QUANTUM TRACE MAP FOR 3-MANIFOLDS AND A 'LENGTH CONJECTURE'

PRARIT AGARWAL, DONGMIN GANG, SANGMIN LEE, AND MAURICIO ROMO

ABSTRACT. We introduce a quantum trace map for an ideally triangulated hyperbolic knot complement $S^3 \setminus \mathcal{K}$. The map assigns a quantum operator to each element of Kauffmann Skein module of the 3-manifold. The quantum operator lives in a module generated by products of quantized edge parameters of the ideal triangulation modulo some equivalence relations determined by gluing equations. Combining the quantum map with a state-integral model of $SL(2, \mathbb{C})$ Chern-Simons theory, one can define perturbative invariants of knot K in the knot complement whose leading part is determined by its complex hyperbolic length. We then conjecture that the perturbative invariants determine an asymptotic expansion of the Jones polynomial for a link composed of \mathcal{K} and K. We propose the explicit quantum trace map for figure-eight knot complement and confirm the *length conjecture* up to the second order in the asymptotic expansion both numerically and analytically.

Contents

1. Introductio	n	2
2. Quantum t	race map	3
2.1. Kauffman	n Skein modules	3
2.2. Quantum	gluing module	6
2.3. Quantum	trace map and Length conjecture	7
3. Perturbativ	ve knot invariants	9
3.1. State-inte	egral model with quantum trace map	9
3.2. Perturbat	tive invariants	11
3.3. Feynman	perturbation theory for the invariants	13
4. Example : figure-eight knot-complement		15
4.1. Explicit of	quantum trace map	16
4.2. Test of th	ne length conjecture - analytic proof up to 1-loop	17
4.3. Test of the length conjecture - numerical check		
5. 3D index with quantum trace map		
6. Quantum trace map in complex Chern-Simons theory		
Acknowledgements		
Appendix A.	Gluing equations of an ideal triangulation	27
Appendix B. Colored Jones polynomial for framed link		28
Appendix C. Quantum dilogarithm function		31
References		

 $\mathbf{2}$

1. INTRODUCTION

Kauffman bracket skein modules (KBSMs) were independently introduced by J. H. Przytycki [27] and V. G. Turaev [32] based on the Kauffman bracket [21] in an attempt to generalize knot polynomials in S^3 to those in arbitrary 3-manifolds. The module becomes a non-commutative algebra when the 3-manifold is chosen to be a thickened surface $S \times I$ with marked points on S. For the case, F. Bonahon and H. Wong in [2] constructed an injective algebra homomorphism, called quantum trace map, from the Skein algebra to the Chechov-Fock algebra of S. The Chechov-Fock algebra is a quantization of Teichmüller space using Thurston's shear coordinates [29, 26] associated to an ideal triangulation of S. In terms of $SL(2,\mathbb{C})$ Chern-Simons theory, the quantum operators correspond to Wilson loop operators whose trajectory is confined on the surface S. Classically the Wilson loops can be represented by a function on the phase space P(S) associated with the surface S. The phase space is the moduli space of $SL(2,\mathbb{C})$ flat connections on S. Shear coordinates provides a natural coordinate of the phase space and the loop operators can be given as a Laurent polynomial of the coordinates. After quantization, the phase space becomes a Hilbert-space on which the quantized loop operators naturally act. Via the 2D quantum trace map, one can express the loop operators in terms of Laurent polynomial of the quantized shear coordinates.

Generalizing the idea of quantum trace map to a general 3-manifold M seems to be pointless since the KBSM is just a module instead of forming an algebra. Contrary to the common belief, we suggest that there is a natural unique 3D quantum trace map. Wilson loop operators in the Chern-Simons theory can be defined along arbitrary links on a 3-manifold. Classically, the loop operator can be regarded as a function on $\mathcal{L}(M)$, the moduli space of $SL(2,\mathbb{C})$ flat connections on M. Unlike P(S), there is no natural symplectic structure on $\mathcal{L}(M)$. Instead, the space $\mathcal{L}(M)$ can be regarded as a Lagrangian subvariety of the phase space $P(\mathbb{T}^2)$ when the 3-manifold has a torus (\mathbb{T}^2) boundary. For an ideally triangulated 3-manifold M, the moduli space can be represented by an algebraic variety determined by gluing equations [30, 25]. Gluing equations are set of algebraic equations among edge parameters of ideal tetrahedra in the triangulation. The gluing equations have a symplectic structure and the quantization of the $\mathcal{L}(M)$ has been well-studied [5, 9]. As a result of the quantization, state-integral models [18, 5] are developed which compute the partition function of $SL(2,\mathbb{C})$ Chern-Simons theory on M. To define a 3D quantum trace map, one needs to quantize functions on $\mathcal{L}(M)$. The functions are given by a Laurent series of the edge parameters subjected to gluing equations. To deal with the 'quantum functions' on $\mathcal{L}(M)$, we introduce a quantum gluing module in Section 2.2 which are the space of them. Our quantum trace map is an injective module homomorphism from the KBSM of M to the quantum gluing module. Once the correct quantum trace map is given, one can generalize the stateintegral models with insertion of the Wilson loop operators along arbitrary knots or links in M.

One big motivation for studying $SL(2,\mathbb{C})$ Chern-Simons theory on a hyperbolic knot complement $S^3 \setminus \mathcal{K}$ is its relation to the volume conjecture [19, 20, 24]. The conjecture relates a large *n* limit of the colored Jones polynomial $J_n(\mathcal{K};q)$ to the hyperbolic volume of its knot complement, $S^3 \setminus \mathcal{K}$. As its strongest version [16], it is conjectured that the asymptotic limit is fully determined by a perturbative expansion of the state-integral on $M = S^3 \setminus \mathcal{K}$. Our 3D quantum trace map can be naturally blended with these developments. We propose that the large n limit of the ratio $J_{n,\tilde{n}=2}(\mathcal{K} \cup K;q)/J_n(\mathcal{K};q)$ can be fully captured by the state-integral model for $S^3 \setminus \mathcal{K}$ with an insertion of quantum trace operator associated with the knot K. We think of \mathcal{K} as a heavy knot with large color n while K as a light knot with fixed color, $\tilde{n} = 2$. The leading order result is given by a simple function of the hyperbolic length of the light knot K in the knot complement $S^3 \setminus \mathcal{K}$, and we call it a *length conjecture* for an obvious reason.

2. QUANTUM TRACE MAP

After a brief review on the trace map of 3-manifold, we introduce a quantum version of the map which we call *quantum trace map*.

2.1. Kauffman Skein modules. We begin by recalling the definitions of Kauffman Skein modules from [28, 4]. Let \mathcal{K} be a hyperbolic knot in S^3 and M be the knot complement: $M = S^3 \setminus \mathcal{K}$. We restrict our discussion to hyperbolic knot complements. Some of the statements below should be modified when M is a general 3-manifold.

We denote the Kauffman bracket Skein module on M by $S_q[M]$ and its 'even' submodule by $S_q^{\text{even}}[M]$. The two modules are defined as

(1)

$$\mathcal{S}_{q}[M] := \frac{\mathbb{C}[q^{\pm 1/4}] \text{-module with basis } \{Y_{K}\}_{(K: \text{ framed link in } M)}}{\langle Y_{K_{+}} - q^{-1/4}Y_{K_{0}} - q^{1/4}Y_{K_{\infty}}, Y_{K \bigsqcup \bigcirc} + (q^{1/2} + q^{-1/2})Y_{K} \rangle}, \\
\mathcal{S}_{q}^{\text{even}}[M] := \frac{\mathbb{C}[q^{\pm 1/4}] \text{-module with basis } \{Y_{K}\}_{(K: \text{ framed 'even' link in } M)}}{\langle Y_{K_{+}} - q^{-1/4}Y_{K_{0}} - q^{1/4}Y_{K_{\infty}}, Y_{K \bigsqcup \bigcirc} + (q^{1/2} + q^{-1/2})Y_{K} \rangle}.$$

The basis includes $Y_{K=\emptyset}$, where \emptyset is the empty link. Here, K_+ , K_0 and K_{∞} , are three framed links which are identical except in a small 3-ball, as depicted in Figure 1. $K \bigsqcup \bigcirc$ is the union of K with an unlinked, 0-framed unknot in the trivial homotopy class. The Y_K in $S_q[M]$ is invariant under Reidemeister moves II, III



