(x_{0}, r\right) \subset A \cap B\left(x_{0}, r\right)$. We then obtain

$$
u\left(x_{0}\right)-\underline{u}(r) \leqq \Psi_{1}\left(x_{0}\right)+m(r)-d(r) \leqq C\left[r^{\alpha}+M_{1}(r)\right] .
$$

This shows that under assumption (4.3)

$$
\underline{\omega}(r) \leqq C\left[r^{\alpha}+M_{1}(r)\right]
$$

for all small $r>0$. A dual argument establishes

$$
\bar{\omega}(r) \equiv \bar{u}(r)-u\left(x_{0}\right) \leqq C\left[r^{\alpha}+M_{2}(r)\right]
$$

under the assumption (4.4).

Corollary 4.2. If $u$ is a solution of (1.2) and both $\psi_{1}$ and $\psi_{2}$ are locally Hölder continuous with $\psi_{1} \leqq \psi_{2}$ on $\Omega$, then $u$ is locally Hölder continuous on $\Omega$.

In order to obtain stronger results, we modify the structure (1.1) of $A$ and $B$ as follows so that in addition, we require

$$
A(x, u, h) \cdot h \geqq 0
$$

$$
\begin{align*}
& A(x, u, 0)=0  \tag{4.7}\\
& B(x, u, 0)=0
\end{align*}
$$

whenever $u \in R^{1}$ and $h \in R^{n}$.
Although this structure is more restrictive, it nevertheless includes a wide class of interesting equations. For example, the $p$-Laplacean

$$
-\operatorname{div}\left(|\nabla u|^{p-2} \nabla u\right)=0,
$$

and the uniformly elliptic linear equations

$$
L u=-\sum_{i, j=1}^{n}\left(a^{i j}(x) u_{x_{j}}\right)_{x_{i}}
$$

considered in [DMV] where $a^{i j} \in L^{\infty}$. Also, the equations

$$
-\operatorname{div} A(x, \nabla u)=0
$$

with $A(x, h) \cdot h \approx|h|^{p}$, associated with a nonlinear potential theory [HKM] satisfy these assumptions.

The key lemma which results from these extra assumptions is the following.
Lemma 4.3. Suppose that $A$ and $B$ satisfy (1.1) and (4.7). Then
(i) If $d \geqq \psi_{1}$ in $G$, then $v=\max (u, d)$ is a subsolution of (1.2) in $G$; that is

$$
\int A(x, v, \nabla v) \cdot \nabla \varphi+B(x, v, \nabla v) \varphi d x \leqq 0
$$

for every nonnegative test function $\varphi$.
(ii) If $d \leqq \psi_{2}$ in $G$, then $w=\min (u, d)$ is a supersolution of (1.2) in $G$.

Proof. Let $\varphi \in C_{0}^{\infty}(G), \varphi \geqq 0$. Fix $\varepsilon>0$ and write

$$
\eta=-\varepsilon \min \left(\varphi, \frac{(u-d)^{+}}{\varepsilon}\right) .
$$

Then $u+\eta \leqq u \leqq \psi_{2}$ and

$$
u+\eta \geqq u-\varepsilon \frac{(u-d)^{+}}{\varepsilon}=\min (u, d) \geqq \psi_{1}
$$

Thus, $\eta$ is admissible for (1.2) and it follows that

$$
\begin{gathered}
0 \geqq \varepsilon \int_{G} A(x, u, \nabla u) \cdot \nabla \min \left(\varphi, \frac{(u-d)^{+}}{\varepsilon}\right) d x \\
+\varepsilon \int_{G} B(x, u, \nabla u) \min \left(\varphi, \frac{(u-d)^{+}}{\varepsilon}\right) d x \\
=\varepsilon \int_{\left\{x \in G:(u-d)^{+} \leqq \varepsilon \varphi\right\}} A(x, u, \nabla u) \cdot \nabla u d x+\varepsilon \int_{\left\{x \in G:(u-d)^{+}>\varepsilon \varphi\right\}} A(x, u, \nabla u) \cdot \nabla \varphi d x \\
+\varepsilon \int_{G} B(x, u, \nabla u) \min \left(\varphi, \frac{(u-d)^{+}}{\varepsilon}\right) d x .
\end{gathered}
$$

