

Spectral asymptotics of a strong δ' interaction supported by a surfacePavel Exner^{a,b,*}, Michal Jex^{a,c}^a Doppler Institute for Mathematical Physics and Applied Mathematics, Czech Technical University in Prague, Břehová 7, 11519 Prague, Czech Republic^b Nuclear Physics Institute ASCR, 25068 Řež near Prague, Czech Republic^c Department of Physics, Faculty of Nuclear Sciences and Physical Engineering, Czech Technical University in Prague, Břehová 7, 11519 Prague, Czech Republic

ARTICLE INFO

Article history:

Received 25 February 2014

Received in revised form 9 June 2014

Accepted 10 June 2014

Available online 12 June 2014

Communicated by P.R. Holland

Keywords:

 δ' surface interaction

Strong coupling expansion

ABSTRACT

We derive asymptotic expansion for the spectrum of Hamiltonians with a strong attractive δ' interaction supported by a smooth surface in \mathbb{R}^3 , either infinite and asymptotically planar, or compact and closed. Its second term is found to be determined by a Schrödinger type operator with an effective potential expressed in terms of the interaction support curvatures.

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1. Introduction

Quantum mechanics of particles confined to curves, graphs, tubes, surfaces, layers, and other geometrically nontrivial objects is a rich and inspirational subject. On one hand it is useful physically, in particular, to describe various nanostructures, and at the same time it offers numerous interesting mathematical problems. Models of “leaky” structures [1] in which the confinement is realized by an attractive potential have the advantage that they take quantum tunneling into account. The potential is often taken singular, of the δ type, because it is easier to handle [2].

Very recently also more singular couplings of the δ' type attracted attention. The corresponding Hamiltonians can be formally written as

$$H_\beta = -\Delta - \beta^{-1}\delta'(\cdot - \Gamma), \quad (1)$$

where Γ is a smooth surface supporting the interaction. Some prefer to write the interaction term as $\beta^{-1}(\delta'(\cdot - \Gamma), \cdot)\delta'(\cdot - \Gamma)$ to stress that the interaction is invariant with respect to mirror reflection. What is important, however, is that either of the expressions is purely formal. A proper definition which employs the standard δ' concept [3] will be given below, here we only note that we write β^{-1} to underline that a strong δ' interaction corresponds to small values of the parameter β . We also note that investigation of such δ' interactions is not just a mathematical exercise. Due to a seminal idea of Cheon and Shigehara [4] made rigorous in [5,6] they

can be approximated by a scaled “triple-layer” potential combination. The possibility of forming such systems with barriers which become more opaque as the energy increases is no doubt physically attractive.

The subject of this letter is the strong coupling asymptotics of bound states of operators (1) with an attractive δ' interaction supported by a finite or infinite surface in \mathbb{R}^3 . The analogous problem for δ interaction supported by infinite surface was solved in [7]. As in this case, we are going to show that the asymptotics is determined by the geometry of Γ . As a byproduct, we will demonstrate the existence of bound states for sufficiently small β for nonplanar infinite surfaces which are asymptotically planar, in a way alternative to the argument proposed recently in [8].

2. The Hamiltonian

The first thing to do is to define properly the operator (1). It acts, of course, as Laplacian outside of the surface Γ

$$(H_\beta\psi)(x) = -(\Delta\psi)(x)$$

for $x \in \mathbb{R}^3 \setminus \Gamma$ and the interaction will be expressed through suitable boundary conditions on the surface which, in accord with [3], would include continuity of the normal derivative together with a jump of the function value. Specifically, the domain of the operator will be

$$\mathcal{D}(H_\beta) = \left\{ \psi \in H^2(\mathbb{R}^3 \setminus \Gamma) \mid \partial_{n_\Gamma} \psi(x) = \partial_{-n_\Gamma} \psi(x) =: \psi'(x)|_\Gamma, \right. \\ \left. -\beta \psi'(x)|_\Gamma = \psi(x)|_{\partial_+ \Gamma} - \psi(x)|_{\partial_- \Gamma} \right\},$$