FIGURE 1. Kauffman triple

and modified Reidemeister move I on K but not under Reidemeister move I. This



FIGURE 2. Y_K under Reidemeister move I on K.

is why the Y_K is labelled by *framed* link K. We call a link $K = \bigcup_{I=1}^{\sharp(K)} K_I$ consisting of $\sharp(K)$ knots an 'even' link if

$$\prod_{I=1}^{\sharp(K)} \operatorname{sgn}(K_I) = +1 \; ,$$

and an 'odd' link otherwise. Here $sgn(K_I) \in \{+1, -1\}$ is defined as

$$\operatorname{sgn}(K_I) := \begin{cases} +1, & \text{if } [K_I] = 1 \text{ as an element of } H_1(M, \mathbb{Z}_2) = \mathbb{Z}_2 = \{+1, -1\}, \\ -1, & \text{otherwise.} \end{cases}$$

For later use, we consider a submodule $S_q^{\text{even}}[M]$ whose basis Y_K are labelled by even links K on M. Note that the Kauffman bracket Skein relation is well-defined in S_q^{even} , i.e. all three links in Kauffman triple share the same evenness/oddness and $K \bigsqcup \bigcirc$ is even if K is even.

2.1.1. Skein algebra and trace map. In the special case $q^{1/2} = 1$, the Skein module gives rise to the Skein algebra, $V[M] := S_q[M]|_{q^{1/2}=1}$.

FIGURE 3. Kauffman bracket Skein relation at $q^{1/2} = 1$.

At $q^{1/2} = 1$, the Skein relation equates the over-crossing and the under-crossing, as in Figure 3, so that the Skein algebra is a indeed a commutative \mathbb{C} -algebra equipped with the multiplication,

(3)
$$Y_{K_1}Y_{K_2} = Y_{K_2}Y_{K_1} := Y_{K_1 \cup K_2}$$

The \mathbb{C} -algebra is generated by elements in $\pi_1(M)$ modulo some relations:

(4)
$$V[M] = \frac{\mathbb{C}[Y_{\gamma}: \gamma \in \pi_1(M)]}{\langle Y_{\gamma_1}Y_{\gamma_2} - Y_{\gamma_2}Y_{\gamma_1}, Y_{\gamma_1}Y_{\gamma_2} + Y_{\gamma_1\gamma_2} + Y_{\gamma_1\gamma_2^{-1}}, Y_e + 2\rangle},$$
$$V^{\text{even}}[M] = \{x \in V[M] : x \text{ is even}\}.$$

Here $e \in \pi_1(M)$ is the identity element. The above relations imply that $Y_{\gamma} = Y_{\gamma^{-1}}$ and Y_{γ} depends only on the conjugacy class of γ in $\pi_1(M)$. General element $x \in V^{\text{even}}[M]$ is given as

$$\sum_{I} C_{I} \prod_{i=1}^{n_{I}} Y_{\gamma_{i}^{I}} \quad \text{with } \gamma_{i}^{I} \in \pi_{1}(M), \ C_{I} \in \mathbb{C} \text{ and } \prod_{i=1}^{n_{I}} \operatorname{sgn}[\gamma_{i}^{I}] = 1 \text{ for all } I.$$

The relations in the denominator come from the relations in the Skein module $S_q[M]$ in eqn.(1) at $q^{1/4} = -1$. We choose $q^{1/4} = -1$ instead of $q^{1/4} = 1$ for a later convenience. There is an algebra isomorphism between $S_q[M]$ at $q^{1/4} = 1$ and $S_q[M]$ at $q^{1/4} = -1$.

One interesting aspect of the Skein algebra is its relation to the coordinate ring of $SL(2, \mathbb{C})$ character variety X[M], which is defined as

(5)
$$X[M] := \operatorname{Hom}[\pi_1(M) \to SL(2, \mathbb{C})] / \sim$$

where $\rho_1 \sim \rho_2$ if $\operatorname{Tr}(\rho_1(\gamma)) = \operatorname{Tr}(\rho_2(\gamma))$ for all $\gamma \in \pi_1(M)$.

For our purposes, we consider the \mathbb{Z}_2 -quotient of the variety, $X[M]/\mathbb{Z}_2$, where the \mathbb{Z}_2 is the action of $H^1(M, \mathbb{Z}_2) = \text{Hom}(\pi_1 M \to \mathbb{Z}_2 = \{\pm 1\}) = \{1, \eta\}$:

(6)
$$X[M]/\mathbb{Z}_2 = \{\rho \in X[M]\}/ \sim$$

where $\rho_1 \sim \rho_2$ if $\rho_1(\gamma) = \eta(\gamma)\rho_2(\gamma)$ for all $\gamma \in \pi_1(M)$.

One can naturally define an algebra homomorphism Φ_{tr} called *trace map* from the Skein algebra V[M] to the coordinate ring $\mathbb{C}[X(M)]$ (or from $V^{\text{even}}[M]$ to $\mathbb{C}[X(M)/\mathbb{Z}_2]$), the algebra of functions on X(M), as follows

(7)
$$\begin{aligned} \Phi_{\mathrm{tr}} &: V(M) \ (\mathrm{or} \ V^{\mathrm{even}}(M)) \to \mathbb{C}[X(M)] \ (\mathrm{or} \ \mathbb{C}[X(M)/\mathbb{Z}_2]) \,, \text{ where} \\ & [\Phi_{\mathrm{tr}}(Y_{\gamma})](\rho) = -\mathrm{Tr}(\rho(\gamma)) \,. \end{aligned}$$

To see that the map is an algebra homomorphism, one needs to use the following property of $SL(2, \mathbb{C})$ matrices

(8)
$$\operatorname{Tr}(A)\operatorname{Tr}(B) = \operatorname{Tr}(AB) + \operatorname{Tr}(AB^{-1}), \quad \operatorname{Tr}(\mathbb{I}) = 2,$$

where \mathbb{I} is the 2 × 2 identity matrix. The trace map is well-defined as a homomorphism from $V^{\text{even}}(M)$ to $\mathbb{C}[X(M)/\mathbb{Z}_2]$ since $\prod_{I=1}^n \eta(\gamma_I) = 1$ for $\prod_{I=1}^n Y_{\gamma_I} \in V^{\text{even}}(M)$. This section can be summarized by the diagram in Figure 4.

$$\begin{array}{c} \mathcal{S}_q^{\text{even}}[M] \\ & \bigvee \\ q^{1/2} = 1 \\ \Phi_{\text{tr}} : V^{\text{even}}[M] \longrightarrow \mathbb{C}[X(M)/\mathbb{Z}_2] \end{array}$$

FIGURE 4. Relations among $S_q^{\text{even}}[M]$ (even Kauffman bracket Skein module), $V^{\text{even}}[M]$ (even Skein algebra) and $\mathbb{C}[X(M)/\mathbb{Z}_2]$ (coordinate ring of character variety). The relations suggest that $S_q^{\text{even}}[M]$ can be considered as a quantization of $\mathbb{C}[X(M)/\mathbb{Z}_2]$.

For later use, we define the canonical component $X_0[M] \subset X[M]$ for a hyperbolic knot complement $M = S^3 \setminus \mathcal{K}$ as

 $X_0[M] :=$ a connected component of X[M] which contain ρ_{hyp} ,

(9)
$$\rho_{\text{hyp}} := \text{an } SL(2, \mathbb{C})\text{-representation}$$

corresponding to the complete hyperbolic structure .

In the same way, one can define $X_0(M)/\mathbb{Z}_2$.

2.2. Quantum gluing module. Following [5, 12], we define the quantum gluing module $\hat{\mathbb{C}}_{q}[\mathcal{T}]$ associated to an ideal triangulation \mathcal{T} of M as

(10)
$$\hat{\mathbb{C}}_{q}[\mathcal{T}] = \frac{\left(\mathbb{C}[q^{\pm 1/4}][\hat{z}_{i}, \hat{z}_{i}^{\prime\prime}, \hat{z}_{i}^{-1}, \hat{z}_{i}^{\prime\prime-1}]\right) / \langle \hat{z}_{i} \hat{z}_{j}^{\prime\prime} - q^{-\delta_{ij}} \hat{z}_{j}^{\prime\prime} \hat{z}_{i} \rangle)}{\text{equivalence relations} \sim}$$

where $\hat{O} \sim e^{\hat{C}_{I}} \cdot \hat{O}$ and $\hat{O} \sim \hat{O} \cdot (\hat{z}_{i}^{-1} + \hat{z}_{i}^{\prime\prime})$.