Thus, assumption (4.7) implies

$$
\begin{gathered}
\int_{G \cap\left\{(u-d)^{+>\varepsilon \varphi\}}\right.} A(x, v, \nabla v) \cdot \nabla \varphi d x \\
+\int_{G \cap\left\{(u-d)^{+>0\}}\right.} B(x, v, \nabla v) \min \left(\varphi, \frac{(u-d)^{+}}{\varepsilon}\right) d x \leqq 0 .
\end{gathered}
$$

Now with $\varepsilon \rightarrow 0$ we obtain

$$
\int_{G \cap\left\{(u-d)^{+}>0\right\}} A(x, v, \nabla v) \cdot \nabla \varphi d x+\int_{G \cap\left\{(u-d)^{+}>0\right\}} B(x, v, \nabla v) \varphi d x \leqq 0
$$

and, since by assumption (4.7), we have $A(x, v, \nabla v)=0$ and $B(x, v, \nabla v)=0$ in the set

$$
G \cap\left\{(u-d)^{+}=0\right\}=G \cap\{v=0\}
$$

we have

$$
\int_{G} A(x, v, \nabla v) \cdot \nabla \varphi d x+\int_{G} B(x, v, \nabla v) \varphi d x \leqq 0
$$

as desired.
The proof of (ii) is similar with $\eta$ defined this time as

$$
\eta=\varepsilon \min \left(\varphi, \frac{(u-d)^{-}}{\varepsilon}\right)
$$

Theorem 4.4. If $u$ is a solution (1.2) with structure (1.1) and (4.7), $x_{0}$ is a point that satisfies Definition 3.5 with $\eta=p-1$, and $\bar{\psi}_{1}\left(x_{0}\right) \leqq \underline{\psi}_{2}\left(x_{0}\right)$, then $u$ is continuous at $x_{0}$ and its modulus of continuity for all small $r>0$ and all $s<r / 2$ is given by

$$
\omega(r) \leqq C\left[\exp \left(-C^{\prime} \int_{2 s}^{r} A_{1}(t) \frac{d t}{t}\right)+\exp \left(-C^{\prime} \int_{2 s}^{r} A_{2}(t) \frac{d t}{t}\right)+M_{1}(r)+M_{2}(r)\right]
$$

where

$$
A_{i}=\left(\frac{\gamma_{p}\left[E_{i} \cap B\left(x_{0}, t\right)\right]}{t^{n-p}}\right)^{1 /(p-1)} \quad i=1,2
$$

and where $M_{i}(r), i=1,2$ are as in Theorem 4.1.

Proof. We proceed as in the proof of Theorem 4.1 and thus first prove

$$
\begin{equation*}
\bar{\psi}_{1}\left(x_{0}\right) \leqq u\left(x_{0}\right) \leqq \psi_{2}\left(x_{0}\right) \tag{4.8}
\end{equation*}
$$

Let $d=d(r)=\Psi_{1}\left(E_{1}, 2 r\right)-r$ and note that since $d<\bar{\psi}_{1}\left(x_{0}\right)$, we have $d<\Psi_{2}(2 r)$ for all small $r>0$. Consequently, we infer from Theorem 4.3 that $\min (u, d)$ is a supersolution in $B\left(x_{0}, 2 r\right)$ of an equation of type (1.1) and (1.2). Let $f=\min (u, d)$. Recall from [MZ1, Corollary 4.4] that $f$ is finely continuous at all points of $B\left(x_{0}, 2 r\right)$. Thus there is a set $E_{0} \subset E_{1}$ such that $f \geqq \min \left(\psi_{1}, d\right)$ on $E_{0}$ and that

$$
\lim _{\substack{x \rightarrow x_{0} \\ x \in E}} f(x)=f\left(x_{0}\right) .
$$