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where n_Γ is the normal to Γ and $\psi(x)|_{\partial_\pm \Gamma}$ are the appropriate traces of the function ψ ; all these quantities exist in view of the Sobolev embedding theorem. Being interested in the attractive δ' interactions, we choose the above form of boundary conditions with $\beta > 0$. Another way to define the operator H_β is by the means of the associated quadratic form as discussed in [2]. Its domain is $H^1(\mathbb{R}^3 \setminus \Gamma)$ and the form value for a function $\psi \in H^1(\mathbb{R}^3 \setminus \Gamma)$ is given by

$$h_\beta[\psi] = \|\nabla\psi\|^2 - \beta^{-1} \|\psi(x)|_{\partial_+ \Gamma} - \psi(x)|_{\partial_- \Gamma}\|_{L^2(\Gamma)}^2. \tag{2}$$

As indicated we are interested in the spectrum of H_β in the strong-coupling regime, $\beta \rightarrow 0_+$, for two kinds of surfaces Γ . The first is an infinite surface of which we assume that:

- (a1) Γ is C^4 smooth and allows a global normal parametrization with uniformly bounded elliptic tensor,
- (a2) Γ has no “near self-intersections”, i.e. there exists its symmetric layer neighborhood of a finite thickness which does not intersect with itself,
- (a3) Γ is asymptotically planar in the sense that its curvatures vanish as the geodetic distance from a fixed point tends to infinity,

and finally

- (a4) trivial case is excluded, Γ is not a plane.

In fact, the assumption (a1) can be weakened in a way similar to [9], however, for the sake of simplicity we stick to the existence of a global normal parametrization. The second class to consider are finite surfaces. The compactness makes the assumptions simpler in this case, on the other hand, we have to require additionally absence of a boundary:

- (b) Γ is a closed C^4 smooth surface of a finite genus.

In this case no global parametrization exists, of course, but the geometry of Γ can be described by an atlas of maps representing normal parameterizations with a uniformly bounded elliptic tensor.

3. Geometric preliminaries

Let us collect now some needed facts about the geometry of the surface and its neighborhoods; for a more complete information we refer, e.g., to [10]. We consider infinite surfaces first and we introduce normal coordinates on Γ starting from a local exponential map $\gamma : T_o\Gamma \rightarrow U_o$ with the origin $o \in \Gamma$ to the neighborhood U_o of the point o ; the coordinates s are given by

$$s = (s_1, s_2) \rightarrow \exp_o \left(\sum_i s_i e_i(o) \right) \tag{3}$$

where $\{e_1(o), e_2(o)\}$ is an orthonormal basis of $T_o\Gamma$. By assumption (a1) one can find a point $o \in \Gamma$ such that the map (3) can be extended to a global diffeomorphism from $T_o\Gamma \simeq \mathbb{R}^2$ to Γ .

Using these coordinates, we express components of the surface metric tensor $g_{\mu\nu}$ as $g_{\mu\nu} = \gamma_{,\mu} \cdot \gamma_{,\nu}$ and denote $g^{\mu\nu} = (g_{\mu\nu})^{-1}$. The invariant surface element is denoted as $d\Gamma = g^{\frac{1}{2}} d^2s$ where $g := \det g_{\mu\nu}$. The unit normal $n(s)$ is defined as the cross product of the linearly independent tangent vectors $\gamma_{,\mu}$, i.e. $n(s) = \frac{\gamma_{,1} \times \gamma_{,2}}{|\gamma_{,1} \times \gamma_{,2}|}$. The Gauss curvature K and mean curvature M can be calculated by means of the Weingarten tensor $h_\mu^\nu := -n_{,\mu} \cdot \gamma_{,\sigma} g^{\sigma\nu}$,

$$K = \det h_\mu^\nu = k_1 k_2, \quad M = \frac{1}{2} \text{Tr} h_\mu^\nu = \frac{1}{2} (k_1 + k_2).$$

We recall that the eigenvalues of h_μ^ν are the principal curvatures $k_{1,2}$ and that the identity $K - M^2 = -\frac{1}{4}(k_1 - k_2)^2$ holds.

We also need neighborhoods of the surface Γ . A layer Ω_d of halfwidth $d > 0$ will be defined as the image of $D_d := \{(s, u) : s \in \mathbb{R}^2, u \in (-d, d)\}$ by the map

$$\mathcal{L}: D_d \ni q \equiv (s, u) \rightarrow \gamma(s) + un(s) \tag{4}$$

This definition provides at the same time a parametrization of Ω_d , and the assumption (a2) can be rephrased as

- (a2) there is a $d_0 > 0$ such that the map (4) is injective for any $d < d_0$.