Here *i* runs from 1 to $|\mathcal{T}|$, the number of tetrahedra in the ideal triangulation. The precise definition of the exponentiated internal edge operators $e^{\hat{C}_I}$ can be found in (130) in appendix A. Before taking the quotient by the equivalence relations, it is a non-commutative algebra. After the quotient, it becomes just a $\mathbb{C}[q^{\pm 1/4}]$ -module since the non-commutative multiplication is not compatible with the equivalence relation and thus can not be well-defined in $\hat{\mathbb{C}}_q[\mathcal{T}]$.

At $q^{1/2} = 1$, $\hat{\mathbb{C}}_q[\mathcal{T}]$ becomes the algebra of functions on the gluing equation variety $\chi[\mathcal{T}]$

(11)

$$\begin{aligned}
\hat{\mathbb{C}}_{q}[\mathcal{T}]|_{q^{1/2}=1} \simeq \mathbb{C}[\chi[\mathcal{T}]], \\
\chi[\mathcal{T}] := \{z_{i}, z_{i}'' : z_{i}^{-1} + z_{i}'' - 1 = 0, e^{C_{I}} = 1\}_{i=1}^{|\mathcal{T}|}.
\end{aligned}$$

For each element $(z_i, z_i'') \in \chi[\mathcal{T}]$, there is an associated representation [30]

(12)
$$\rho_{(z,z'')}^{\mathcal{T}} \in \operatorname{Hom}[\pi_1(S^3 \backslash \mathcal{K}) \to PSL(2,\mathbb{C})]$$

We choose an ideal triangulation \mathcal{T} which is a ρ_{hyp} -regular:

(13) \mathcal{T} is called ρ_{hyp} -regular if $\exists (z, z'')_{\text{hyp}} \in \chi[\mathcal{T}]$ such that $\rho_{(z, z'')_{\text{hyp}}}^{\mathcal{T}} = \rho_{\text{hyp}}$.

 ρ_{hyp} is an $SL(2, \mathbb{C})$ -representation in (9) but also can be regarded as an $PSL(2, \mathbb{C}) = SL(2, \mathbb{C})/\mathbb{Z}_2$ -representation. Let us denote by $\chi_0[\mathcal{T}]$ the connected component of $\chi[\mathcal{T}]$ which contains $(z, z'')_{\text{hyp}}$. In general, the gluing equation variety $\chi[\mathcal{T}]$ depends on the choice of ρ_{hyp} -regular triangulation \mathcal{T} while the $\chi_0[\mathcal{T}]$ does not [31]. The component can be identified with

(14)
$$\chi_0[\mathcal{T}] \simeq X_0(M)/\mathbb{Z}_2$$

Another nice reason for restricting to the component $\chi_0[\mathcal{T}]$ is that any $PSL(2, \mathbb{C}) = SL(2, \mathbb{C})/\mathbb{Z}_2$ representation $\rho_{(z,z'')}^{\mathcal{T}}$ with $(z, z'') \in \chi_0[\mathcal{T}]$ can be lifted to a $SL(2, \mathbb{C})$ representation. Thus, one can define the classical trace map $\Phi_{\rm tr} : V^{\rm even}[M] \to \mathbb{C}[\chi_0[\mathcal{T}]]$ as follows

(15)
$$[\Phi(Y_{\gamma})](z, z'') = -\operatorname{Tr}(\rho_{(z, z'')}^{\mathcal{T}}(\gamma)) .$$

The trace map is identical to the map in (7) under the identification $\chi_0(\mathcal{T}) \simeq X_0[M]/\mathbb{Z}_2$. The lifting from $PSL(2, \mathbb{C})$ -representation to $SL(2, \mathbb{C})$ -representation is not unique and the two different upliftings, ρ and $\tilde{\rho}$, are related by the action of $H^1(S^3 \setminus \mathcal{K}, \mathbb{Z}_2) = \text{Hom}[\pi_1(S^3 \setminus \mathcal{K}) \to \mathbb{Z}_2] = \{1, \eta\}$, i.e.

(16)
$$\tilde{\rho}(\gamma) = \eta(\gamma)\rho(\gamma) ,$$

For even $Y_{\gamma} \in V^{\text{even}}[M]$, the trace map is well-defined, i.e. independent of the choice of liftings, since $\eta(\gamma) = 1$.



FIGURE 5. Quantum trace map.

2.3. Quantum trace map and Length conjecture. The quantum trace map for 2D surfaces with SL(2) was introduced in [2], and was generalized to groups of higher rank in [11]. Here we introduce a 3D version of the quantum trace map.

Conjecture 2.1 (main conjecture). There exists a unique injective module homomorphism $\operatorname{tr}_q^{\mathcal{T}}: \mathcal{S}_q^{\operatorname{even}}[M] \to \hat{\mathbb{C}}_q[\mathcal{T}]$ which satisfies the following properties:

I. At $q^{1/4} = -1$, the quantum trace map is identical to the trace map $\Phi_{\rm tr}$ (17) $({\rm tr}_q^{\mathcal{T}})|_{q^{1/4}=-1} = \Phi_{\rm tr}$ as a map $V^{\rm even}(M) \to \mathbb{C}[X_0(M)/\mathbb{Z}_2]$.

The relation between $\operatorname{tr}_{q}^{\mathcal{T}}$ and Φ_{tr} is summarized in Figure 5.

II. (All-order length conjecture) For an even link $K = \bigcup_{I=1}^{\sharp(K)} K_I \subset M = S^3 \setminus \mathcal{K}$, (18)

$$\frac{\left.J_{n,\tilde{n}_I=2}(\mathcal{K}\cup K;q)\right|_{q^{1/4}=-\exp\left(\frac{\pi i}{2k}\right)}}{J_n(\mathcal{K};q)} \xrightarrow{\quad n=k;\,k\to\infty} \sum_{s=0}^\infty \mathcal{Z}^{(\mathrm{hyp})}_s(\hat{O}_K;M)\left(\frac{2\pi i}{k}\right)^s \;.$$

or equivalently,

$$\log \frac{J_{n,\tilde{n}_I=2}(\mathcal{K}\cup K;q)}{J_n(\mathcal{K};q)}\Big|_{q^{1/4}=-\exp(\frac{\pi i}{2k})} \xrightarrow{n=k; k\to\infty} \sum_{s=0}^{\infty} \mathcal{W}_s^{(\mathrm{hyp})}(\hat{O}_K;M)\left(\frac{2\pi i}{k}\right)^s$$

Here J denotes a variation of the colored Jones polynomial for framed links. While the conventional colored Jones polynomial is an invariant of oriented links, our Jis an invariant of unoriented framed links. See Appendix **B** for the definition. In the above, $I = 1, \ldots, \sharp(K)$ and \tilde{n}_I denotes the 'color' of the *I*-th component.

The perturbative invariant $\mathcal{W}_s^{(hyp)}(\hat{O}_K; M)$ is determined by the quantum trace operator

(19)
$$\hat{O}_K := \operatorname{tr}_q^{\mathcal{T}}(Y_K) \,.$$

More generally, for an element $Y = \sum_{a} C_a Y_{K_a} \in \mathcal{S}_q^{\text{even}}[M]$,

(20)
$$\hat{O}_Y := \operatorname{tr}_q^{\mathcal{T}}(Y) = \sum_a C_a \hat{O}_{K_a}$$

The operator \hat{O}_K is a $\hat{\mathbb{C}}_q[\mathcal{T}]$ -valued invariant of unoriented framed links.

By incorporating the \hat{O} with the state-integral model [5], the perturbative invariants $\{\mathcal{W}_s^{(\text{hyp})}(\hat{O}_K; M), \mathcal{Z}_s^{(\text{hyp})}(\hat{O}_K; M)\}_{s=0}^{\infty}$ can be obtained from the perturbative expansion of the state-integral model. The explicit form of the state-integral model and its perturbative expansion using Feynman diagram will be given in section 3. Two perturbative expansions are related to each other by exponential or logarithm, i.e.