Then since $f$ is approximately continuous

$$
u\left(x_{0}\right) \geqq f\left(x_{0}\right)=\lim _{\substack{x \rightarrow x_{0} \\ x \in E_{0}}} f(x) \geqq \liminf _{\substack{x \rightarrow x_{0} \\ x \in E_{0}}} \min \left(\psi_{1}(x), d\right)=d
$$

and by letting $r \rightarrow 0$ we have

$$
u\left(x_{0}\right) \geqq \bar{\psi}_{1}\left(x_{0}\right) .
$$

Thus, the first inequality of (4.8) is established and the second is obtained by a dual argument. Now let

$$
\begin{align*}
m(r) & =\sup _{B\left(x_{0}, r\right)}(u-d)^{-}  \tag{4.9}\\
v & =(u-d)^{-} \\
w & =m(2 r)-v \\
w^{\prime} & =w+\lambda r
\end{align*}
$$

where $\lambda$ is a number that dominates the sum of all the coefficients that appear in (1.1).
We now collect some often used estimates that arise when dealing with sub and supersolutions of (1.2). Since $v$ is a subsolution of an equation of type (1.2), one may use $\varphi=\eta^{p} v$ as a test function and employ standard estimates to obtain

$$
\begin{gather*}
C^{-1} \int_{B\left(x_{0}, 2 r\right)} \eta^{p}|\nabla v|^{p} d x \leqq m(2 r) \int_{B\left(x_{0}, 2 r\right)} \eta^{p-1}|\nabla \eta||\nabla v|^{p-1} d x  \tag{4.10}\\
+m(2 r) \int_{B\left(x_{0}, 2 r\right)} \eta^{p}+|\nabla \eta|^{p} d x .
\end{gather*}
$$

Also, since $w$ is a supersolution, so is $w^{\prime}$ and because of the definition of $\lambda$ one obtains the following estimate, cf. [GZ, equation (23)]:

$$
\begin{equation*}
\int_{B\left(x_{0}, 2 r\right)} w^{\prime}(x)^{-\beta}\left|\nabla w^{\prime}(x)\right|^{p} \eta(x)^{p} d x \leqq C^{\prime} \int_{B\left(x_{0}, 2 r\right)} w^{\prime}(x)^{p-\beta}\left[\eta(x)^{p}+|\nabla \eta(x)|^{p}\right] d x \tag{4.11}
\end{equation*}
$$

for every $\beta>1$ and every nonnegative Lipschitz function $\eta$ with compact support in $B\left(x_{0}, 2 r\right)$. Finally, we will need the following weak Harnack inequality for supersolutions:

$$
\begin{equation*}
\inf _{B\left(x_{0}, \sigma r\right)} v \geqq C^{\prime}\left\{f_{B\left(x_{0}, r\right)} v(x)^{y} d x\right\}^{1 / y}-C^{\prime \prime} \lambda^{1 / p} r \tag{4.12}
\end{equation*}
$$

for any $\gamma<n(p-1) /(n-p)$ and any $\sigma \in(0,1)$, cf. [T]. Proceeding as in the proof of Theorem 3.7, we let $A=\{u>d\}$ and since $E_{1} \cap B\left(x_{0}, r\right) \subset A \cap B\left(x_{0}, r\right)$ for small $r>0$ we obtain

$$
\begin{align*}
& \quad m(2 r)^{p} \gamma_{p}\left[E_{1} \cap B\left(x_{0}, r\right)\right] \\
& \leqq m(2 r)^{p} \gamma_{p}\left[A \cap B\left(x_{0}, r\right)\right] \\
& \leqq \int_{B\left(x_{0}, 2 r\right)}|\nabla(\eta w)|^{p} d x \\
& \leqq C\left(\int_{B\left(x_{0}, 2 r\right)}|\nabla w|^{p} \eta^{p}+w^{p}|\nabla \eta|^{p} d x\right) \\
& =\left(\int_{B\left(x_{0}, 2 r\right)}\left|\nabla w^{\prime}\right|^{p} \eta^{p}+w^{p}|\nabla \eta|^{p} d x\right) \\
& \leqq C\left(\frac{m(2 r)}{r} \int_{B\left(x_{0}, 2 r\right)}\left|\nabla w^{\prime}\right|^{p-1} \eta^{p-1} d x+r^{n-p} f_{B\left(x_{0}, 2 r\right)} w^{p} d x+m(2 r) r^{n-p}\right)  \tag{4.10}\\
& \leqq C\left(\frac{m(2 r)}{r} \int_{B\left(x_{0}, 2 r\right)}\left(\eta\left(w^{\prime}\right)^{-(1-\vartheta)}\left|\nabla w^{\prime}\right|\right)^{p-1}\left(\left(w^{\prime}\right)^{(1-\vartheta)(p-1)}\right) d x+I(r)\right)
\end{align*}
$$

