# ESTIMATES FOR SUMS AND GAPS OF EIGENVALUES OF LAPLACIANS ON MEASURE SPACES 

DA-WEN DENG AND SZE-MAN NGAI


#### Abstract

For Laplacians defined by measures on a bounded domain in $\mathbb{R}^{n}$, we prove analogs of the classical eigenvalue estimates for the standard Laplacian: lower bound of sums of eigenvalues by Li and Yau, and gaps of consecutive eigenvalues by Payne, Pólya and Weinberger. This work is motivated by the study of spectral gaps for Laplacians on fractals.


## 1. Introduction

One of the anomalous behaviors of Laplacians on fractals is the existence of spectral gaps, i.e., $\varlimsup_{k \rightarrow \infty} \lambda_{k+1} / \lambda_{k}>1$, where $\lambda_{k}$ are the eigenvalues. This is not possible for bounded domains on $\mathbb{R}^{n}$ or on Riemannian manifolds. In fact, according to the Weyl law, on a compact connected oriented $n$-dimensional Riemannian manifold $M$,

$$
\left(\lambda_{k}\right)^{n / 2} \sim \frac{(2 \pi)^{n} k}{\omega_{n} \operatorname{vol}(M)} \quad \text { as } k \rightarrow \infty
$$

where $\operatorname{vol}(\mathrm{M})$ is the volume of $M$, and $\omega_{n}$ is volume of the unit ball in $\mathbb{R}^{n}$ (see, e.g., [2]). Consequently, $\lim _{k \rightarrow \infty} \lambda_{k+1} / \lambda_{k}=1$. Strichartz [20] showed that the existence of spectral gaps implies better convergence of Fourier series. Rigorous proofs for the existence of spectral gaps have been obtained for only a limited number of fractals, such as the Sierpiński gasket and the Vicsek set (see [5, 8,21$]$ ). For Laplacians defined by most self-similar measures, especially those with overlaps, it is not clear whether spectral gap exists. This is the main motivation of the present paper. This paper is also a continuation of the work by the authors' [4] and by Pinasco and Scarola [17] on estimating the first eigenvalue of Laplacians with respect to fractal measures.

To describe some classical results, let $\Omega$ be a bounded domain on $\mathbb{R}^{n}$ and let $\lambda_{k}$ be the $k$-th Dirichlet eigenvalue. Li and Yau [10] obtained the following lower estimate for the sum

[^0]of the first $k$ eigenvalues
\[

$$
\begin{equation*}
\sum_{i=1}^{k} \lambda_{i} \geq \frac{n C_{n}}{n+2} k^{(n+2) / n} \operatorname{vol}(\Omega)^{-2 / n} \tag{1.1}
\end{equation*}
$$

\]

where $\operatorname{vol}(\Omega)$ denotes the volume of $\Omega$, and $C_{n}=(2 \pi)^{2} B_{n}^{-2 / n}$, with $B_{n}$ being the volume of the unit ball in $\mathbb{R}^{n}$.

An upper estimate was obtained by Kröger [16]. The results of Li-Yau and Kröger have been extended to homogeneous Riemannian manifolds by Strichartz [19].

For the gaps between consecutive eigenvalues, Payne, Pólya and Weinberger [15] (see also [18]) proved the following estimate for the gaps between two consecutive eigenvalues:

$$
\begin{equation*}
\lambda_{k+1}-\lambda_{k} \leq \frac{4 \sum_{i=1}^{k} \lambda_{i}}{n k} \tag{1.2}
\end{equation*}
$$

The goal of this paper is to prove analogs of (1.1) and (1.2) for Laplacians defined by a measure $\mu$. Let $\Omega \subset \mathbb{R}^{n}$ be a bounded open subset of $\mathbb{R}^{n}$ and $\mu$ be a positive finite Borel measure with $\mu(\Omega)>0$ and with support being contained in $\bar{\Omega}$. Under suitable conditions (see Section 2), $\mu$ defines a Dirichlet Laplacian $-\Delta_{\mu}$; moreover, there exists an orthonormal basis $\left\{\varphi_{n}\right\}$ consisting of eigenfunctions of $-\Delta_{\mu}$ and the eigenvalues $\lambda_{n}$ satisfy $0<\lambda_{1} \leq \lambda_{2} \leq \cdots$, with $\lim _{n \rightarrow \infty} \lambda_{n}=\infty$. We remark that if $\mu$ is the restriction of Lebesgue measure to $\Omega$, then $\Delta_{\mu}$ is the classical Dirichlet Laplacian.

We first prove an analog of the classical lower estimate of the sum of eigenvalues of the standard Laplacian obtained by Li and Yau [10]. We let $L_{\mu}^{2}(\Omega)$ denote the Hilbert space of square-integrable functions with respect to $\mu$. For $u \in L_{\mu}^{2}(\Omega)$, if there is no confusion of what $\Omega$ is, we let

$$
\|u\|_{\mu}=\left(\int_{\Omega}|u|^{2} d \mu\right)^{1 / 2}
$$

If $\mu$ is the restriction of Lebesgue measure to $\Omega$, we denote the corresponding $L^{2}$-space and norm respectively by $L^{2}(\Omega)$ and $\|\cdot\|$.

Theorem 1.1. Let $\Omega \subseteq \mathbb{R}^{n}$ be a bounded domain, $\mu$ be a positive finite Borel measure on $\Omega$ with $\operatorname{supp}(\mu) \subseteq \bar{\Omega},-\Delta_{\mu}$ be the Dirichlet Laplacian defined by $\mu$ described in Section 2, $\lambda_{k}$ be the $k$-th eigenvalue of $-\Delta_{\mu}$, and $\varphi_{k}$ be the corresponding $L_{\mu}^{2}(\Omega)$-normalized eigenfunction. Then

$$
\begin{aligned}
\sum_{j=1}^{k} \lambda_{j} & \geq\left(\sum_{j=1}^{k}\left\|\varphi_{j}\right\|^{2}\right)^{\frac{n+2}{2}}\left(B_{n} \sup _{z \in \mathbb{R}^{n}} \sum_{j=1}^{k}\left|\hat{\varphi}_{j}(z)\right|^{2}\right)^{-\frac{2}{n}} \frac{n}{n+2} \\
& \geq\left(\sum_{j=1}^{k}\left\|\varphi_{j}\right\|^{2}\right)\left((2 \pi)^{-n} \operatorname{vol}(\Omega) B_{n}\right)^{-\frac{2}{n}} \frac{n}{n+2}
\end{aligned}
$$

where $\operatorname{vol}(\Omega)$ is the volume of $\Omega$ and $B_{n}$ is the volume of the unit ball in $\mathbb{R}^{n}$.

Finally, we generalize the classical theorem by Payne, Pólya and Weinberger [15] on the gaps between two consecutive eigenvalues.

Theorem 1.2. Assume the hypotheses of Theorem 1.1. Then for all $k \geq 1$,

$$
\begin{equation*}
\lambda_{k+1}-\lambda_{k} \leq \frac{4 \sum_{i=1}^{k} \lambda_{i}}{n \sum_{i=1}^{k}\left\|\varphi_{i}\right\|^{2}} \tag{1.3}
\end{equation*}
$$

Remark 1.1. We note that in (1.2), $\left\|\varphi_{j}\right\|>0$ for all $i$. In fact, if $\nabla \varphi_{j}=0$, then, in view of the Poincaré inequality for measures (see (2.1)), we would get $\left\|\varphi_{j}\right\|_{\mu}=0$, a contradiction.

Both Theorems 1.1 and 1.2 involve the sum $\sum_{i=1}^{k}\left\|\varphi_{i}\right\|^{2}$, which suggests that it is necessary to study the eigenfunctions. We are not able to obtain a good estimate for this sum. Properties of eigenfunctions in one-dimensional, especially when the support of $\mu$ is an interval, have been studied in [1]). In Section 3, we focus on the case when the support of $\mu$ is not an interval, such as a Cantor-type measure.