Moreover, in view of (a1) such an \mathcal{L} is a diffeomorphism, which will be crucial for the considerations to follow. The layer Ω_d can be regarded as a manifold with a boundary and characterized by the metric tensor which can be expressed in the parametrization (4) as

$$G_{ij} = \begin{pmatrix} (G_{\mu\nu}) & 0 \\ 0 & 1 \end{pmatrix},$$

where $G_{\mu\nu} = (\delta_\mu^\sigma - uh_\mu^\sigma)(\delta_\rho^\sigma - uh_\rho^\sigma)g_{\rho\nu}$. We use here the convention in which the Latin indices run through 1, 2, 3, numbering the coordinates (s_1, s_2, u) in Ω_d , and the Greek ones through 1, 2. The volume element of the manifold Ω_d can be written in the form $d\Omega_d := \sqrt{G} d^2s du$ with

$$G := \det G_{ij} = g[(1 - uk_1)(1 - uk_2)]^2 = g(1 - 2Mu + Ku^2)^2;$$

with the future purpose in mind we introduce a shorthand for the last factor, $\xi(s, u) := 1 - 2M(s)u + K(s)u^2$. The curvatures also allow us to express more explicitly the next assumption:

- (a3) $K, M \rightarrow 0$ as $|s| := \sqrt{s_1^2 + s_2^2} \rightarrow \infty$.

Recall next a few useful estimates made possible by the assumption (a3), cf. [11]. In combination with (a1) and (a2) it implies that the principal curvatures k_1 and k_2 are uniformly bounded. We set

$$\rho := (\max\{\|k_1\|_\infty, \|k_2\|_\infty\})^{-1};$$

note that $\rho > d_0$ holds for the critical halfwidth of assumption (a2). It can be checked easily that for a given $d < \rho$ the following inequalities are satisfied in the layer neighborhood Ω_d of Γ ,

$$C_-(d) \leq \xi \leq C_+(d), \tag{5}$$

where $C_\pm := (1 \pm d\rho^{-1})^2$, and this in turn implies

$$C_-(d)g_{\mu\nu} \leq G_{\mu\nu} \leq C_+(d)g_{\mu\nu}. \tag{6}$$

Since the metric tensor $g_{\mu\nu}$ uniformly elliptic by assumption, we also have

$$c_- \delta_{\mu\nu} \leq g_{\mu\nu} \leq c_+(d) \delta_{\mu\nu} \tag{7}$$

as a matrix inequality for some positive constants c_\pm .

Let us briefly describe modifications needed if we pass to closed surfaces. As we have indicated a global parametrization is replaced now by a finite atlas \mathcal{A} of maps; in each part \mathcal{M}_i we introduce normal coordinates and define layer neighborhoods by the maps $\hat{\mathcal{M}}_i$ on $D_{i,d} := \{(s, u) : s \in \text{dom}\mathcal{M}_i, u \in (-d, d)\}$ with a given $d > 0$,

$$\hat{\mathcal{M}}_i: D_{i,d} \ni q \equiv (s, u) \rightarrow \gamma_i(s) + un(s) \tag{8}$$

In view of assumption (b) there is a critical $d_0 > 0$ such that every map $\hat{\mathcal{M}}_i: D_{i,d} \rightarrow \Omega_d$ from \mathcal{A} is injective provided $d < d_0$ and

a diffeomorphism. Furthermore, $\hat{\mathcal{M}}_i(s_i, u_i) = \hat{\mathcal{M}}_j(s_j, u_j)$ implies $\mathcal{M}_i(s_i) = \mathcal{M}_j(s_j)$. The above estimates of the metric tensor remains valid also for compact Γ .

4. The results

As in the case of a δ interaction supported by a surface, the asymptotics is determined by the geometry of Γ . To state the results, we introduce the following comparison operator,

$$S = -\Delta_\Gamma - \frac{1}{4}(k_1 - k_2)^2 = -\Delta_\Gamma + K - M^2, \tag{9}$$

where Δ_Γ is the Laplace–Bertrami operator on the surface Γ and $k_{1,2}$ are the principal curvatures of Γ . The spectrum of S is purely discrete if Γ is compact. In the noncompact case the potential vanishes at infinity and has negative values unless Γ is a plane which is, however, excluded by assumption (a4). Consequently, $\sigma_{\text{ess}}(S) = [0, \infty)$ and the discrete spectrum is nonempty. We denote the eigenvalues of S , arranged in the ascending order with the multiplicity taken into account, as μ_j .