where $\vartheta$ is such that $1<(1-\vartheta) p<n /(n-p)$ and

$$
\begin{align*}
& I(r)=r^{n-p} f_{B\left(x_{0}, 2 r\right)} w^{p} d x+m(2 r) r^{n-p} \\
& \leqq C\left(\frac{m(2 r)}{r}\left(\int_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{-(1-\vartheta) p}\left|\nabla w^{\prime}\right|^{p} \eta^{p} d x\right)^{(p-1) / p}\right. \\
& \left.\times\left(\int_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{(1-\vartheta)(p-1) p} d x\right)^{1 / p}+I(r)\right) \\
& \leqq C\left(\frac{m(2 r)}{r}\left(\int_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{\vartheta p}\left[\eta^{p}+|\nabla \eta|^{p}\right] d x\right)^{(p-1) / p}\left(\int_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{(1-\vartheta)(p-1) p} d x\right)^{1 / p}+I(r)\right) \\
& \leqq C\left(\frac{m(2 r)}{r} r^{n(p-1) / p}\left(f_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{q p} d x\right)^{(p-1) / p}\right.  \tag{4.11}\\
& \left.\times r^{n / p}\left(f_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{(1-\vartheta)(p-1) p} d x\right)^{1 / p}+I(r)\right) \\
& \leqq C\left(m(2 r) r^{n-p}\left(f_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{9 p} d x\right)^{(p-1) / p}\left(f_{B\left(x_{0}, 2 r\right)}\left(w^{\prime}\right)^{(1-\vartheta)(p-1) p} d x\right)^{1 / p}+I(r)\right) \\
& \leqq C\left(m(2 r) r^{n-p}\left(\inf _{B\left(x_{0}, r\right)} w^{\prime \vartheta(p-1)}\right)\left(\inf _{B\left(x_{0}, r\right)} w^{\prime(1-\vartheta)(p-1)}\right)+I(r)\right) \quad \text { by }(4.12) \\
& \leqq C m(2 r) r^{n-p} \inf _{B\left(x_{0}, r\right)}\left(w^{\prime}\right)^{p-1}+C m(2 r) r^{n-p}\left(\inf _{B\left(x_{0}, r\right)} w+C^{\prime} r\right)^{p-1} \\
& \leqq C m(2 r)\left(r^{n-p}\right)[m(2 r)-m(r)+\lambda r]^{p-1}+r^{n-p} \quad \text { for } \quad 0<r<1 .
\end{align*}
$$

That is,

$$
\begin{equation*}
m(2 r)^{p-1}\left(\frac{\gamma_{p}\left[E_{1} \cap B\left(x_{0}, r\right)\right]}{r^{n-p}}\right) \leqq C[m(2 r)-m(r)+\lambda r]^{p-1} \text { for } 0<r<1 \tag{4.13}
\end{equation*}
$$

which implies

$$
\int_{0}^{1} m(2 r)\left(\frac{\gamma_{p}\left[E_{1} \cap B\left(x_{0}, r\right)\right]}{r^{n-p}}\right)^{1 /(p-1)} \frac{d r}{r}<\infty .
$$