This paper is organized as follows. In Section 2, we recall the definition and some elementary properties of the Dirichlet Laplacian $\Delta_{\mu}$ defined on a domain by a measure $\mu$. In Section 3 we prove the min-max principle for $-\Delta_{\mu}$ and some properties of the eigenfunctions in one-dimension. Theorem 1.1 is proved in Section 4. Section 5 is devoted to the proof of Theorem 1.2. Finally in Section 6 we state some comments and open questions.

## 2. Preliminaries

For convenience, we summarize the definition of the Dirichlet Laplacian with respect to a measure $\mu$; details can be found in [9]. Let $\Omega \subset \mathbb{R}^{n}$ be a bounded open subset and $\mu$ be a positive finite Borel measure with $\operatorname{supp}(\mu) \subseteq \bar{\Omega}$ and $\mu(\Omega)>0$. We further suppose $\mu$ satisfies the following Poincaré inequality (PI) for measures: There exists a constant $C>0$ such that

$$
\begin{equation*}
\int_{\Omega}|u|^{2} d \mu \leq C \int_{\Omega}|\nabla u|^{2} d x \quad \text { for all } u \in C_{c}^{\infty}(\Omega) \tag{2.1}
\end{equation*}
$$

Notice that (PI) cannot be immediately extended to $H_{0}^{1}(\Omega)$ functions. For example, let $\Omega=(0,1) \subseteq \mathbb{R}$ and $\mu$ be the standard Cantor measure, which is supported on the Cantor set. For any $u \in H_{0}^{1}(\Omega)$, if we increase the value of $u$ on the Cantor set, $\int_{0}^{1}|\nabla u|^{2} d x$ remains unchanged but $\int_{0}^{1}|u|^{2} d \mu$ can be increased within the same equivalence class of $u$ without bound and hence the inequality cannot hold. However, the following is true. (PI) implies that each equivalence class $u \in H_{0}^{1}(\Omega)$ contains a unique (in $L_{\mu}^{2}(\Omega)$ sense) member $\bar{u}$ that belongs to $L_{\mu}^{2}(\Omega)$ and satisfies both conditions below:
(1) There exists a sequence $\left\{u_{n}\right\}$ in $C_{c}^{\infty}(\Omega)$ such that $u_{n} \rightarrow \bar{u}$ in $H_{0}^{1}(\Omega)$ and $u_{n} \rightarrow \bar{u}$ in $L_{\mu}^{2}(\Omega) ;$
(2) $\bar{u}$ satisfies the inequality in (2.1).

We call $\bar{u}$ the $L_{\mu}^{2}(\Omega)$-representative of $u$. Consider our Cantor set example above. For $u \in H_{0}^{1}(\Omega)$, let $\left\{u_{n}\right\} \subseteq C_{c}^{\infty}(\Omega)$ be a sequence convergent to $u$ and hence Cauchy in $H_{0}^{1}(\Omega)$. By (PI), $\left\{u_{n}\right\}$ is Cauchy and hence convergent in $L_{\mu}^{2}(\Omega)$. Then $\bar{u}$ is the function obtained by redefining $u$ on the Cantor set to be the $L_{\mu}^{2}(\Omega)$ limit of $u_{n}$.

Assume (PI) holds and define a mapping $\iota: H_{0}^{1}(\Omega) \rightarrow L_{\mu}^{2}(\Omega)$ by

$$
\iota(u)=\bar{u}
$$

$\iota$ is a bounded linear operator, but not necessarily injective. Consider the subspace $\mathcal{N}$ of $H_{0}^{1}(\Omega)$ defined as

$$
\mathcal{N}:=\left\{u \in H_{0}^{1}(\Omega):\|\iota(u)\|_{\mu}=0\right\} .
$$

Now let $\mathcal{N}^{\perp}$ be the orthogonal complement of $\mathcal{N}$ in $H_{0}^{1}(\Omega)$. Then $\iota: \mathcal{N}^{\perp} \rightarrow L_{\mu}^{2}(\Omega)$ is injective. Throughout the rest of this paper, unless explicitly stated otherwise, we will use the $L_{\mu}^{2}(\Omega)$-representative $\bar{u}$ of $u$ and denote it simply by $u$.

Consider a nonnegative bilinear form $\mathcal{E}(\cdot, \cdot)$ in $L_{\mu}^{2}(\Omega)$ given by

$$
\begin{equation*}
\mathcal{E}(u, v):=\int_{\Omega} \nabla u \cdot \nabla v d x \tag{2.2}
\end{equation*}
$$

with domain $\operatorname{dom}(\mathcal{E})=\mathcal{N}^{\perp}$, or more precisely, $\iota\left(\mathcal{N}^{\perp}\right)$. (PI) implies that $(\mathcal{E}, \operatorname{dom}(\mathcal{E}))$ is a closed quadratic form on $L_{\mu}^{2}(\Omega)$. Hence, there exists a nonnegative self-adjoint operator $-\Delta_{\mu}$ in $L_{\mu}^{2}(\Omega)$ such that $\operatorname{dom}(\mathcal{E})=\operatorname{dom}\left(\left(-\Delta_{\mu}\right)^{1 / 2}\right)$ and

$$
\mathcal{E}(u, v)=\left\langle\left(-\Delta_{\mu}\right)^{1 / 2} u,\left(-\Delta_{\mu}\right)^{1 / 2} v\right\rangle_{L_{\mu}^{2}(\Omega)} \quad \text { for all } u, v \in \operatorname{dom}(\mathcal{E})
$$

(see [3]). We call $\Delta_{\mu}$ the (Dirichlet) Laplacian with respect to $\mu$. It follows that $u \in \operatorname{dom}\left(\Delta_{\mu}\right)$ and $-\Delta_{\mu} u=f$ if and only if $-\Delta u=f d \mu$ in the sense of distribution: for all $\varphi \in C_{c}^{\infty}(\Omega)$, $\int_{\Omega} \nabla u \cdot \nabla \varphi d x=\int_{\Omega} f \varphi d \mu$ (see [9, Proposition 2.2]). A real number $\lambda \in \mathbb{R}$ is a (Dirichlet) eigenvalue of $-\Delta_{\mu}$ with eigenfunction $f$ if for all $\varphi \in C_{c}^{\infty}(\Omega)$,

$$
\begin{equation*}
\int_{\Omega} \nabla f \cdot \nabla \varphi d x=\lambda \int_{\Omega} f \varphi d \mu \tag{2.3}
\end{equation*}
$$

From [9, Theorem 1.2], when $\mu$ satisfies (PI), there exists an orthonormal basis $\left\{\varphi_{n}\right\}_{n=1}^{\infty}$ of $L_{\mu}^{2}(\Omega)$ consisting of (Dirichlet) eigenfunctions of $-\Delta_{\mu}$. The eigenvalues $\left\{\lambda_{n}\right\}_{n=1}^{\infty}$ satisfy $0<\lambda_{1} \leq \lambda_{2} \leq \cdots$. Moreover, if $\operatorname{dim}(\operatorname{dom} \mathcal{E})=\infty$, then $\lim _{n \rightarrow \infty} \lambda_{n}=\infty$. We have the following characterizations of $\operatorname{dom} \mathcal{E}$ and $\operatorname{dom}\left(-\Delta_{\mu}\right)$ :

$$
\begin{aligned}
& \operatorname{dom} \mathcal{E}=\mathcal{N}^{\perp}=\left\{\sum_{n=1}^{\infty} c_{n} \varphi_{n}: \sum_{n=1}^{\infty} c_{n}^{2} \lambda_{n}<\infty\right\}, \\
& \operatorname{dom}\left(-\Delta_{\mu}\right)=\left\{\sum_{n=1}^{\infty} c_{n} \varphi_{n}: \sum_{n=1}^{\infty} c_{n}^{2} \lambda_{n}^{2}<\infty\right\} .
\end{aligned}
$$