(21)
$$\exp\left(\sum_{s=0}^{\infty} \mathcal{W}_{s}^{(\text{hyp})} \hbar^{s}\right) = \sum_{s=0}^{\infty} \mathcal{Z}_{s}^{(\text{hyp})} \hbar^{s} \text{ as a formal power series in } \hbar.$$

The leading "classical" coefficient $\mathcal{Z}_{s=0}^{\text{hyp}}(\hat{O}_K; M)$ is simply given by

(22)
$$\mathcal{Z}_{s=0}^{(\text{hyp})}(\hat{O};M) := \left(\hat{O}|_{q^{1/4} \to -1}\right) \Big|_{(z_i, z_i'') = (z_i', z_i'')_{\text{hyp}}}$$

Here $\hat{O}|_{q^{1/4} \to -1}$ is regarded as an element in $\mathbb{C}[\chi_0(\mathcal{T})]$. The classical value is

(23)
$$\mathcal{Z}_{s=0}^{\text{hyp}}(\hat{O}_{K};M) = \text{Tr}\left([\Phi_{\text{tr}}(Y_{\gamma_{K}})](\rho_{\text{hyp}})\right) ,$$

or equivalently, $\mathcal{W}_{s=0}^{(\text{hyp})}(\hat{O}_{K};M) = \log\left([\Phi_{\text{tr}}(Y_{\gamma_{K}})](\rho_{\text{hyp}})\right)$

Here $Y_{\gamma_K} = \prod_{I=1}^n Y_{\gamma_I}$ is an element of the Skein algebra obtained from Y_K at $q^{-1/4} = -1$ and $\rho_{\text{hyp}} \in \text{Hom}[\pi_1(M) \to SL(2,\mathbb{C})]$ is an $SL(2,\mathbb{C})$ flat connection corresponding to the complete hyperbolic structure on M. The classical part is related to the complex length of the geodesic (hence the name "length conjecture") by

(24)
$$[\Phi_{\rm tr}(Y_{\gamma_K})](\rho_{\rm hyp}) := \prod_{I=1}^{\sharp(K)} \left(-\operatorname{Tr}(\rho_{\rm hyp}(\gamma_I)) \right) = \prod_{I=1}^{\sharp(K)} \left(-e^{\frac{1}{2}\ell_{\mathbb{C}}(\gamma_I)} - e^{-\frac{1}{2}\ell_{\mathbb{C}}(\gamma_I)} \right) .$$

 $\ell_{\mathbb{C}}(\gamma_I)$ denotes the complexified length of a geodesic in the same homotopy class as γ_I of K_I

To prove the length conjecture in (18) and (23), we only need to prove them for basis $\{Y_{K_a}\}$ of $\mathcal{S}_q^{\text{even}}[M]$. It is obvious for (18) since for $Y_K = \sum_a C_a(q^{1/4})Y_{K_a}$,

(25)
$$J_{n,\tilde{n}_I=2}(\mathcal{K}\cup K;q) = \sum_a C_a(q^{1/4})J_{n,\tilde{n}_I=2}(\mathcal{K}\cup K_a;q)$$

and the perturbative invariants $\mathcal{Z}_{s}^{(\mathrm{hyp})}(\hat{O}_{Y_{K}})$ satisfy

(26)
$$\sum_{s=0}^{\infty} \mathcal{Z}_{s}^{(\text{hyp})}(\hat{O}_{K}; M) \hbar^{s} = \sum_{a} C_{a}(q^{1/4} = -e^{\frac{\hbar}{4}}) \left(\sum_{s} \mathcal{Z}_{s}^{(\text{hyp})}(\hat{O}_{K_{a}}; M) \hbar^{s} \right) .$$

The relation in (23) is valid for arbitrary \hat{O}_K if it holds for every basis since

(27)
$$\Phi_{\rm tr}(Y_{\gamma_K})(\rho) = \sum_a C_a(q^{1/4} = -1)\Phi_{\rm tr}(Y_{\gamma_K})(\rho)$$
for arbitrary $\rho \in {\rm Hom}\left[\pi_1 M \to SL(2,\mathbb{C})\right],$

which follows from the fact that the quantum trace map $\operatorname{tr}_q^{\mathcal{T}}$ is a module homomorphism and it becomes the classical trace map Φ_{tr} in (7) at $q^{1/4} = -1$.

III. For an even link K and a meridian knot \bigcirc_{m} , a knot linking the heavy knot \mathcal{K} ,

(28)
$$\hat{O}_{K \sqcup \bigcirc_{\mathrm{m}}^{2}} = \hat{O}_{\bigcirc_{\mathrm{m}}^{2}} \cdot \hat{O}_{K}$$
, where $\hat{O}_{\bigcirc_{\mathrm{m}}^{2}} := 2 + e^{\hat{\mathfrak{M}}} + e^{-\hat{\mathfrak{M}}}$

See (131) for the definition of $e^{\pm \hat{\mathfrak{M}}}$ and Figure 8 for the knot \bigcirc_{m} . The operator $\hat{O}_{\bigcirc_{\mathrm{m}}^2}$ always commutes with $e^{\hat{C}_I}$ and thus the left multiplication is well-defined in

 $\hat{\mathbb{C}}_q[\mathcal{T}]$. In section 3, we will see that

(29)
$$\sum_{s=0}^{\infty} \mathcal{Z}_s^{(\text{hyp})} (\hat{O}_{K \sqcup \bigcirc_{\mathrm{m}}^2}) \hbar^s = 4 \sum_{s=0}^{\infty} \mathcal{Z}_s^{(\text{hyp})} (\hat{O}_K) \hbar^s$$

for an arbitrary knot K in $M = S^3 \setminus \mathcal{K}$. It is compatible with the all-order length conjecture (18) since

(30)
$$J_{n,\tilde{n}_I=2}(\mathcal{K} \cup (\mathcal{K} \sqcup \bigcirc_{\mathrm{m}}^2)) = (q^{n/2} + q^{-n/2})^2 \times J_{n,\tilde{n}_I=2}(\mathcal{K} \cup \mathcal{K}) ,$$
and $(q^{n/2} + q^{-n/2})^2|_{q=\exp(2\pi i/k), n=k} = 4 .$

Here we use the property of our Jones polynomial depicted in Figure 8.

3. Perturbative knot invariants

We recall and extend state-integral models from [5, 9, 12] to define the perturbative knot invariants $\{\mathcal{Z}_s^{(hyp)}(\hat{O}_K; M)\}_{s=0}^{\infty}$.

3.1. State-integral model with quantum trace map. The state-integral model is based on an ideal triangulation \mathcal{T} of $M = S^3 \setminus \mathcal{K}$ and can be written as in the following form using Dirac brackets $(k := |\mathcal{T}|)$

$$(31)$$

$$Z_{\hbar}[M + \hat{O}; X_{\mathbf{m}}]$$

$$= \langle X_{\mathbf{m}}, C_{1}, \dots, C_{k-1}; \Pi_{M,C} | \hat{O}(q^{\frac{1}{4}} = -e^{\frac{\hbar}{4}}) | \Delta^{\otimes k} \rangle |_{C_{1} = \dots = C_{k-1} = 0},$$

$$= \int \prod_{i=1}^{k} \frac{dZ_{i}}{\sqrt{2\pi\hbar}} \langle X_{\mathbf{m}}, C_{1}, \dots, C_{k-1}; \Pi_{M,C} | Z_{1}, \dots, Z_{k}; \Pi_{Z} \rangle \langle Z_{1}, \dots, Z_{k}; \Pi_{Z} | \hat{O} | \Delta^{\otimes k} \rangle$$

The state-integral is a function on a meridian variable $X_{\mathbf{m}}$. Here $|X_1, \ldots, X_k; \Pi_X \rangle$ denotes the position basis of $(\mathcal{H}(\partial \Delta))^{\otimes k}$, the Hilbert-space associated with the *k*-tetrahedra [5], with respect to a polarization choice $\Pi_X = (X_1, \ldots, X_k, P_1, \ldots, P_k)^T$:

(32)

$$\langle X_1, \dots, X_k; \Pi_X | e^{\hat{X}_i} | \psi \rangle = e^{X_i} \langle X_1, \dots, X_k; \Pi_X | \psi \rangle \text{ and } \langle X_1, \dots, X_k; \Pi_X | e^{\hat{P}_i} | \psi \rangle = \langle X_1, \dots, X_k; \Pi_X | \psi \rangle |_{X_i \to X_i + \hbar}, \ \forall | \psi \rangle \in (\mathcal{H}(\partial \Delta))^{\otimes k}$$

Here two polarizations $\Pi_{M,C}$ and Π_Z are

(33)
$$\Pi_{M,C} = \begin{pmatrix} \mathfrak{M} \\ C_1 \\ \cdots \\ C_{k-1} \\ \frac{1}{2} \mathfrak{L} \\ \Gamma_1 \\ \cdots \\ \Gamma_{k-1} \end{pmatrix}, \quad \Pi_Z = \begin{pmatrix} Z_1 \\ Z_2 \\ \cdots \\ Z_k \\ Z''_1 \\ Z''_2 \\ \cdots \\ Z''_k \end{pmatrix}$$

and they are related to each other by a linear canonical transformation

(34)
$$\Pi_{M,C} = g \cdot \Pi_Z - (i\pi + \frac{\hbar}{2}) \begin{pmatrix} \vec{\nu} \\ \vec{\nu}_p \end{pmatrix}, \text{ where } g \in Sp(2k,\mathbb{Z}) \text{ and } \vec{\nu}, \vec{\nu}_p \in \mathbb{Z}^k.$$

 Γ 's are chosen such that the linear transformation becomes a canonical transformation. The choice is not unique but the final state-integral does not depend on it. The two position bases are related to each other by following unitary transformation

$$\langle X_{\mathbf{m}}, C_1, \dots, C_{k-1}; \Pi_{M,C} | Z_1, \dots, Z_k; \Pi_Z \rangle$$
$$= \frac{1}{\sqrt{\det B}} \exp\left(\frac{1}{2\hbar}Q(\vec{X}, \vec{Z})\right) \Big|_{\vec{X} = (X_{\mathbf{m}}, C_1, \dots, C_{k-1})}^{\vec{Z} = (Z_1, \dots, Z_k)}$$

(35) with

$$\begin{split} Q(\vec{X},\vec{Z}) &= \vec{Z}B^{-1}A\vec{Z} + \vec{X}DB^{-1}\vec{X} + (2\pi i + \hbar)\vec{f}B^{-1}\vec{X} + (i\pi + \frac{\hbar}{2})^2\vec{f}B^{-1}\vec{\nu} \\ &- \vec{Z}B^{-1}((2i\pi + \hbar)\vec{\nu} + 2\vec{X}) \;. \end{split}$$

Here A, B, C and D are the four $(k \times k)$ block matrices of g:

$$(36) g = \begin{pmatrix} A & B \\ C & D \end{pmatrix}$$

The vectors $(\vec{f}, \vec{f}'') = (f_i, f''_i)_{i=1}^k$ are known as combinatorial flattening, and chosen to satisfy the following relation

(37)
$$\begin{pmatrix} A & B \\ C & D \end{pmatrix} \cdot \begin{pmatrix} \vec{f} \\ \vec{f''} \end{pmatrix} = \begin{pmatrix} \vec{\nu} \\ \vec{\nu}_p \end{pmatrix} +$$

 $|\Delta^{\otimes k}\rangle \in (\mathcal{H}(\partial \Delta))^{\otimes k}$ is the wave-function for k-tetrahedra satisfying the following difference equations

(38)
$$(e^{\hat{Z}''_i} + e^{-\hat{Z}_i} - 1) |\Delta^{\otimes k}\rangle = 0, \quad i = 1, \dots k.$$

In the polarization Π_Z , the wave-function is given by a product of quantum dilogarithms (see Appendix C for our convention for the quantum dilogarithm)

(39)
$$\langle Z_1, \dots, Z_k; \Pi_Z | \Delta^{\otimes k} \rangle = \prod_{i=1}^k \psi_{\hbar}(Z_i) .$$

Gathering all the expressions above, one has

$$\begin{aligned} &(40)\\ &Z_{\hbar}[M+\hat{O};X_{\mathbf{m}}] \\ &= \frac{1}{\sqrt{\det B}} \sum_{\alpha} C_{\alpha}(q^{\frac{1}{4}}) \big|_{q^{\frac{1}{4}} = -e^{\frac{\hbar}{4}}} \int \prod_{i=1}^{k} \frac{dZ_{i}}{\sqrt{2\pi\hbar}} \exp\left(\frac{1}{2\hbar}Q(\vec{X},\vec{Z})\right) \Big|_{\vec{X}=(X_{\mathbf{m}},\cdots,C_{k-1})}^{\vec{Z}=(Z_{1},\ldots,Z_{k})} \\ &\times \langle Z_{1},\ldots,Z_{k};\Pi_{Z}|\prod_{i} \hat{z}_{i}^{a_{i}^{(\alpha)}}(\hat{z}_{i}^{\prime\prime})^{b_{i}^{(\alpha)}}|\Delta^{\otimes k}\rangle , \\ &= \frac{1}{\sqrt{\det B}} \sum_{\alpha} C_{\alpha}(q^{\frac{1}{4}}) \big|_{q^{\frac{1}{4}} = -e^{\frac{\hbar}{4}}} \int \prod_{i=1}^{k} \frac{dZ_{i}}{\sqrt{2\pi\hbar}} \exp\left(\frac{1}{2\hbar}Q(\vec{X},\vec{Z})\right) \Big|_{\vec{X}=(X_{\mathbf{m}},\cdots,C_{k-1})}^{\vec{Z}=(Z_{1},\ldots,Z_{k})} \\ &\times \prod_{i=1}^{k} \psi_{\hbar}(Z_{i}+b_{i}^{(\alpha)}\hbar) \exp\left(a_{i}^{(\alpha)}Z_{i}\right) , \end{aligned}$$

when $\hat{O} = \sum_{\alpha} C_{\alpha}(q^{\frac{1}{4}}) \prod_{i} \hat{z}_{i}^{a_{i}^{(\alpha)}}(\hat{z}_{i}'')^{b_{i}^{(\alpha)}} \in \hat{\mathbb{C}}_{q}[\mathcal{T}]$. From the expression in (31), it is not difficult to see that the $Z_{\hbar}[M + \hat{O}; X_{\mathbf{m}}]$ with $\hat{O} \in \hat{\mathbb{C}}_{q}[\mathcal{T}]$ is well-defined (recall

the definition of $\hat{\mathbb{C}}_q[\mathcal{T}]$ in (10)), i.e.

(41)
$$Z_{\hbar}[M+\hat{O}] = Z_{\hbar}[M+e^{\hat{C}_{I}}\cdot\hat{O}] = Z_{\hbar}[M+\hat{O}\cdot(\hat{z}_{i}^{-1}+\hat{z}_{i}'')]$$

It is also straightforward to see that

(42)
$$Z_{\hbar}[M + F(e^{\mathfrak{M}}) \cdot \hat{O}] = F(e^{X_{\mathbf{m}}})Z_{\hbar}[M + \hat{O}],$$

for arbitrary Laurent polynomial F(z) and $\hat{O} \in \hat{\mathbb{C}}_q[\mathcal{T}]$.