First we inspect the essential spectrum in the strong-coupling regime:

Theorem 1. *Let an infinite surface Γ satisfy assumptions (a1)–(a4), then $\sigma_{\text{ess}}(H_\beta) \subseteq [\epsilon(\beta), \infty)$, where $\epsilon(\beta) = -\frac{4}{\beta^2} + \mathcal{O}(e^{-c/\beta})$ holds as $\beta \rightarrow 0_+$ for some constant $c > 0$.*

We note that in case of a compact Γ we have $\sigma_{\text{ess}}(H_\beta) = [0, \infty)$; a proof can be found in [8]. The next two theorems describe the asymptotics of the negative point spectrum of H_β .

Theorem 2. *Let an infinite surface Γ satisfy assumptions (a1)–(a4), then H_β has at least one isolated eigenvalue below the threshold of the essential spectrum for all sufficiently small $\beta > 0$, and the j -th eigenvalue behaves in the limit $\beta \rightarrow 0_+$ as*

$$\lambda_j = -\frac{4}{\beta^2} + \mu_j + \mathcal{O}(-\beta \ln \beta).$$

Theorem 3. *Let a compact surface Γ satisfy assumption (b), then H_β has at least one isolated eigenvalue below the threshold of the essential spectrum for all $\beta > 0$, and the j -th eigenvalue behaves in the limit $\beta \rightarrow 0_+$ as*

$$\lambda_j = -\frac{4}{\beta^2} + \mu_j + \mathcal{O}(-\beta \ln \beta).$$

5. Bracketing estimates

The basic idea is analogous to the one used in [7], namely to estimate the operator H_β from above and below, in a tight enough manner, by suitable operators for which we are able to calculate the spectrum directly. The starting point for such estimates is the bracketing trick, that is, imposing additional Dirichlet/Neumann conditions at the boundary of the neighborhood Ω_d of the surface Γ . We introduce quadratic forms h_β^+ and h_β^- , both of them given by the formula

$$\|\nabla\psi\|_{L^2(\Omega_d)}^2 - \beta^{-1} \int_\Gamma |\psi(s, 0_+) - \psi(s, 0_-)|^2 d\Gamma$$

with the domains $\mathcal{D}(h_\beta^+) = \tilde{H}_0^1(\Omega_d \setminus \Gamma)$ and $\mathcal{D}(h_\beta^-) = H^1(\Omega_d \setminus \Gamma)$, respectively, the former being understood as a set of functions which are locally H^1 and vanish at the boundary of Ω_d . We denote the self-adjointed operators associated with these forms as H_β^\pm . By the standard bracketing argument we get

$$-\Delta_{\mathbb{R}^3 \setminus \Omega_d}^N \oplus H_\beta^- \leq H_\beta \leq -\Delta_{\mathbb{R}^3 \setminus \Omega_d}^D \oplus H_\beta^+, \tag{10}$$

where $-\Delta_{\mathbb{R}^3 \setminus \Omega_d}^{D,N}$ is the Dirichlet–Laplacian and Neumann–Laplacian respectively on the set $\mathbb{R}^3 \setminus \Omega_d$. The operators $-\Delta_{\mathbb{R}^3 \setminus \Omega_d}^{D,N}$ are positive, hence all the information about the negative spectrum is encoded in the operators H_β^\pm .

The next step is to transform the operators H_β^\pm into the curvilinear coordinates (s, u) . This is done by means of the unitary transformation

$$U\psi = \psi \circ \mathcal{L} : L^2(\Omega_d) \rightarrow L^2(D_d, d\Omega).$$

By $(\cdot, \cdot)_G$ we denote the scalar product in $L^2(D_d, d\Omega)$. The operators $UH_\beta^\pm U^{-1}$ acting on this space are associated with the forms

$$\begin{aligned} &h_\beta^\pm(U^{-1}\psi) \\ &= (\partial_i\psi, G^{ij}\partial_j\psi)_G - \beta^{-1} \int_\Gamma |\psi(s, 0_+) - \psi(s, 0_-)|^2 d\Gamma \end{aligned}$$

having the domains $\tilde{H}_0^1(D_d \setminus \Gamma, d\Omega)$ and $\tilde{H}^1(D_d \setminus \Gamma, d\Omega)$, respectively. Next we employ another unitary transformation, inspired by [11], with the aim to get rid of the transverse coordinate dependence, i.e. switch from the metric $d\Omega$ to $d\Gamma du$ by

$$\tilde{U}\psi = \xi^{\frac{1}{2}}\psi : L^2(D_d, d\Omega) \rightarrow L^2(D_d, d\Gamma du).$$