To state a sufficient condition for (PI), we recall that the lower $L^{\infty}$-dimension of a measure $\mu$ is defined by

$$
\underline{\operatorname{dim}}_{\infty}(\mu)=\liminf _{\delta \rightarrow 0^{+}} \frac{\ln \left(\sup _{x} \mu\left(B_{\delta}(x)\right)\right)}{\ln \delta}
$$

where the supremum is taken over all $x \in \operatorname{supp}(\mu)$.
Theorem 2.1. ( [9, Theorems 1.1, 1.2]) Let $\Omega \subseteq \mathbb{R}^{n}$ be a bounded open set and $\mu$ be a positive finite Borel measure on $\mathbb{R}^{n}$ with $\operatorname{supp}(\mu) \subseteq \bar{\Omega}$ and $\mu(\Omega)>0$. Assume $\operatorname{dim}_{\infty}(\mu)>n-2$.
(a) (PI) holds. In particular, if $n=1$, or $n=2$ and $\mu$ is upper $s$-regular with $s>0$, or $\mu$ is absolutely continuous with bounded density, (PI) holds.
(b) The set of eigenvalues of $-\Delta_{\mu}$ is contained in $(0, \infty)$ and has no accumulation point. Hence $-\Delta_{\mu}$ has a positive smallest eigenvalue $\lambda_{1}^{\mu}$.

## 3. The min-max principle and properties of eigenfunctions

Let $\Omega$ be a bounded domain in $\mathbb{R}^{n}$ and $\mu$ be a positive finite Borel measure with $\operatorname{supp}(\mu) \subseteq$ $\bar{\Omega}$. In this section we extend the variational principle for the principal eigenvalue and Courant's min-max principle for the $k$-th eigenvalue to the Laplacians $\Delta_{\mu}$. This will be needed in the proof of Theorem 1.2. We first introduce some notation that will be needed in the proof of the theorem. Let $(\cdot, \cdot)_{\mu}$ denote the inner product in $L_{\mu}^{2}(\Omega)$. Also, for any subset $S \subseteq L_{\mu}^{2}(\Omega)$, let $\langle S\rangle$ be the vector subspace of $L_{\mu}^{2}(\Omega)$ spanned by $S$, and let $S^{\perp}$ be the orthogonal complement of $S$ in $L_{\mu}^{2}(\Omega)$.

Theorem 3.1. (Min-Max principle) Let $\Delta_{\mu}$ be the Dirichlet Laplacian defined on a bounded domain $\Omega \subseteq \mathbb{R}^{n}$ and let $0<\lambda_{1} \leq \lambda_{2} \leq \cdots$ be the eigenvalues. Then for $k=1,2, \ldots$, the $k$-th eigenvalue satisfies

$$
\begin{equation*}
\lambda_{k}=\max _{S \in \Sigma_{k-1}} \min \left\{\mathcal{E}(u, u): u \in S^{\perp},\|u\|_{\mu}=1\right\} \tag{3.1}
\end{equation*}
$$

where $\Sigma_{k-1}$ is the collection of all $(k-1)$-dimensional subspaces of $\operatorname{dom}\left(-\Delta_{\mu}\right)$. In particular, for $k=1$ we have the variational principle for the principal eigenvalue:

$$
\begin{equation*}
\lambda_{1}=\min \left\{\mathcal{E}(u, u): u \in \operatorname{dom}\left(-\Delta_{\mu}\right),\|u\|_{\mu}=1\right\} . \tag{3.2}
\end{equation*}
$$

Proof. Step 1. Let $\left\{\varphi_{k}\right\} \subseteq L_{\mu}^{2}(\Omega)$ be an orthonormal basis of $L_{\mu}^{2}(\Omega)$ with $\left\{\varphi_{k}\right\} \subset \operatorname{dom}\left(-\Delta_{\mu}\right)$ satisfying

$$
\begin{cases}-\Delta_{\mu} \varphi_{k}=\lambda_{k} \varphi_{k} & \text { in } \Omega  \tag{3.3}\\ \varphi_{k}=0 & \text { in } \partial \Omega\end{cases}
$$

for $k=1,2, \ldots$ Hence

$$
\begin{equation*}
\mathcal{E}\left(\varphi_{k}, \varphi_{l}\right)=\left(\lambda_{k} \varphi_{k}, \varphi_{l}\right)_{\mu}=\lambda_{k} \delta_{k l}, \tag{3.4}
\end{equation*}
$$

where $\delta_{k l}$ is the Kronecker delta. If $u \in \operatorname{dom} \mathcal{E}$ and $\|u\|_{\mu}=1$, then

$$
\begin{equation*}
u=\sum_{k=1}^{\infty} d_{k} \varphi_{k} \quad \text { in } L_{\mu}^{2}(\Omega) \tag{3.5}
\end{equation*}
$$

where $d_{k}=\left(u, \varphi_{k}\right)_{\mu}$, and the equality holds in the sense that $\left\|u-\sum_{k=1}^{N} d_{k} \varphi_{k}\right\|_{\mu} \rightarrow 0$ as $N \rightarrow \infty$. Also,

$$
\begin{equation*}
\sum_{k=1}^{\infty} d_{k}^{2}=\|u\|_{\mu}^{2}=1 \tag{3.6}
\end{equation*}
$$

Step 2. Equation (3.4) implies that $\left\{\varphi_{k} / \lambda_{k}^{1 / 2}\right\}_{k=1}^{\infty} \subseteq \operatorname{dom}\left(-\Delta_{\mu}\right)$ is an orthonormal set with respect to the inner product $\mathcal{E}(\cdot, \cdot)$. We claim that $\left\{\varphi_{k} / \lambda_{k}^{1 / 2}\right\}_{k=1}^{\infty}$ is an orthonormal basis of $\operatorname{dom}\left(-\Delta_{\mu}\right)$ with respect to $\mathcal{E}(\cdot, \cdot)$. To see this, we let $u \in \operatorname{dom}\left(-\Delta_{\mu}\right)$ such that $\mathcal{E}\left(\varphi_{k}, u\right)=0$ for all $k \geq 1$. Then $\left(\lambda_{k} \varphi_{k}, u\right)_{\mu}=0$ for all $k \geq 1$ implies that $\left(u_{k}, u\right)_{\mu}=0$ for all $k \geq 1$. Thus, $u=0 \mu$-a.e. on $\Omega$, which implies that $u=0$ in $\operatorname{dom}\left(-\Delta_{\mu}\right)$ (in the $H_{0}^{1}$-norm), since $\iota: \mathcal{N}^{\perp} \rightarrow L_{\mu}^{2}(\Omega)$ is an injection. Thus, for all $u \in \operatorname{dom}\left(-\Delta_{\mu}\right)$,

$$
\begin{equation*}
u=\sum_{k=1}^{\infty} a_{k} \frac{\varphi_{k}}{\lambda_{k}^{1 / 2}}, \tag{3.7}
\end{equation*}
$$

where $a_{k}=\mathcal{E}\left(u, \varphi_{k} / \lambda_{k}^{1 / 2}\right)$. Observe that

$$
a_{k}=\mathcal{E}\left(\frac{\varphi_{k}}{\lambda_{k}^{1 / 2}}, u\right)=\frac{1}{\lambda_{k}^{1 / 2}}\left(\lambda_{k} \varphi_{k}, u\right)_{\mu}=\frac{1}{\lambda_{k}^{1 / 2}}\left(\lambda_{k} \varphi_{k}, \sum_{\ell=1}^{\infty} d_{\ell} \varphi_{\ell}\right)_{\mu}=\lambda_{k}^{1 / 2} d_{k} .
$$