3.2. **Perturbative invariants.** We are ready to define and explain the perturbative invariants $\{\mathcal{Z}_s^{(\text{hyp})}(\hat{O}; M)\}_{s=0}^{\infty}$. By expanding the state-integral in the limit $\hbar \to 0$ around the saddle point $\vec{Z}^{(\text{hyp})}$ associated to complete hyperbolic structure of M, one have

(43)
$$Z_{\hbar}[M+\hat{O};X_{\mathbf{m}}=0] \xrightarrow{\hbar \to 0} \exp\left(\sum_{n=0}^{\infty} S_n^{(\mathrm{hyp})}(\hat{O};M)\hbar^{n-1}\right) .$$

In the expansion, one can use the asymptotic expansion of the quantum dilogarithm function in (155). Then, we define the perturbative invariants $\{\mathcal{Z}_s^{(hyp)}, \mathcal{W}_s^{(hyp)}\}_{s=0}^{\infty}$ as

(44)
$$\frac{\exp\left(\sum_{n=0}^{\infty} S_n^{(\text{hyp})}(\hat{O}; M)\hbar^{n-1}\right)}{\exp\left(\sum_{n=0}^{\infty} S_n^{(\text{hyp})}(M)\hbar^{n-1}\right)} = \sum_{s=0}^{\infty} \mathcal{Z}_s^{(\text{hyp})}(\hat{O}; M)\hbar^s$$
$$= \exp\left(\sum_{s=0}^{\infty} \mathcal{W}_s^{(\text{hyp})}(\hat{O}; M)\hbar^s\right)$$

Here $S_n^{(\text{hyp})}(M) := S_n^{(\text{hyp})}(\hat{O}_{K=\emptyset}; M)$ is the perturbative invariant of state-integral model without any insertion of loop operator. From the definition, one can see that the relation in (29) simply follows from (42) with $F(x) = 2 + x + \frac{1}{x}$ and $X_{\mathbf{m}} = 0$. In the classical limit $\hbar \to 0$, saddle points $\{\vec{Z} = \vec{Z}^{(\alpha)}\}$ of the state-integral at $X_{\mathbf{m}} = 0$ satisfy following equations

(45)
$$\prod_{j=1}^{\kappa} z_j^{A_{ij}} (1 - z_j^{-1})^{B_{ij}} \Big|_{z_i = \exp(Z_i^{(\alpha)})} = (-1)^{\nu_i} .$$

For ρ_{hyp} -regular ideal triangulation \mathcal{T} , the saddle point $\vec{Z}^{(\text{hyp})}$ is uniquely characterized by following conditions $(z_i^{(\text{hyp})} := \exp(Z_i^{(\text{hyp})}))$

(46)
$$\operatorname{Im}[z_i^{(\operatorname{hyp})}] < 0 \quad \text{for all } i = 1, \dots, k$$

Example 3.1. As an example, consider the case with $M = S^3 \setminus \mathbf{4}_1$. Its simplest ideal triangulation consists of two tetrahedra, say Δ_Y and Δ_Z , with edge parameters

$$(y, y', y'') = (e^Y, e^{Y'}, e^{Y''})$$
 and $(z, z', z'') = (e^Z, e^{Z'}, e^{Z''}).$

Using the gluing data given in Appendix A, we have

(47)
$$\begin{pmatrix} \mathfrak{M} \\ C_1 \\ \frac{1}{2}\mathfrak{L} \\ \Gamma_1 \end{pmatrix} = \begin{pmatrix} 0 & -1 & 1 & 0 \\ 1 & 1 & -1 & -1 \\ 0 & 1 & 0 & -1 \\ 0 & 0 & 1 & 0 \end{pmatrix} \cdot \begin{pmatrix} Y \\ Z \\ Y'' \\ Z'' \end{pmatrix}.$$

12 PRARIT AGARWAL, DONGMIN GANG, SANGMIN LEE, AND MAURICIO ROMO

The corresponding state-integral model is [5, 9]

(48)
$$Z_{\hbar}(S^{3}\backslash \mathbf{4}_{1}; X_{\mathbf{m}}) = \langle X_{\mathbf{m}}, C_{1} = 0 | \Delta^{\otimes 2} \rangle ,$$
$$= \int \frac{dY dZ}{2\pi\hbar} \exp\left(\frac{1}{2\hbar} (X_{\mathbf{m}}^{2} - 2X_{\mathbf{m}}Y + 2X_{\mathbf{m}}Z - 2YZ)\right) \psi_{\hbar}(Y) \psi_{\hbar}(Z) .$$

The saddle points of the integral are the solutions of the following equations:

(49)
$$\partial_{Y,Z} \left(\operatorname{Li}_2(e^{-Y}) + \operatorname{Li}_2(e^{-Z}) + \frac{1}{2}X_{\mathbf{m}}^2 + X_{\mathbf{m}}Z - X_{\mathbf{m}}Y - YZ \right) = 0 ,$$
$$\Longrightarrow \quad -Z + \log(1 - e^{-Y}) = X_{\mathbf{m}} , \quad -Y + \log(1 - e^{-Z}) = -X_{\mathbf{m}} .$$

The saddle point equations coincide with the (logarithmic) gluing equations of the ideal triangulation. At $X_{\mathbf{m}} = 0$, there are two saddle points,

(50)

$$(\alpha) = (\overline{\text{hyp}}) : Y = Z = \frac{i\pi}{3},$$

$$(\alpha) = (\text{hyp}) : Y = Z = -\frac{i\pi}{3}.$$

We now extend the state-integral model for $S^3 \setminus \mathbf{4}_1$ to include a quantum trace map \hat{O}_{K_b} , where K_b is the geodesic knot in the homotopy class $b \in \pi_1(S^3 \setminus \mathbf{4}_1)$, see Figure 6. The classical trace map is

(51)
$$\Phi_{\rm tr}(Y_{K_b}) = -(y^{-1} + z^{-1} - y^{-1}z^{-1}) \in \mathbb{C}[\chi[\mathcal{T}]]$$

The expression is obtained using the holonomy matrices in (140) and the gluing equation varieity $\chi[\mathcal{T}]$ is given in (139). After quantization, we assume that the quantum loop operator is given as

(52)
$$\hat{O}_{K_b} = \operatorname{tr}_q(Y_{K_b}) = -C_1 \hat{y}^{-1} - C_2 \hat{z}^{-1} + C_3 \hat{y}^{-1} \hat{z}^{-1} \in \hat{\mathbb{C}}_q[\mathcal{T}]$$

where $\{C_{\alpha=1,2,3}\}$ are Laurent polynomials in $q^{1/4}$ which all become 1 in the limit $q^{1/4} \rightarrow -1$. Later in section 4, we will propose that $C_1 = C_2 = C_3 = q^{1/2} = e^{\frac{\hbar}{2}}$ and check it against the all-order length conjecture. Under the quantization, the state-integral is

(53)
$$Z_{\hbar} \left[S^{3} \backslash \mathbf{4}_{1} + \hat{O}_{K_{b}}; X_{\mathbf{m}} = 0 \right]$$
$$= \int \frac{dY dZ}{2\pi\hbar} \exp\left(-\frac{YZ}{\hbar}\right) \psi_{\hbar}(Y) \psi_{\hbar}(Z) \times e^{\hbar/2} \left(-e^{-Y} - e^{-Z} + e^{-Y-Z}\right)$$

For K_b^2 (2-cabling of K_b), the classical trace map is

(54)
$$\Phi_{\rm tr}(Y_{K_b^2}) = (\Phi_{\rm tr}(Y_{K_b}))^2 = (y^{-1} + z^{-1} - y^{-1}z^{-1})^2 \in \mathbb{C}[\chi_0[\mathcal{T}]].$$

After quantization, as will be proposed in section 4, the quantum loop operator is

(55)
$$\hat{O}_{K_b^2} = q(\hat{y}^{-1} + \hat{z}^{-1} - \hat{y}^{-1}\hat{z}^{-1})^2 + 1 - q^2 \in \hat{\mathbb{C}}_q[\mathcal{T}] .$$

The state-integral with the loop operator is

(56)

$$Z_{\hbar} \left[S^{3} \backslash \mathbf{4}_{1} + \hat{O}_{K_{b}^{2}}; X_{\mathbf{m}} = 0 \right]$$

$$= \int \frac{dYdZ}{2\pi\hbar} \exp\left(-\frac{YZ}{\hbar}\right) \psi_{\hbar}(Y) \psi_{\hbar}(Z) \times \left(e^{\hbar} \left(e^{-Y} + e^{-Z} - e^{-Y-Z}\right)^{2} + 1 - e^{2\hbar}\right).$$

3.3. Feynman perturbation theory for the invariants. In this section we will give an expression for the invariants $\{\mathcal{W}_s^{\text{hyp}}(\hat{O}_K; M)\}_{s=0}^{\infty}$ using Feynman diagrams following [9]. The quantum loop operator \hat{O} can be generally given in the following form