Similarly as before, we denote the scalar product in $L^2(D_d, d\Gamma du)$ as $(\cdot, \cdot)_g$ and consider the operators

$$F_\beta^\pm := \tilde{U}UH_\beta^\pm U^{-1}\tilde{U}^{-1}$$

which act in $L^2(D_d, d\Gamma du)$. The quadratic forms ζ_β^\pm associated with F_β^\pm can be calculated as $h_\beta^\pm(\tilde{U}^{-1}U^{-1}\psi)$ with the result

$$\begin{aligned} \zeta_\beta^+[\psi] &= (\partial_\mu\psi, G^{\mu\nu}\partial_\nu\psi)_g + (\psi, (V_1 + V_2)\psi)_g + \|\partial_3\psi\|_g \\ &\quad - \beta^{-1} \int_\Gamma |\psi(s, 0_+) - \psi(s, 0_-)|^2 d\Gamma \\ &\quad - \int_\Gamma M(|\psi(s, 0_+)|^2 - |\psi(s, 0_-)|^2) d\Gamma \\ \zeta_\beta^-[\psi] &= \zeta_\beta^+[\psi] + \int_\Gamma \varsigma(s, d)|\psi(s, d)|^2 d\Gamma \\ &\quad - \int_\Gamma \varsigma(s, -d)|\psi(s, -d)|^2 d\Gamma, \end{aligned}$$

where $\varsigma = \frac{M-Ku}{\xi}$, the two curvature-induced potentials are

$$V_1 = g^{-\frac{1}{2}}(g^{\frac{1}{2}}G^{\mu\nu}J_{,\mu})_{, \nu} + J_{,\mu}G^{\mu\nu}J_{,\nu}, \quad V_2 = \frac{K - M^2}{\xi^2}$$

with $J = \frac{\ln \xi}{2}$. The corresponding form domains are $\tilde{H}_0^1(D_d \setminus \Gamma, d\Gamma du)$ and $\tilde{H}^1(D_d \setminus \Gamma, d\Gamma du)$, respectively.

6. Proof of Theorem 1

In the excluded case when Γ is a plane, the spectrum is easily found by separation of variables. Since a δ' interaction in one dimension has a single eigenvalue equal to $-\frac{4}{\beta^2}$, cf. [3, Sec. 1.4], we get $\sigma(H_\beta) = \sigma_{\text{ess}}(H_\beta) = [-\frac{4}{\beta^2}, \infty)$. We want to show that under the assumption (a3) the essential spectrum does not change, at least asymptotically. We employ an estimate which follows from Lemma 4 that we will prove below, namely

$$\int_{-d}^d \left| \frac{df}{du} \right|^2 du - \beta^{-1} |f(0_+) - f(0_-)|^2 \geq \left(-\frac{4}{\beta^2} - \frac{16}{\beta^2} \exp\left(-\frac{4d}{\beta}\right) \right) \|f\|_{L^2(-d,d)}. \tag{11}$$

As we shall see the inequality holds for sufficiently small β and $\frac{d}{\beta} > 2$. The inclusion $\sigma_{\text{ess}}(H_\beta) \subseteq [\epsilon(\beta), \infty)$ is equivalent to

$$\inf \sigma_{\text{ess}}(H_\beta) \geq \epsilon(\beta)$$

which will be satisfied if $\inf \sigma_{\text{ess}}(H_\beta^-) \geq \epsilon(\beta)$ for H_β^- acting in $L^2(\Omega_d)$ for $d < g_0 < \rho$. This is obvious from inequalities (10) and the fact that the operator $-\Delta_{\mathbb{R}^3 \setminus \Omega_d}^N$ is positive and cannot thus contribute to the negative part of the spectrum. In the next step we divide the surface Γ into two parts, namely $\Gamma_\tau^{\text{int}} := \{s \in \Gamma \mid r(s) < \tau\}$ and $\Gamma_\tau^{\text{ext}} := \Gamma \setminus \overline{\Gamma_\tau^{\text{int}}}$. The layer neighborhoods corresponding to Γ_τ^{int} and Γ_τ^{ext} are $D_\tau^{\text{int}} = \{(s, u) \in D_d \mid s \in \Gamma_\tau^{\text{int}}\}$ and $D_\tau^{\text{ext}} = D_d \setminus \overline{D_\tau^{\text{int}}}$. We introduce the Neumann operators on respective neighborhoods, $H_{\beta,\tau}^{-,z}$ for $z = \text{int}, \text{ext}$ associated with the forms