Hence,

$$
\begin{equation*}
u=\sum_{k=1}^{\infty}\left(\lambda_{k}^{1 / 2} d_{k}\right) \frac{\varphi_{k}}{\lambda_{k}^{1 / 2}}=\sum_{k=1}^{\infty} d_{k} \varphi_{k} \quad \text { in } \operatorname{dom}\left(-\Delta_{\mu}\right) \tag{3.8}
\end{equation*}
$$

Equations (3.8) and (3.4) imply that

$$
\mathcal{E}(u, u)=\mathcal{E}\left(\sum_{k=1}^{\infty} d_{k} \varphi_{k}, \sum_{k=1}^{\infty} d_{k} \varphi_{k}\right)=\sum_{k=1}^{\infty} d_{k}^{2} \lambda_{k} \geq \lambda_{1} .
$$

Since $\mathcal{E}\left(\varphi_{1}, \varphi_{1}\right)=\lambda_{1}$, (3.2) follows.
Step 3. Now let $\left\{\varphi_{k}\right\}_{k=1}^{\infty} \subseteq L_{\mu}^{2}(\Omega)$ be an orthonormal basis of $L_{\mu}^{2}(\Omega)$ satisfying (3.3). Let $v \in \operatorname{dom}\left(-\Delta_{\mu}\right)$ with $\|v\|_{\mu}=1$. Write $v=\sum_{i=1}^{\infty} c_{i} \varphi_{i}$, where the equality holds in both $L_{\mu}^{2}(\Omega)$ and $\operatorname{dom}\left(-\Delta_{\mu}\right)$. As $\mathcal{E}(v, v)=\sum_{i=1}^{\infty} \lambda_{i} c_{i}^{2}$,

$$
\begin{aligned}
\lambda_{k} & =\min \left\{\sum_{i=1}^{\infty} \lambda_{i} c_{i}^{2}: c_{1}=\cdots=c_{k-1}=0, \sum_{i=1}^{\infty} c_{i}^{2}=1\right\} \\
& =\min \left\{\mathcal{E}(v, v): v \in\left\langle\varphi_{1}, \ldots, \varphi_{k-1}\right\rangle^{\perp},\|v\|_{\mu}=1\right\} .
\end{aligned}
$$

We claim that this is equal to

$$
\max _{S \in \Sigma_{k-1}} \min \left\{\mathcal{E}(v, v): v \in S^{\perp},\|v\|_{\mu}=1\right\}
$$

To prove this, let $S \in \Sigma_{k-1}$ and consider the following two cases.

Case 1. $S=\left\langle v_{1}, \ldots, v_{k-1}\right\rangle$ with $v_{\ell}=\sum_{i=1}^{\infty} d_{\ell i} \varphi_{i}$ for $\ell=1, \ldots, k-1$ and $\operatorname{det}\left(d_{\ell i}\right)_{\ell, i=1}^{k-1}=0$.
In this case, there exist $c_{1}, \ldots, c_{k-1}$, not all zero, such that

$$
\sum_{i=1}^{k-1} c_{i}^{2}=1 \quad \text { and } \quad \sum_{i=1}^{k-1} d_{\ell i} c_{i}=0, \quad \ell=1, \ldots, k-1
$$

Let $\bar{v}:=\sum_{i=1}^{k-1} c_{i} \varphi_{i} \in S^{\perp}$. Then $\|\bar{v}\|_{\mu}=1$ and

$$
\min _{v \in S^{\perp},\|v\|_{\mu}=1} \mathcal{E}(v, v) \leq(\bar{v}, \bar{v})=\sum_{i=1}^{k-1} \lambda_{i} c_{i}^{2} \leq \lambda_{k}=\min _{v \in\left\langle\varphi_{1}, \ldots, \varphi_{k-1}\right\rangle^{\perp}} \mathcal{E}(v, v)
$$

Case 2. $S=\left\langle v_{1}, \ldots, v_{k-1}\right\rangle, v_{\ell}=\sum_{i=1}^{\infty} d_{\ell i} \varphi_{i}$ for $\ell=1, \ldots, k-1$ with $\operatorname{det}\left(d_{\ell i}\right)_{\ell, i=1}^{k-1} \neq 0$.
Let $v \in\left\langle\varphi_{1}, \ldots, \varphi_{k-1}\right\rangle^{\perp}$ with $\|v\|_{\mu}=1$, i.e., $v=\sum_{i=k}^{\infty} c_{i} \varphi_{i}$ with $\sum_{i=k}^{\infty} c_{i}^{2}=1$ and $\mathcal{E}(v, v)=$ $\sum_{i=k}^{\infty} \lambda_{i} c_{i}^{2}$. We will find $v^{\prime} \in S^{\perp}$ with $\left\|v^{\prime}\right\|_{\mu}=1$ such that $\mathcal{E}\left(v^{\prime}, v^{\prime}\right) \leq \mathcal{E}(v, v)$.

Let $\bar{v}=\sum_{i=1}^{\infty} c_{i} \varphi_{i}$ (i.e., $\bar{v}$ has the same $\varphi_{i}$ components as $v$ for $i \geq k$ ) such that $\bar{v} \cdot v_{\ell}=0$ for all $\ell=1, \ldots, k$. That is,

$$
\sum_{i=1}^{k-1} d_{\ell i} c_{i}=\sum_{i=k}^{\infty} d_{\ell i} c_{i}, \quad \ell=1, \ldots, k-1
$$

Since $\operatorname{det}\left(d_{\ell i}\right)_{\ell, i=1}^{k-1} \neq 0$, a solution $c_{1}, \ldots, c_{k-1}$ exists (can be all 0 ). Let $v^{\prime}:=\bar{v} /\|\bar{v}\|_{\mu}$. Note that

$$
\|\bar{v}\|_{\mu}^{2}=\sum_{i=1}^{k-1} c_{i}^{2}+\sum_{i=k}^{\infty} c_{i}^{2}=1+\sum_{i=1}^{k-1} c_{i}^{2}
$$

and

$$
\sum_{j=1}^{k-1} \lambda_{j} c_{j}^{2} \leq\left(\sum_{j=1}^{k-1} c_{j}^{2}\right) \lambda_{k}=\left(\sum_{j=1}^{k-1} c_{j}^{2}\right)\left(\sum_{j=k}^{\infty} c_{j}^{2}\right) \lambda_{k} \leq\left(\sum_{j=1}^{k-1} c_{j}^{2}\right)\left(\sum_{j=k}^{\infty} \lambda_{j} c_{j}^{2}\right)
$$

Thus,

$$
\begin{aligned}
\mathcal{E}\left(v^{\prime}, v^{\prime}\right) & =\frac{1}{\|\bar{v}\|_{\mu}^{2}} \mathcal{E}(\bar{v}, \bar{v})=\frac{1}{1+\sum_{i=1}^{k-1} c_{i}^{2}} \sum_{j=1}^{\infty} \lambda_{j} c_{j}^{2}=\frac{1}{1+\sum_{i=1}^{k-1} c_{i}^{2}}\left(\sum_{j=1}^{k-1} \lambda_{j} c_{j}^{2}+\sum_{j=k}^{\infty} \lambda_{j} c_{j}^{2}\right) \\
& \leq \sum_{j=k}^{\infty} \lambda_{j} c_{j}^{2}=\mathcal{E}(v, v) .
\end{aligned}
$$

It follows that

$$
\min _{v^{\prime} \in S^{\perp},\left\|v^{\prime}\right\|_{\mu}=1} \mathcal{E}\left(v^{\prime}, v^{\prime}\right) \leq \min _{v \in\left\langle\varphi_{1}, \ldots, \varphi_{k}\right\rangle^{\perp},\|v\|_{\mu}=1} \mathcal{E}(v, v) .
$$

This completes the proof.