Now refer to [GZ, Theorem 2.7] to see that this implies the existence of constants $C_{1}$ and $C_{2}$ such that for sufficiently small $r>0$

$$
m(r) \leqq C_{1} \exp \left(-C_{2} \int_{2 s}^{r} A_{1}(t) \frac{d t}{t}\right)
$$

whenever $s \leqq r / 2$ and where

$$
A_{1}(t)=\left(\frac{\gamma_{p}\left[E_{1} \cap B\left(x_{0}, t\right)\right]}{t^{n-p}}\right)^{1 /(p-1)}
$$

As in the proof of Theorem 4.1 we have

$$
\begin{aligned}
\underline{\omega}(r)=u\left(x_{0}\right)-\underline{u}(r) & \leqq m(r)+\underline{\psi}_{1}\left(x_{0}\right)-\underline{\psi}_{2}\left(E_{1}, r\right)+r \\
& \leqq C_{1} \exp \left(-C_{2} \int_{25}^{r} A_{1}(t) \frac{d t}{t}\right)+M_{1}(r)
\end{aligned}
$$

for every $s \leqq r / 2$. This is the desired conclusion under (4.8). In case

$$
\bar{\psi}_{1}\left(x_{0}\right)=u\left(x_{0}\right)=\underline{\psi}_{2}\left(x_{0}\right)
$$

an estimate for the lower oscillation is obtained in the same way except that $d(r)$ is taken as

$$
d(r)=\inf \left\{\underline{\psi}_{1}\left(E_{1}, r\right), \underline{\psi}_{2}(2 r)\right\}-r
$$

An estimate for the upper oscillation is obtained in a similar way.

## 5. Sharpness of results

In this section we show that the Wiener conditions for the obstacles $\psi_{1}$ and $\psi_{2}$ assumed in Theorem 4.4 are necessary for a non-trivial class of equations, including the linear equations of [DMV]. For this we assume that $A: R^{n} \times R^{n} \rightarrow R^{n}$ is a Borel function satisfying the following assumptions for almost every $x$, all $h$, and for
some $0<\gamma \leqq \mu<\infty$ :

$$
\begin{gather*}
|A(x, h)| \leqq \mu|h|^{p-1}  \tag{5.1}\\
A(x, h) \cdot h \geqq \gamma|h|^{p} \\
\left(A\left(x, h_{1}\right)-A\left(x, h_{2}\right)\right) \cdot\left(h_{1}-h_{2}\right)>0 \quad \text { for } h_{1} \neq h_{2} \\
A(x, \lambda h)=|\lambda|^{p-2} \lambda A(x, h) \quad \lambda \in R^{1} .
\end{gather*}
$$

Lemma 5.1. Suppose that $p>n-1$ or that $p=2$ and the equation

$$
-\operatorname{div} A(x, \nabla u)=0
$$

is linear. If $x_{0}$ is not a p-Wiener point of $\psi_{1}$, then there is a solution $v$ of the unilateral obstacle problem in a neighborhood $U$ of $x_{0}$ such that $v$ is not continuous at $x_{0}$. Moreover,

$$
\text { ess } \liminf _{x \rightarrow x_{0}} v(x)<\bar{\psi}_{1}\left(x_{0}\right)
$$

Proof. The cases $p>n-1$ are proved in [HK, Theorem 1.16]. The linear case is done by Mosco [M, Thm. 5.2]. This linear case can be easily treated without using Green's function by using arguments of [HK, Theorem 1.16] and [HKM, Theorem 3.2].

Theorem 5.2. Suppose $A$ satisfies (5.1) and that $p>n-1$ or that $-\operatorname{div} A(x, \nabla u)=0$ is linear ( $p=2$ ). Suppose, further, that there is a (quasicontinuous) Soboler function $u_{0}$ such that $\psi_{1} \leqq u_{0} \leqq \psi_{2}$ in a neighborhood $U$ of $x_{0}$. If $x_{0}$ is neither an upper ( $p, p-1$ )-Wiener point for $\psi_{1}$ nor a lower ( $p, p-1$ )-Wiener point for $\psi_{2}$, then there is a solution $u$ which is discontinuous at $x_{0}$.