$$(\partial_i \psi, G^{ij} \partial_j \psi)_G - \beta^{-1} \int_{\Gamma_\tau^z} |\psi(s, 0_+) - \psi(s, 0_-)|^2 d\Gamma$$

defined on $\tilde{H}^1(D_\tau^z \setminus \Gamma, d\Omega)$. Using once more Neumann bracketing we get $H_\beta^- \geq H_{\beta,\tau}^{-,\text{int}} \oplus H_{\beta,\tau}^{-,\text{ext}}$. The inner part is compact, hence the spectrum of $H_{\beta,\tau}^{-,\text{int}}$ is purely discrete. Consequently, the min-max principle implies

$$\inf \sigma_{\text{ess}}(H_\beta^-) \geq \inf \sigma_{\text{ess}}(H_{\beta,\tau}^{-,\text{ext}}),$$

and it is sufficient to check that the right-hand side cannot be smaller than $\epsilon(\beta)$. The quantities $m_\tau^+ := \sup_{\Gamma_\tau^{\text{ext}}} \xi$ and $m_\tau^- := \inf_{\Gamma_\tau^{\text{ext}}} \xi$ tend to one as $\tau \rightarrow \infty$ in view of assumption (a3). We have the following estimate,

$$\begin{aligned} (\psi, H_{\beta,\tau}^{-,\text{ext}} \psi)_G &\geq \int_{D_\tau^{\text{ext}}} |\partial_3 \psi(q)|^2 d\Omega \\ &\quad - \beta^{-1} \int_{\Gamma_\tau^{\text{ext}}} |\psi(s, 0_+) - \psi(s, 0_-)|^2 d\Gamma \\ &\geq m_\tau^- \int_{D_\tau^{\text{ext}}} |\partial_3 \psi(q)|^2 d\Gamma du \\ &\quad - \beta^{-1} \int_{\Gamma_\tau^{\text{ext}}} |\psi(s, 0_+) - \psi(s, 0_-)|^2 d\Gamma \\ &\geq \frac{1}{\beta^2 m_\tau^+ m_\tau^-} \left[-4 - 16 \exp\left(-\frac{4d}{\beta}\right) \right] \\ &\quad \times \int_{D_\tau^{\text{ext}}} |\psi(q)|^2 d\Omega, \end{aligned}$$

and since τ is arbitrary, we obtain $\epsilon(\beta) \geq -\frac{4}{\beta^2} - \frac{16}{\beta^2} \exp(-\frac{4d}{\beta})$.

7. Proof of Theorem 2

To prove the second theorem, we will need several auxiliary results. The operators F_β^\pm are still not suitable to work with and so we replace them with a slightly cruder bounds. First we estimate the values of the potentials V_1 and V_2 . With the help of inequalities (5)–(7) we are able to check that

$$dv^- \leq V_1 \leq dv^+$$

holds for suitable numbers v^\pm and $d < d_0 < \rho$. On the other hand, V_2 can be estimated as

$$C_-^{-2}(K - M^2) \leq V_2 \leq C_+^{-2}(K - M^2),$$

where C_\pm are the same as in (5). This allows us to replace (10) with the estimates using operators D_β^\pm ,

$$\begin{aligned} D_{d,\beta}^- &:= U_d^- \otimes I + \int_{\Gamma}^{\oplus} T_{d,\beta}^-(s) d\Gamma \leq F_\beta^- \leq H_\beta \\ H_\beta &\leq F_\beta^+ \leq U_d^+ \otimes I + \int_{\Gamma}^{\oplus} T_{d,\beta}^+(s) d\Gamma =: D_{d,\beta}^+ \end{aligned} \tag{12}$$