The following proposition establishes some properties of eigenfunctions in one-dimension, some of them being specific for measures on bounded domains in $\mathbb{R}$. Additional properties
of eigenfunctions can be found in [1]. Let $\mathcal{L}^{1}(E)$ be the one-dimensional Lebesgue measure of a subset $E \subseteq \mathbb{R}$.

Proposition 3.2. Let $\Omega \subset \mathbb{R}$ be a bounded open interval, $\mu$ be a positive finite Borel measure on $\Omega$ with $\operatorname{supp}(\mu) \subseteq \bar{\Omega}, \Delta_{\mu}$ be the Dirichlet Laplacian with respect to $\mu$, and $\varphi \in H_{0}^{1}(\Omega)$ be an eigenfunction of $-\Delta_{\mu}$, i.e., there exists $\lambda \in \mathbb{R}$ such that $-\Delta_{\mu} \varphi=\lambda \varphi$. Then
(a) $\varphi \in C^{0,1 / 2}(\Omega)$;
(b) $\varphi$ is linear over any component of $\Omega \backslash \operatorname{supp}(\mu)$;
(c) if $\mathcal{L}^{1}(\operatorname{supp}(\mu))=0$, then $\varphi \notin C^{2}(\Omega)$, and in fact, $\varphi^{\prime}$ is not absolutely continuous (with respect to Lebesgue measure);
(d) eigenfunctions corresponding to the first eigenvalue do not change sign;
(e) the first eigenvalue is simple.

Proof. (a) It follows directly from Sobolev's embedding theorem that $H_{0}^{1}(\Omega) \hookrightarrow C^{0,1 / 2}(\Omega)$.
(b) Consider a component $(a, b)$ of $\Omega \backslash \operatorname{supp}(\mu)$. For all $v \in C_{c}^{\infty}(a, b) \subseteq C_{c}^{\infty}(\Omega)$,

$$
\int_{\Omega} \varphi^{\prime} v^{\prime} d x=\lambda \int_{\Omega} \varphi v d \mu=0
$$

and hence it also holds for continuous piecewise linear $v$ with $\operatorname{supp}(v) \subset(a, b)$. Note that $\left.\varphi^{\prime}\right|_{(a, b)} \in L^{2}(a, b) \subset L^{1}(a, b)$. Let

$$
a<x-\delta<x_{1}<x_{1}+\delta<x_{2}-\delta<x_{2}<x_{2}+\delta<b
$$

Let $v \in C(a, b)$ that is equal to 0 on $\left(a, x_{1}-\delta\right) \cup\left(x_{2}+\delta, b\right)$, equal to 1 on $\left(x_{1}+\delta, x_{2}-\delta\right)$, and linear over $\left(x_{1}-\delta, x_{1}+\delta\right)$ and $\left(x_{2}-\delta, x_{2}+\delta\right)$. Then

$$
\begin{aligned}
0 & =\int_{a}^{b} \varphi^{\prime} v^{\prime} d y=\int_{x_{1}-\delta}^{x_{1}+\delta} \varphi^{\prime}(y) v^{\prime}(y) d y+\int_{x_{2}-\delta}^{x_{2}+\delta} \varphi^{\prime}(y) v^{\prime}(y) d y \\
& =\frac{1}{2 \delta} \int_{x_{1}-\delta}^{x_{1}+\delta} \varphi^{\prime}(y) d y-\frac{1}{2 \delta} \int_{x_{2}-\delta}^{x_{2}+\delta} \varphi^{\prime}(y) d y
\end{aligned}
$$

By the Lebesgue differentiation theorem, for Lebesgue a.e. $x \in(a, b), \varphi^{\prime}(x)=c$, a constant. As $\varphi \in H_{0}^{1}(\Omega)$ is absolutely continuous, for all $x \in(a, b)$,

$$
\varphi(x)=\varphi(a)+\int_{a}^{x} \varphi^{\prime}(y) d y=\varphi(a)+c(x-a)
$$

completing the proof of (b).
(c) By part (b), $\varphi^{\prime \prime}=0$ Lebesgue a.e. Hence, if $\varphi^{\prime}$ is absolutely continuous, then for Lebesgue a.e. $a, b \in \Omega$,

$$
\varphi^{\prime}(a)-\varphi^{\prime}(b)=\int_{a}^{b} \varphi^{\prime \prime}(y) d y=0
$$

Thus, $\varphi^{\prime}$ is a constant. Since $\varphi \in H_{0}^{1}(\Omega)$, we conclude that $\varphi \equiv 0$, contradicting that $\varphi$ is an eigenfunction.
(d) By [1, Proposition 3.4], eigenfunctions corresponding to the first eigenvalue are concave or convex. As they vanish at the end points, they do not change sign.
(e) Let $\varphi_{1}$ and $\varphi_{2}$ be normalized functions corresponding to the first eigenvalue of $-\Delta_{\mu}$. Then $\varphi_{1}, \varphi_{2} \in C^{0, \alpha}(\Omega) \subset C(\bar{\Omega})$. If $\varphi_{1} \equiv \varphi_{2}$ on $\operatorname{supp}(\mu)$, then by linearity on components of $\Omega \backslash \operatorname{supp}(\mu)$ and continuity, $\varphi_{1} \equiv \varphi_{2}$ on $\Omega$. Thus $\varphi_{1} \not \equiv \varphi_{2}$ if and only if $\varphi_{1} \not \equiv \varphi_{2}$ on $\operatorname{supp}(\mu)$.

Suppose that $\varphi_{1}$ and $\varphi_{2}$ are of the same sign, say positive, and $\varphi_{1} \not \equiv \varphi_{2}$. If $\varphi_{1} \geq \varphi_{2}$, then $\varphi_{1}>\varphi_{2}$ on some subset $E \subset \Omega$ with $\mu(E)>0$. Hence

$$
1=\int_{\Omega}\left|\varphi_{1}\right|^{2} d \mu>\int_{\Omega}\left|\varphi_{2}\right|^{2} d \mu=1
$$

a contradiction. Thus, there exist $x_{1}, x_{2} \in \Omega$ such that $\varphi_{1}\left(x_{1}\right)>\varphi_{2}\left(x_{1}\right)$ and $\varphi_{1}\left(x_{2}\right)<\varphi_{2}\left(x_{2}\right)$. Now $\varphi=\varphi_{1}-\varphi_{2}$ is an eigenfunction with $\varphi\left(x_{1}\right)>0$ and $\varphi\left(x_{2}\right)<0$, contradicting (d).

## 4. Lower estimate of sums of eigenvalues

We will use the a lemma from [10] which says that if $f$ is a real-valued function defined on $\mathbb{R}^{n}$ with $0 \leq f \leq M_{1}$, and

$$
\int_{\mathbb{R}^{n}}|z|^{2} d z \leq M_{2}
$$

then

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} f(z) d z \leq\left(M_{1} B_{n}\right)^{\frac{2}{n+2}} M_{2}^{\frac{n}{n+2}}\left(\frac{n+2}{n}\right)^{\frac{n}{n+2}} \tag{4.1}
\end{equation*}
$$

Proof of Theorem 1.1. Let

$$
\begin{equation*}
\Phi(x, y)=\sum_{j=1}^{k} \varphi_{j}(x) \varphi_{j}(y), x, y \in \Omega \quad \text { and } \quad f(z):=\int_{\Omega}|\hat{\Phi}(z, y)|^{2} d \mu(y), z \in \mathbb{R}^{n} \tag{4.2}
\end{equation*}
$$

where

$$
\hat{\Phi}(z, y)=(2 \pi)^{-n / 2} \int_{x \in \mathbb{R}^{n}} \Phi(x, y) e^{-i x \cdot z} d x
$$

is the Fourier transform of $\Phi$ at $z$ and each $\varphi_{j}$ is extended to $\mathbb{R}^{n}$ by setting it equal to 0 on $\mathbb{R}^{d} \backslash \Omega$. Then, using the linearity of the Fourier transform,

$$
\begin{align*}
& f(z)=\int_{\Omega}\left|\sum_{j=1}^{k} \hat{\varphi}_{j}(z) \varphi_{j}(y)\right|^{2} d \mu(y) \\
&=\int_{\Omega} \sum_{j, \ell=1}^{k} \hat{\varphi}_{j}(z) \varphi_{j}(y) \overline{\hat{\varphi}_{\ell}}(z)  \tag{4.3}\\
& \varphi_{\ell}(y) d \mu(y) \\
&=\sum_{j=1}^{k}\left|\hat{\varphi}_{j}(z)\right|^{2}
\end{align*}
$$