Proof. Let us assume that $x_{0}$ is not a regular point for $\psi_{1}$. If $u$ is a continuous solution of the double obstacle problem, then $u\left(x_{0}\right) \geqq \bar{\psi}_{1}\left(x_{0}\right)$, (see Remark 3.8). Thus, it is enough to show that for some solution $u$

$$
\text { ess } \liminf _{x \rightarrow x_{0}} u(x)<\Psi_{1}\left(x_{0}\right)
$$

To this end, let $v$ be the discontinuous solution to the single obstacle problem given by Lemma 5.1 with

$$
\text { ess } \liminf _{x \rightarrow x_{0}} v(x)<\Psi_{1}\left(x_{0}\right)
$$

Let $u$ be a solution to the double obstacle problem with

$$
u-\min \left(v, u_{0}\right) \in W_{0}^{1, p}(U)
$$

To conclude the proof, we show that $u \leqq v$ in $U$. For this, let

$$
\eta=\min (v-u, 0) \in W_{0}^{1, p}(U)
$$

Then, $u+\eta \leqq u \leqq \psi_{2}$ and $u+\eta=\min (v, u) \geqq \psi_{1}$. Thus,

$$
\begin{aligned}
0 & \leqq \int_{U}(A(x, \nabla u)-A(x, \nabla u)) \cdot \nabla \eta d x \\
& =-\int_{\{0<u\}}(A(x, \nabla u)-A(x, \nabla v)) \cdot(\nabla u-\nabla v) d x=0 .
\end{aligned}
$$

Thus, $\eta=0$ by (5.1), thus proving that $u \leqq v$ a.e. as desired.
Remark 5.3. By refining the argument of [HK] we could dispense with the homogeneity assumption in (5.1).

## 6. Monotone operators

In this section we study operators that satisfy (1.1) as well as the following:

$$
\begin{gather*}
\left(A\left(x, \eta, h_{1}\right)-A\left(x, \eta, h_{2}\right)\right) \cdot\left(h_{1}-h_{2}\right)>0 \text { for } h_{1} \neq h_{2}  \tag{6.1}\\
|B(x, \eta, h)| \leqq \mu|\eta|^{p-1}+v .
\end{gather*}
$$

These operators are not covered by the structure imposed by (4.7). The object of this section is to prove the following result.

Theorem 6.1. Suppose $x_{0}$ is a ( $p, p-1$ )-Wiener point for both obstacles $\psi_{1}$ and $\psi_{2}$. Then, if $\Psi_{1}\left(x_{0}\right) \leqq \Psi_{2}\left(x_{0}\right)$ and (6.1) holds, the solution $u$ to the double obstacle problem is continuous at $x_{0}$.

Before proving this result, we will need the following two Lemmas. Note that since $u$ is bounded, the operator defined by

$$
\begin{aligned}
\tilde{A}(x, h) & =A(x, u(x), h) \\
\tilde{B}(x) & =B(x, u(x), \nabla u(x))
\end{aligned}
$$

satisfy the same structural assumptions.
The following is a special case of a result proved in [MZ2].
Lemma 6.2. Let $w \in W^{1, p}(\Omega), w \geqq \psi_{1}$ (or $w \leqq \psi_{2}$ ). Then there exists $v \in W^{1, p}(\Omega)$ such that $v-w \in W_{0}^{1, p}(\Omega), v \geqq \psi_{1}$ (or $v \leqq \psi_{2}$ ) and

$$
\int_{\Omega} \tilde{A}(x, \nabla v) \cdot \nabla \varphi d x+\int_{\Omega} \tilde{B}(x) \varphi d x \geqq 0
$$

whenerer $\varphi \in W_{0}^{1, p}(\Omega), \varphi \geqq \psi_{1}-u$ (or $\varphi \leqq \psi_{2}-u$ ).
The following lemma is crucial in establishing Theorem 6.1.
Lemma 6.3. Suppose that $x_{0}$ is a $(p, p-1)$-Wiener point for $\psi_{1}$ and $\psi_{2}$ and that
$\Psi_{1}\left(x_{0}\right) \leqq \underline{\psi}_{2}\left(x_{0}\right)$. Then for each $d>\Psi_{2}\left(x_{0}\right)$,
and for each $d<\bar{\psi}_{1}\left(x_{0}\right)$,

$$
\lim _{r \rightarrow 0} f_{B\left(x_{0}, r\right)}(u-d)^{+} d x=0
$$

$$
\lim _{r \rightarrow 0} f_{B\left(x_{0}, r\right)}(u-d)-d x=0
$$

Proof. We will prove only the first assertion, the proof of the second being similar.