(57)
$$\hat{O} = \sum_{\alpha} C_{\alpha}(q^{1/4}) \hat{O}_{\alpha} = \sum_{\alpha} C_{\alpha}(q^{1/4}) \prod_{i} \hat{z}_{i}^{a_{i}^{(\alpha)}}(\hat{z}_{i}'')^{b_{i}^{(\alpha)}}.$$

Giving an explicit formula for each of the terms $\mathcal{W}_s^{\text{hyp}}(\hat{O}_K; M)$ is quite involved. We will instead give an implicit expression in terms of the perturbative expansion of each term $Z_{\hbar}(M + \hat{O}_{\alpha}; X_{\mathbf{m}})$ in the expression

(58)
$$Z_{\hbar}(M+\hat{O};X_{\mathbf{m}}) = \sum_{\alpha} C_{\alpha}(q^{\frac{1}{4}}) Z_{\hbar}(M+\hat{O}_{\alpha};X_{\mathbf{m}}) \big|_{q^{\frac{1}{4}} = -e^{\frac{\hbar}{4}}},$$

which we denote them as $S_{n;\alpha}^{(c)}$:

(59)
$$Z_{\hbar}(M + O_{\alpha}; X_{\mathbf{m}}) = \frac{1}{\sqrt{\det B}} \int \prod_{i=1}^{k} \frac{dZ_{i}}{\sqrt{2\pi\hbar}} \exp\left(\frac{1}{2\hbar}Q(Z, X_{\mathbf{m}})\right) e^{a^{(\alpha)} \cdot Z} \prod_{i=1}^{k} \psi_{\hbar}(Z_{i} + b_{i}^{(\alpha)}\hbar)$$
$$\xrightarrow{\text{around a saddle point } Z^{(c)} \text{ in } \hbar \to 0} \exp\left(\sum_{n=0}^{\infty} \hbar^{n-1}S_{n;\alpha}^{(c)}\right).$$

The expansion can be obtained by considering the expansion $Z_i = Z_i^{(c)} + Y_i$. So, the integral expression becomes

(60)
$$Z_{\hbar}(M + \hat{O}_{\alpha}; X_{\mathbf{m}}) = \frac{e^{\Gamma^{(0)}}}{\sqrt{\det B}} \int \prod_{i=1}^{k} \frac{dY_{i}}{\sqrt{2\pi\hbar}} e^{\frac{1}{2\hbar}Y_{i}H^{ij}Y_{j}} \prod_{i=1}^{k} e^{\sum_{s=1}^{\infty} \Gamma_{i}^{(s)}Y_{i}^{s}/s!},$$

where we defined

(61)
$$H^{ij} := (B^{-1}A)_{ij} + (z_j z_j'')^{-1} \delta_{ij},$$

the 'vacuum energy':

(62)
$$\Gamma^{(0)} = F(X_{\mathbf{m}}, Z^{(c)}) + \sum_{i=1}^{k} \sum_{n=1}^{\infty} \frac{B_n (1 + b_i^{(\alpha)}) \hbar^{n-1}}{n!} \operatorname{Li}_{2-n} \left(e^{-Z_i^{(c)}} \right),$$

where $B_n(1+b_i^{(\alpha)})$ denotes the Bernoulli polynomials evaluated at $1+b_i^{(\alpha)}$ and

$$F(X_{\mathbf{m}}, Z^{(c)}) := \frac{1}{2\hbar} Q(Z^{(c)}, X_A) + a \cdot Z^{(c)} + \frac{1}{\hbar} \sum_i \operatorname{Li}_2(e^{-Z_i^{(c)}}),$$

the linear vertex

(63)
$$\Gamma_i^{(1)} = a_i^{(\alpha)} + \frac{1}{2} - \sum_{n=1}^{\infty} \frac{B_n (1 + b_i^{(\alpha)}) \hbar^{n-1}}{n!} \operatorname{Li}_{1-n} \left(e^{-Z_i^{(c)}} \right),$$

and higher valence vertices $(k \ge 2)$:

(64)
$$\Gamma_i^{(k)} = (-1)^k \sum_{n=0}^{\infty} \frac{B_n (1 + b_i^{(\alpha)}) \hbar^{n-1}}{n!} \operatorname{Li}_{2-n-k} \left(e^{-Z_i^{(c)}} \right).$$

14 PRARIT AGARWAL, DONGMIN GANG, SANGMIN LEE, AND MAURICIO ROMO

Then, the coefficients $S_{n;\alpha}^{(c)}$ can be computed as sums over Feynman diagrams corresponding to non-directed graphs, without open edges, which we denote as \mathcal{G}_{Γ} . The situation is completely analogous to [9], with a slight modification on the valence of the vertices. We associate a weight to each \mathcal{G}_{Γ} :

(65)
$$W(\mathcal{G}_{\Gamma}) := \frac{1}{|\operatorname{Aut}(\mathcal{G}_{\Gamma})|} \sum_{\text{labels}} \prod_{v \in \text{vertices}} \Gamma_{l_v}^{(k_v)} \prod_{e \in \text{edges}} \Pi_e$$

here l_v is the set of labels corresponding to the legs of the vertex v in the graph and k_v its valence. Given a graph \mathcal{G}_{Γ} we denote \mathcal{L} the number of loops, V_a the number of vertices of valence a. We define

(66)
$$\mathcal{G}_n := \{\mathcal{G}_{\Gamma} : \mathcal{L} + V_1 + V_2 \le n\}$$

then, the $S_{n;\alpha}^{(c)}$ coefficients are given by

(67)
$$S_{0;\alpha}^{(c)} = \operatorname{coeff}[\Gamma^{(0)}, \hbar^{-1}]$$

the 1-loop term:

(68)
$$e^{S_{1,\alpha}^{(c)}} = \frac{e^{\operatorname{coeff}[\Gamma^{(0)},\hbar^0]}}{\sqrt{\det B}} \frac{i^k}{\sqrt{\det H}}$$

and for $n \geq 2$:

(69)
$$S_{n,k}^{\alpha} = \operatorname{coeff}\left\{\Gamma^{(0)} + \sum_{\mathcal{G}_{\Gamma} \in \mathcal{G}_{n}} W(\mathcal{G}_{\Gamma}), \hbar^{n-1}\right\} \qquad n \ge 2.$$

Now we can use these results to compute the invariants $\{\mathcal{W}^{hyp}_s(\hat{O};M)\}_{s=0}^{\infty}$. First denote

(70)
$$\frac{Z_{\hbar}(M+\hat{O};X_{\mathbf{m}})}{Z_{\hbar}(M;X_{\mathbf{m}})} = \frac{1}{Z_{\hbar}(M;X_{\mathbf{m}})} \sum_{\alpha} C_{\alpha}(q^{\frac{1}{4}}) Z_{\hbar}(M+\hat{O}_{\alpha};X_{\mathbf{m}})$$
$$\xrightarrow{\text{around a saddle point } Z^{(c)} \text{ in } \hbar \to 0} \sum_{\alpha} C_{\alpha}(q^{\frac{1}{4}} = -e^{\frac{\hbar}{4}}) \exp\left(\sum_{n=1}^{\infty} \hbar^{n-1} \widetilde{S}_{n,\alpha}^{(c)}\right)$$

where

(71)
$$\widetilde{S}_{n,\alpha}^{(c)} := S_{n,\alpha}^{(c)} - S_n^{(c)}(M) \qquad n \ge 1$$

and we have used the fact that

(72)
$$S_0^{(c)}(M) = S_{0,\alpha}^{(c)}$$
 for all α .