Let $M_{1}:=\sup _{z \in \mathbb{R}^{n}} \sum_{j=1}^{k}\left|\hat{\varphi}_{j}(z)\right|^{2}$. Then it follows that for all $z \in \mathbb{R}^{n}$,

$$
\begin{equation*}
0 \leq f(z) \leq M_{1} \leq(2 \pi)^{-n} \sum_{j=1}^{k}\left\|\varphi_{j}\right\|_{L^{1}(\Omega)}^{2} \leq(2 \pi)^{-n} \operatorname{vol}(\Omega) \sum_{j=1}^{k}\left\|\varphi_{j}\right\|^{2} \tag{4.4}
\end{equation*}
$$

Also,

$$
\begin{align*}
\int_{\mathbb{R}^{n}}|z|^{2} f(z) d z & =\int_{\Omega} \int_{\mathbb{R}^{n}}|z|^{2}|\hat{\Phi}(z, y)|^{2} d z d \mu(y) \\
& =\int_{\Omega} \int_{\mathbb{R}^{n}}\left|\widehat{\nabla_{x} \Phi}(z, y)\right|^{2} d z d \mu(y) \\
& =\int_{\Omega} \int_{\Omega}\left|\nabla_{x} \Phi(x, y)\right|^{2} d x d \mu(y) \quad \text { (Plancherel's theorem) }  \tag{4.5}\\
& =\int_{\Omega} \int_{\Omega}\left(\sum_{j=1}^{k} \lambda_{j} \varphi_{j}(x) \varphi_{j}(y)\right)\left(\sum_{\ell=1}^{k} \varphi_{\ell}(x) \varphi_{\ell}(y)\right) d \mu(x) d \mu(y) \\
& =\sum_{j=1}^{k} \lambda_{j}=: M_{2}
\end{align*}
$$

where the fourth equality follows from (2.3). Using (4.3) followed by the Plancherel theorem, we get

$$
\begin{equation*}
\int_{\mathbb{R}^{n}} f(z) d z=\sum_{j=1}^{k}\left\|\hat{\varphi}_{j}(z)\right\|^{2}=\sum_{j=1}^{k}\left\|\varphi_{j}\right\|^{2} \tag{4.6}
\end{equation*}
$$

By applying [10, Lemma 1] (see (4.1)) to the function $f$ in (4.2), and using (4.5) and (4.6), we get

$$
\begin{aligned}
\sum_{j=1}^{k}\left\|\varphi_{j}\right\|^{2} & =\int_{\mathbb{R}^{n}} f(z) d z \leq\left(M_{1} B_{n}\right)^{\frac{2}{n+2}} M_{2}^{\frac{n}{n+2}}\left(\frac{n+2}{n}\right)^{\frac{n}{n+2}} \\
& =\left(B_{n} \sup _{z \in \mathbb{R}^{n}} \sum_{j=1}^{k}\left|\hat{\varphi}_{j}(z)\right|^{2}\right)^{\frac{2}{n+2}}\left(\sum_{j=1}^{k} \lambda_{j}\right)^{\frac{n}{n+2}}\left(\frac{n+2}{n}\right)^{\frac{n}{n+2}}
\end{aligned}
$$

Thus, using (4.4), we get

$$
\begin{aligned}
\sum_{j=1}^{k} \lambda_{j} & \geq\left(\sum_{j=1}^{k}\left\|\varphi_{j}\right\|^{2}\right)^{\frac{n+2}{n}}\left(B_{n} \sup _{z \in \mathbb{R}^{n}} \sum_{j=1}^{k}\left|\hat{\varphi}_{j}(z)\right|^{2}\right)^{-\frac{2}{n}}\left(\frac{n}{n+2}\right) \\
& \geq\left(\sum_{j=1}^{k}\left\|\varphi_{j}\right\|^{2}\right)\left((2 \pi)^{-n} \operatorname{vol}(\Omega) B_{n}\right)^{-\frac{2}{n}}\left(\frac{n}{n+2}\right)
\end{aligned}
$$

which completes the proof.

## 5. Upper estimate of gaps of eigenvalues

This section is devoted to generalizing the estimate of Payne, Pólya and Weinberger in 1.2 to Laplacians with respect to measures. We use the same notation as in Theorem 1.1.

Proof of Theorem 1.2. By the min-max principle,

$$
\begin{equation*}
\lambda_{k+1}=\inf \left\{\frac{\int_{\Omega}|\nabla v|^{2} d x}{\int_{\Omega} v^{2} d \mu}: \int_{\Omega} v \varphi_{i} d \mu=0, i=1, \ldots, k, v \in \operatorname{dom}\left(-\Delta_{\mu}\right)\right\} \tag{5.1}
\end{equation*}
$$

Choose test functions $v_{i}=g \varphi_{i}-\sum_{j=1}^{k} a_{i j} \varphi_{j}$, where $g$ will be chosen later. Obviously, $v_{i} \in \operatorname{dom} \mathcal{E}$. Suppose $\int_{\Omega} v_{i} \varphi_{\ell} d \mu=0$. Then

$$
0=\int_{\Omega} g \varphi_{i} \varphi_{\ell} d \mu-\sum_{j=1}^{k} a_{i j} \int_{\Omega} \varphi_{j} \varphi_{\ell} d \mu=\int_{\Omega} g \varphi_{i} \varphi_{\ell} d \mu-a_{i \ell}
$$

Therefore, $a_{i \ell}=a_{\ell i}$ and

$$
\int_{\Omega} v_{i}^{2} d \mu=\int_{\Omega}\left(g \varphi_{i} v_{i}-\sum_{j=1}^{k} a_{i j} \varphi_{j} v_{i}\right) d \mu=\int_{\Omega} g \varphi_{i} v_{i} d \mu
$$

Notice that

$$
\int_{\Omega}\left|\nabla v_{i}\right|^{2} d x=\left\langle-\Delta v_{i}, v_{i}\right\rangle_{H^{-1}, H_{0}^{1}}
$$

Since

$$
\begin{aligned}
\Delta v_{i} & =(\Delta g) \varphi_{i}+2 \nabla g \cdot \nabla \varphi_{i}+g \Delta \varphi_{i}-\sum_{j=1}^{k} a_{i j} \Delta \varphi_{j} \\
& =(\Delta g) \varphi_{i}+2 \nabla g \cdot \nabla \varphi_{i}-\lambda_{i} g \varphi_{i} d \mu+\sum_{i=1}^{k} a_{i j} \lambda_{j} \varphi_{j} d \mu
\end{aligned}
$$

we have

$$
\int_{\Omega}\left|\nabla v_{i}\right|^{2} d x=-\int_{\Omega}(\Delta g) v_{i} \varphi_{i} d x-2 \int_{\Omega} v_{i} \nabla g \cdot \nabla \varphi_{i} d x+\lambda_{i} \int_{\Omega} g \varphi_{i} v_{i} d \mu
$$