Let $d>\underline{\psi}_{2}\left(x_{0}\right)$. Then, for some $\varepsilon>0, \psi_{1}<d-\varepsilon$ q.e. in some neighborhood $U$ of $x_{0}$. Furthermore, the set

$$
E=U \cap\left\{\psi_{2}(x) \leqq d\right\}
$$

is not ( $p, p-1$ )-thin at $x_{0}$. Choose a compact set $K \subset E \cup\left\{x_{0}\right\}$ such that $K$ is not ( $p, p-1$ )-thin at $x_{0}$ and let $G=U-K$. Let $w$ be a quasicontinuous function in $W^{1, p}(\Omega)$ such that $w-u \in W_{0}^{1, p}(U)$ and $w=d$ in $K$. Let $v$ be the solution to the single obstacle problem in $G$ with obstacle $\psi_{1}$ and with $v-w \in W_{0}^{1, P}(G)$; the existence of $v$ being assured by Lemma 6.2. We claim that $v \geqq u$ in $G$. Indeed, let

$$
\eta=\min (v-u, 0)
$$

Then $\eta \in W_{0}^{1, p}(G)$ is nonpositive and admissible for (1.2) since $\psi_{1} \leqq u+\eta \leqq \psi_{2}$ in $G$. Thus,

$$
\begin{aligned}
0 & \leqq \int_{G}(A(x, u(x), \nabla u(x))-\tilde{A}(x, \nabla v)) \cdot \nabla \eta d x \\
& +\int_{G}(B(x, u(x), \nabla u(x))-\tilde{B}(x) \eta d x \\
& =-\int_{\{u<v\}}(A(x, u(x), \nabla u(x))-A(x, u(x), \nabla v(x))) \cdot(\nabla u-\nabla v) d x \\
& \leqq 0
\end{aligned}
$$

Thus, $\eta=0$ by (6.1) and hence $v \geqq u$ which establishes our claim.
Now $v$ is the solution to a single obstacle problem and is therefore a supersolution in $G$ of the associated equation. Since $v-w \in W_{0}^{1, p}(G)$ and $w=d$ in $K$, it follows from [GZ, Theorem 2.2] that

$$
\liminf _{x \rightarrow x_{0}} v(x) \geqq d
$$

whence $v \geqq \psi_{1}+\varepsilon$ q.e. in $B\left(x_{0}, r\right) \cap G$ for some $r>0$. This implies that $v$ in fact is a solution of the equation

$$
-\operatorname{div} \tilde{A}(x, \nabla v)+\tilde{B}(x)=0
$$

in $B\left(x_{0}, r\right) \cap G$. Now since $v$ is a solution, we may employ [GZ, Theorem 2.2] again to obtain

$$
\lim _{x \rightarrow x_{0}} v(x)=d
$$

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Thus,

$$
\limsup _{r \rightarrow 0} f_{B\left(x_{0}, r\right)}(u-d)^{+} d x \leqq \lim _{r \rightarrow 0} f_{B\left(x_{0}, r\right)}(v-d)^{+} d x=0
$$

which completes the proof.
We now return to the proof of Theorem 6.1. Under the hypotheses of this theorem, we know from Theorem 3.2 that $u$ has a Lebesgue point at $x_{0}$. The result above implies that $\bar{\psi}_{1}\left(x_{0}\right) \leqq u\left(x_{0}\right) \leqq \psi_{2}\left(x_{0}\right)$. Now appeal to Corollary 3.3.

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