where

$$U_d^\pm = -C_\pm \Delta_\Gamma + C_\pm^{-2}(K - M^2) + v^\pm d$$

with the domain $\mathcal{D}(U_d^\pm) = L^2(\mathbb{R}^2, d\Gamma)$ and the transverse part acts as

$$T_{d,\beta}^\pm(s) \psi = -\Delta \psi$$

with the domains

$$\begin{aligned} \mathcal{D}(T_{a,\beta}^+(s)) &= \{f \in H^2(((-a, a) \setminus \{0\}) \mid f(a) = f(-a) = 0, \\ &\quad f'(0_-) = f'(0_+) = -\beta^{-1}(f(0_+) - f(0_-)) \\ &\quad + M(f(0_+) + f(0_-))\} \end{aligned}$$

and

$$\begin{aligned} \mathcal{D}(T_{a,\beta}^-(s)) &= \left\{ f \in H^2(((-a, a) \setminus \{0\}) \mid \mp \frac{\|M\|_\infty + d\|K\|_\infty}{C_-} f(\pm a) \right. \\ &= f'(\pm a), \\ &\quad \left. f'(0_-) = f'(0_+) = -\beta^{-1}(f(0_+) - f(0_-)) \right. \\ &\quad \left. + M(f(0_+) + f(0_-)) \right\}, \end{aligned}$$

respectively. The negative spectrum is described by the following result the proof of which can be found in [12].

Lemma 4. Each of the operators $T_{d,\beta}^\pm(s)$ has exactly one negative eigenvalue $t_\pm(d, \beta)$, respectively, which is independent of s provided that $\frac{d}{\beta} > 2$ and $\beta(\|M\|_\infty + d\|K\|_\infty) < 1$. For all $\beta > 0$ sufficiently small these eigenvalues satisfy the inequalities

$$\begin{aligned} -\frac{4}{\beta^2} - \frac{16}{\beta^2} \exp\left(-\frac{4d}{\beta}\right) &\leq t_-(d, \beta) \leq -\frac{4}{\beta^2} \leq t_+(d, \beta) \\ &\leq -\frac{4}{\beta^2} + \frac{16}{\beta^2} \exp\left(-\frac{4d}{\beta}\right). \end{aligned}$$

On the other hand, the spectrum of the operators U_d^\pm has the asymptotic expansion governed by the operator S which we can adopt from [7]:

Lemma 5. The eigenvalues of U_d^\pm satisfy the relations

$$\mu_j^\pm(d) = \mu_j + C_j^\pm d + \mathcal{O}(d^2) \quad \text{for } d \rightarrow 0,$$

where μ_j is the j -th eigenvalue of the operator S and the constants C_j^\pm are independent on d .

With these prerequisites we are ready to prove the second theorem. We put $d(\beta) = -\beta \ln \beta$. Using the fact that each of the operators $T_{d,\beta}^{\pm}(s)$ has exactly one negative eigenvalue $t_{\pm}(d(\beta), \beta)$ together with the explicit form of $D_{d,\beta}^{\pm}$ we can write their spectra as $t_{\pm}(d(\beta), \beta) + \mu_j^{\pm}(d(\beta))$, where μ_j^{\pm} are the eigenvalues of the operators U_d^{\pm} . Using now Lemmata 4 and 5 we are able to rewrite this as

$$t_{\pm}(d(\beta), \beta) + \mu_j^{\pm}(d(\beta)) = -\frac{4}{\beta^2} + \mu_j + \mathcal{O}(\beta |\ln \beta|),$$

hence the min–max principle in combination with inequalities (12) conclude the argument.

8. Proof of Theorem 3

The existence of isolated eigenvalues can be checked variationally as in [8]. For a test function ξ one chooses characteristic function of the volume enclosed by the surface Γ ; this yields an estimate of the ground state energy from above,

$$\lambda_0 \leq \frac{h_{\beta}(\xi)}{\|\xi\|^2} = \beta^{-1} \frac{S}{V} \quad (13)$$

where h_{β} is quadratic (2), S is the area of the surface Γ and V is the volume enclosed by Γ . The proof of the asymptotic expansion proceed with minimum modifications as for the infinite surface, hence we omit the details.

Acknowledgements

The research was supported by the Czech Science Foundation

within the project 14-06818S and by Grant Agency of the Czech Technical University in Prague, grant No. SGS13/217/OHK4/3T/14.

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