Now,

$$
\begin{align*}
-2 \sum_{i=1}^{k} \int_{\Omega} v_{i} \nabla g \cdot \varphi_{i} d x & =-2 \sum_{i=1}^{k} g \nabla u \cdot \varphi_{i} \nabla \varphi_{i} d x+2 \sum_{j, i} a_{i j} \int_{\Omega} \varphi_{j} \nabla \varphi_{i} \cdot \nabla g d x \\
& =\sum_{i=1}^{k}\left(-\frac{1}{2} \int_{\Omega} \nabla g^{2} \cdot \nabla \varphi_{i}^{2} d x\right)+\sum_{i, j} a_{i j} \int_{\Omega} \nabla\left(\varphi_{i} \varphi_{j}\right) \nabla g d x  \tag{5.2}\\
& =\sum_{i=1}^{k} \frac{1}{2} \int \varphi_{i}^{2} \Delta g^{2} d x-\sum_{i, j=1}^{k} a_{i j} \int_{\Omega} \varphi_{i} \varphi_{j} \Delta g d x
\end{align*}
$$

Hence,

$$
\begin{aligned}
\lambda_{k+1} \int_{\Omega} v_{i}^{2} d \mu \leq & \sum_{i=1}^{k} \int|\nabla v|^{2} d x=\sum_{i=1}^{k}\left(-\int v_{i} \varphi_{i} \Delta g d x\right)+\frac{1}{2} \sum_{i=1}^{k} \int \varphi_{i}^{2} \Delta g^{2} d x \\
& -\sum_{i, j=1}^{k} a_{i j} \int_{\Omega} \varphi_{i} \varphi_{j} \Delta g d x+\sum_{i=1}^{k} \lambda_{i} \int_{\Omega} g \varphi_{i} v_{i} d \mu \\
= & -\sum_{i=1}^{k} \int_{\Omega} \varphi_{i}^{2} g \Delta g d x+\sum_{i, j=1}^{k} a_{i j} \int_{\Omega} \varphi_{i} \varphi_{m} j \Delta g d x+\frac{1}{2} \sum_{i=1}^{k} \int_{\Omega} \varphi_{i}^{2} \Delta g^{2} d x \\
& \quad-\sum_{i, j=1}^{k} a_{i j} \int_{\Omega} \varphi_{i} \varphi_{j} \Delta g d x+\sum_{i=1}^{k} \lambda_{i} \int g \varphi_{i} v_{i} d \mu \\
= & \sum_{i=1}^{k} \int_{\Omega} \varphi_{i}^{2}|\nabla g|^{2}+\sum_{i=1}^{k} \lambda_{i} \int v_{i}^{2} d \mu \\
\leq & \sum_{i=1}^{k} \int_{\Omega} \varphi_{i}^{2}|\nabla g|^{2} d x+\lambda_{k} \sum_{i=1}^{k} \int_{\Omega} v_{i}^{2} d \mu .
\end{aligned}
$$

Since for all $i$,

$$
\lambda_{k+1} \int_{\Omega} v_{i}^{2} d \mu \leq \int_{\Omega}\left|\nabla v_{i}\right|^{2} d x
$$

we have

$$
\lambda_{k+1}-\lambda_{k} \leq \frac{\sum_{i=1}^{k} \int_{\Omega} \varphi_{i}^{2}|\nabla g|^{2} d x}{\sum_{i=1}^{k} \int_{\Omega} v_{i}^{2} d \mu}
$$

Now take $g=g_{a}(x)=\sum_{i=1}^{n} a_{i} x_{i}$ with $\sum_{i=1}^{n} a_{i}^{2}=1$. Then $\Delta g=0$ and $|\nabla g|=1$. It follows that

$$
\begin{equation*}
\lambda_{k+1}-\lambda_{k} \leq \frac{\sum_{i=1}^{k} \int_{\Omega} \varphi_{i}^{2} d x}{\sum_{i=1}^{k} v_{i a}^{2} d \mu}=\frac{\sum_{i=1}^{k}\left\|\varphi_{i}\right\|^{2}}{\sum_{i=1}^{k} \int_{\Omega} v_{i a}^{2} d \mu} \tag{5.3}
\end{equation*}
$$

where $v_{i a}=g_{a}(x) \varphi_{i}-\sum_{j=1}^{k} a_{i j} \varphi_{j}$. From (5.2),

$$
\begin{aligned}
\sum_{i=1}^{k} \int_{\Omega} \varphi_{i}^{2} d x & =\sum_{i=1}^{k}\left\|\varphi_{i}\right\|^{2}=-2 \sum_{i=1}^{k} \int_{\Omega} v_{i a}\left(\nabla g_{a} \cdot \nabla \varphi_{i}\right) d x \\
& =-2 \sum_{i=1}^{k} \int_{\Omega} v_{i a}\left(\sum_{j=1}^{n} a_{j} \frac{\partial \varphi_{i}}{\partial x_{j}}\right) d x
\end{aligned}
$$

Normalize the measure on $S^{n-1}$ so that $\int_{S^{n-1}} d S_{a}=1$. Then

$$
\begin{aligned}
\frac{1}{2} \sum_{i=1}^{k}\left\|\varphi_{i}\right\|^{2} & =-\sum_{i=1}^{k} \int_{S^{n-1}} \int_{\Omega} v_{i a}\left(\sum_{j=1}^{n} a_{j} \frac{\partial \varphi_{i}}{\partial x_{j}}\right) d x d S_{a} \\
& =-\int_{S^{n-1}} \int_{\Omega}\left(\sum_{i=1}^{k} v_{i a}\left(\sum_{j=1}^{n} a_{j} \frac{\partial \varphi_{i}}{\partial x_{j}}\right)\right) d x d S_{a} \\
& \leq\left(\int_{\Omega} \int_{S^{n-1}} \sum_{i=1}^{k} v_{i a}^{2} d S_{a} d x\right)^{1 / 2}\left(\int_{\Omega} \int_{S^{n-1}} \sum_{i=1}^{k}\left(\sum_{j=1}^{n} a_{j} \frac{\partial \varphi_{i}}{\partial x_{j}}\right)^{2} d S_{a} d x\right)^{1 / 2}
\end{aligned}
$$

Obviously, $\int_{S^{n-1}} a_{j} a_{\ell}=\delta_{j \ell} / N$. Hence

$$
\begin{aligned}
\int_{\Omega} \int_{S^{n-1}} \sum_{i=1}^{k}\left(\sum_{j=1}^{n} a_{j} \frac{\partial \varphi_{i}}{\partial x_{j}}\right)^{2} d S_{a} d x & =\sum_{i} \int_{\Omega} \frac{1}{n} \sum_{j=1}^{n}\left(\frac{\partial \varphi_{i}}{\partial x_{j}}\right)^{2} d x=\sum_{i=1}^{k} \frac{1}{n} \int_{\Omega}\left|\nabla \varphi_{i}\right|^{2} d x \\
& =\frac{1}{n} \sum_{i=1}^{k} \lambda_{i} \int_{\Omega} \varphi_{i}^{2} d \mu=\frac{1}{n} \sum_{i=1}^{k} \lambda_{i}
\end{aligned}
$$

It follows that

$$
\frac{1}{4}\left(\sum_{i=1}^{k}\left\|\varphi_{i}\right\|^{2}\right)^{2} \leq \frac{1}{n}\left(\sum_{i=1}^{k} \lambda_{i}\right) \int_{\Omega} \int_{S^{n-1}} \sum_{i=1}^{k} v_{i a}^{2} d S_{a} d x
$$

Finally, from (5.3) we get

$$
\left(\lambda_{k+1}-\lambda_{k}\right) \sum_{i=1}^{k} \int_{S^{n-1}} \int_{\Omega} v_{i a}^{2} d x d S_{a} \leq \sum_{i=1}^{k}\left\|\varphi_{i}\right\|^{2}
$$

Hence

$$
\frac{1}{4}\left(\sum_{i=1}^{k}\left\|\varphi_{i}\right\|^{2}\right)^{2} \leq \frac{1}{n}\left(\sum_{i=1}^{k} \lambda_{i}\right) \frac{\sum_{i=1}^{k}\left\|\varphi_{i}\right\|^{2}}{\lambda_{k+1}-\lambda_{k}}
$$

which implies (1.3) and completes the proof.
Remark 5.1. In the case $\mu$ is Lebesgue measure, the inequality in Theorem 1.2 reduces to

$$
\lambda_{k+1}-\lambda_{k} \leq \frac{4 \sum_{i=1}^{k} \lambda_{i}}{n k}
$$

which coincides with the classical Payne, Pólya and Weinberger inequality (see [15, 18]).

## 6. Comments and open problems

In view of Theorems 1.1 and 1.2, it is of interest to estimate the bound of the norm $\left\|\varphi_{i}\right\|$ of the eigenfunctions $\varphi_{i}$ that satisfy $\left\|\varphi_{i}\right\|_{\mu}=1$.

An upper estimate for the sum of eigenvalues was obtained by Kröger [16]. Let dist $(x, \partial \Omega)$ denote the distance from a point $x \in \Omega$ to the boundary of $\Omega$. Let $\Omega_{r}=\{x \in \Omega$ : $\operatorname{dist}(x, \partial \Omega)<1 / r\}$ and $B_{1}$ be a unit ball in $\mathbb{R}^{n}$. Kröger proved that if there exists a
constant $C_{\Omega}^{(0)}$ such that $\operatorname{vol}\left(\Omega_{r}\right) \leq\left(C_{\Omega}^{(0)} / r\right) \operatorname{vol}(\Omega)^{(n-2) n}$ for every $r>\operatorname{vol}(\Omega)^{-1 / n}$, then for every $k \geq\left(C_{\Omega}^{(0)}\right)^{n}$,

$$
\begin{equation*}
\sum_{j=1}^{k} \lambda_{j} \leq(2 \pi)^{2} \frac{n}{n+2}\left(\operatorname{vol}(\Omega) \operatorname{vol}\left(B_{1}\right)\right)^{-2 / n}\left(k^{(n+2) / n}+C_{n}^{(1)} C_{\Omega}^{(1)} k^{(n+1) / n}\right) \tag{6.1}
\end{equation*}
$$

where $C_{n}^{(1)}$ is a constant depending only on the dimension $n$. It is of interest to generalize this result to Laplacians with respect to measures.

The spectral asymptotics of Laplacians defined on domains by fractal measures have been investigated and obtained by a number of authors (see $[6,7,11-14]$ and the references therein). It is of interest to find examples among these measure for which spectral gap exist.

Acknowledgements. Part of this work was carried out while the first author was visiting the Hunan Normal University, and the second author was visiting Xiangtan University and Harvard University. They are very grateful to the hosting institutions for their hospitality. The second authors thanks Professor Shing-Tung Yau for the possibility to visit CMSA of Harvard University and for providing some valuable comments.

## References

[1] E. J. Bird, S.-M. Ngai and A. Teplyaev, Fractal Laplacians on the unit interval, Ann. Sci. Math. Québec 27 (2003), 135-168.
[2] I. Chavel, Eigenvalues in Riemannian geometry, Academic Press, Inc, Orlando, FL, 1984.
[3] E. B. Davies, Spectral theory and differential operators, Cambridge Studies in Advance Mathematics, vol. 42, Cambridge University Press, Cambridge 1995.
[4] D.-W. Deng and S.-M. Ngai, Eigenvalue estimates for Laplacians on measure spaces, J. Funct. Anal. 268 (2015), 2231-2260.
[5] S. Drenning and R. S. Strichartz, Spectral decimation on Hambly's homogeneous hierarchical gaskets, Illinois J. Math. 53 (2009), 915-937.
[6] U. Freiberg, Spectral asymptotics of generalized measure geometric Laplacians on Cantor like sets, Forum Math. 17 (2005), 87-104.
[7] T. Fujita, A fractional dimension, self-similarity and a generalized diffusion operator, Probabilistic methods in mathematical physics (Katata/Kyoto, 1985), 83-90, Academic Press, Boston, MA, 1987.
[8] K.-E. Hare and D. Zhou, Gaps in the ratios of the spectra of Laplacians on fractals, Fractals 17 (2009), 523-535.
[9] J. Hu, K.-S. Lau and S.-M. Ngai, Laplace operators related to self-similar measures on $\mathbb{R}^{d}$, J. Funct. Anal. 239 (2006), 542-565.
[10] P. Li and S.-T. Yau, On the Schrödinger equation and the eigenvalue problem, Comm. Math. Phys. 88 (1983), 309-318.
[11] H. P. McKean and D. B. Ray, Spectral distribution of a differential operator, Duke Math. J. 29 (1962) 281-292.
[12] K. Naimark and M. Solomyak, The eigenvalue behaviour for the boundary value problems related to self-similar measures on $\mathbb{R}^{d}$, Math. Res. Lett. 2 (1995), 279-298.
[13] S.-M. Ngai, Spectral asymptotics of Laplacians associated with one-dimensional iterated function systems with overlaps, Canad. J. Math. 63 (2011), 648-688.
[14] S.-M. Ngai, W. Tang, and Y. Xie, Spectral asymptotics of one-dimensional fractal Laplacians in the absence of second-order identities, Discrete Contin. Dyn. Syst. 38 (2018), 1849-1887.
[15] L. E. Payne, G. Pólya and H. F. Weinberger, On the ratio of consecutive eigenvalues, Stud. Appl. Math. 35 (1956), 289-298.
[16] P. Kröger, Estimates for sums of eigenvalues of the Laplacian, J. Funct. Anal. 126 (1994), 217-227.
[17] J. P. Pinasco and C. Scarola, Eigenvalue bounds and spectral asymptotics for fractal Laplacians, preprint.
[18] R. Schoen and S.-T. Yau, Lectures on Differential Geometry, Conference Proceedings and Lecture Notes in Geometry and Topology, Vol. 1, International Press, Cambridge, 1994.
[19] R. S. Strichartz, Estimates for sums of eigenvalues for domains in homogeneous spaces, J. Funct. Anal. 137 (1996), 152-190.
[20] R. S. Strichartz, Laplacians on fractals with spectral gaps have nicer Fourier series, Math. Res. Lett. 12 (2005), 269-274.
[21] D. Zhou, Spectral analysis of Laplacians on the Vicsek set, Pacific J. Math. 241 (2009), 369-398.

Hunan Key Laboratory for Computational and Simulation in Science and Engineering, School of Mathematics and Computational Sciences, Xiangtan University, Hunan 411105, People's Republic of China.

E-mail address: taimantang@gmail.com, tmtang@xtu.edu.cn

Key Laboratory of High Performance Computing and Stochastic Information Processing (HPCSIP) (Ministry of Education of China), College of Mathematics and Computer Science, Hunan Normal University, Changsha, Hunan 410081, P. R. China, and Department of Mathematical Sciences, Georgia Southern University, Statesboro, GA 30460-8093, USA.

E-mail address: smngai@georgiasouthern.edu


[^0]:    Date: January 17, 2019.
    2010 Mathematics Subject Classification. Primary: 35P15, 28A80, 35J05. Secondary: 34L16, 65L15, 65L60.

    Key words and phrases. Eigenvalue estimate; Fractal; measure; Laplacian.
    This work is supported in part by the Hunan Provincial Natural Science Foundation of China 12JJ60007 and the National Natural Science Foundation of China, grants 11771136 and 11271122. The second author was also supported by the Center of Mathematical Sciences and Applications (CMSA) of Harvard University, the Hunan Province Hundred Talents Program, Construct Program of the Key Discipline in Hunan Province, and a Faculty Research Scholarly Pursuit Award from Georgia Southern University.