Therefore, the expression of $\widetilde{S}_{n,\alpha}^{(c)}$, $n \geq 1$ in terms of Feynman diagrams reduce exactly to the same computation than $S_n^{(c)}(M)$, presented in [9] with two modifications in the vertices:

(1) The vertices of valence 0 and 1 becomes

$$\Gamma^{(0)} = a^{(\alpha)} \cdot Z^{(c)} + \sum_{i=1}^{k} \sum_{n=1}^{\infty} \frac{(B_n(1+b_i^{(\alpha)}) - B_n)\hbar^{n-1}}{n!} \operatorname{Li}_{2-n} \left(e^{-Z_i^{(c)}}\right),$$
(73) $\Gamma_i^{(1)} = a_i^{(\alpha)} - \sum_{n=1}^{\infty} \frac{(B_n(1+b_i^{(\alpha)}) - B_n)\hbar^{n-1}}{n!} \operatorname{Li}_{1-n} \left(e^{-Z_i^{(c)}}\right),$

(2) The coefficient $B_n = B_n(1)$ is replaced by $B_n(1 - b_i^{(\alpha)}) - B_n$ in the vertex $\Gamma_i^{(s)}, s \ge 2$

The coefficients C_{α} are Laurent polynomials in $q^{1/4}$, so they can be expanded in power series in \hbar :

(74)
$$C_{\alpha}(q^{\frac{1}{4}} = -e^{\frac{\hbar}{4}}) = \sum_{s \ge 0} \hbar^s f_{s,\alpha}.$$

Therefore, the series expression in (70) can be written as a formal power series in \hbar whose $\mathcal{O}(\hbar^0)$ term is given by

(75)
$$\frac{Z_{\hbar}(M+O;X_{\mathbf{m}})}{Z_{\hbar}(M;X_{\mathbf{m}})} = \sum_{c} W_{\gamma}^{(c)} + \mathcal{O}(\hbar)$$

where $W_K^{(c)} := \sum_{\alpha} f_{0,\alpha} e^{\tilde{S}_{1,\alpha}^{(c)}}$ is the classical expectation value of the light knot K at the saddle point (c). Therefore, for a fixed saddle point (c) we can write (75) as

(76)
$$\ln\left(W_K^{(c)}\right) + \ln\left(1 + P^{(c)}(\hbar)\right)$$

where $P^{(c)}(\hbar)$ is a power series in \hbar , without a constant term. Hence, the Taylor series of the logarithm defines the invariants $\{\mathcal{W}_s^{(\mathrm{hyp})}(\hat{O}_K; M)\}_{s=0}^{\infty}$. They become very involved as s grows larger and we find it not very illuminating to include a general formula here. Instead we we will present explicit expressions for s = 0, 1.

For s = 0, we get the classical expectation value, as remarked above

(77)

$$W_{K}^{\text{class}} := \exp\left(\mathcal{W}_{0}^{(\text{hyp})}(\hat{O}; M)\right) = \sum_{\alpha} f_{0,\alpha} e^{\tilde{S}_{1,\alpha}^{(\text{hyp})}} = \sum_{\alpha} C_{\alpha}(q^{1/4} = -1) \prod_{i} (z_{i}'')^{b_{i}^{(\alpha)}} z_{i}^{a_{i}^{(\alpha)}},$$

where $z_i'' = (1 - e^{-Z_i})$ and $z_i = e^{Z_i}$ are evaluated at the saddle point $Z^{(\text{hyp})}$ in (46). Hence the value of $\exp\left(\mathcal{W}_0^{(\text{hyp})}(\hat{O}; M)\right)$ coincides with the classical expectation value of \hat{O} , as expected.

For s = 1, we get

(78)
$$\mathcal{W}_{1}^{(\text{hyp})}(\hat{O};M) = \frac{\sum_{\alpha} (f_{1,\alpha} + f_{0,\alpha} \tilde{S}_{2,\alpha}^{(\text{hyp})}) e^{\tilde{S}_{1,\alpha}^{(\text{hyp})}}}{W_{K}^{\text{class}}}$$

where the explicit expressions for $\widetilde{S}_{1,\alpha}^{(\text{hyp})}$, $\widetilde{S}_{2,\alpha}^{(\text{hyp})}$ can be obtained with the help of the Feynman diagram expansion as described above. In the next section we will provide explicit exampled for links in the figure-eight knot-complement.

4. Example : figure-eight knot-complement

We explicitly define the quantum trace map for $M = S^3 \setminus \mathbf{4}_1$ and prove (or numerically check) some parts of conjecture in (2.1).

The Skein module for $M = S \setminus \mathbf{4}_1$ is studied in [3]:

$$\mathcal{S}_q(S^3 \setminus \mathbf{4}_1) = \mathbb{C}[q^{\pm 1/4}] \text{-module with basis } \{Y_{K_b^n \sqcup \bigcirc_m^m}\}_{m \ge 0, 0 \le n \le 2}$$

(79)

Where \bigcirc_{m} is the meridian knot, i.e. the knot in the red color, and K_b is the knot in green color depicted in Figure 6. K_b is an even knot while \bigcirc_{m} is an odd knot. K^n denotes the *n*-cabling of a knot K.



FIGURE 6. Three generators of fundamental group, given in (133), of the figure-eight knot complement.

4.1. Explicit quantum trace map. The quantum gluing module $\hat{\mathbb{C}}_q[\mathcal{T}]$ for an ideal triangulation \mathcal{T} in Appendix A is

$$(80) \quad \hat{\mathbb{C}}_{q}[\mathcal{T}] = \frac{\mathbb{C}[q^{\pm 1/4}][\hat{y}^{\pm 1}, \hat{z}^{\pm 1}, (\hat{y}'')^{\pm 1}, (\hat{z}'')^{\pm 1}] / \langle \hat{z}\hat{z}'' - q^{-1}\hat{z}''\hat{z}, \hat{y}\hat{y}'' - q^{-1}\hat{y}''\hat{y} \rangle}{\text{equivalence relations} \sim} ,$$
where $\hat{O} \sim q^{-1}\hat{y}\hat{z}(\hat{y}''\hat{z}'')^{-1}\hat{O} , \ \hat{O} \sim \hat{O} \cdot (\hat{y}^{-1} + \hat{y}'') \text{ and } \hat{O} \sim \hat{O} \cdot (\hat{z}^{-1} + \hat{z}'') .$

where $O = q g_{\mathcal{L}}(g_{\mathcal{L}}) = 0$, O = O = (g + g) and O = O = (z + f)

We propose following quantum trace map $\operatorname{tr}_q^{\mathcal{T}}$: $\mathcal{S}_q^{\operatorname{even}}[M] \to \hat{\mathbb{C}}_q[\mathcal{T}]$

$$\operatorname{tr}_{q}^{\tau}(Y_{\bigcirc_{\mathrm{m}}^{2m}}) = (\hat{y}''/\hat{z} + \hat{z}/\hat{y}'' + 2)^{m} ,$$

$$(81) \quad \operatorname{tr}_{q}^{\tau}(Y_{K_{b}\sqcup\bigcirc_{\mathrm{m}}^{2m}}) = -q^{1/2} \left(\hat{y}''/\hat{z} + \hat{z}/\hat{y}'' + 2\right)^{m} \left(\hat{y}^{-1} + \hat{z}^{-1} - \hat{y}^{-1}\hat{z}^{-1}\right) ,$$

$$\operatorname{tr}_{q}^{\tau}(Y_{K_{b}^{2}\sqcup\bigcirc_{\mathrm{m}}^{2m}}) = \left(\hat{y}''/\hat{z} + \hat{z}/\hat{y}'' + 2\right)^{m} \left(q(\hat{y}^{-1} + \hat{z}^{-1} - \hat{y}^{-1}\hat{z}^{-1})^{2} - q^{2} + 1\right) .$$

Assuming the conjecture III in (2.1) and using the fact $e^{\hat{\mathfrak{M}}} = \hat{y}''/\hat{z}$, the only nontrivial parts of the proposal are $\operatorname{tr}_q^{\mathcal{T}}(Y_{K_b})$ and $\operatorname{tr}_q^{\mathcal{T}}(Y_{K_b^2})$. We support the proposal by checking the all-order length conjecture.

by checking the all-order length conjecture. The perturbative invariant $\mathcal{W}_s^{(\text{hyp})}(\hat{O}_K; S^3 \setminus \mathbf{4}_1)$ defined in (44) for $K = K_b$ and $K = K_b^2$ can be computed by expanding the state-integral models in (53) and (56) around the saddle point in (50). Up to s = 2, one has

(82)
$$\sum_{s=0}^{2} \mathcal{W}_{s}^{(\text{hyp})}(\hat{O}_{K_{b}}; S^{3} \setminus \mathbf{4}_{1})\hbar^{s} = \frac{7i\pi + 3\log 3}{6} + \frac{i\sqrt{3} + 6}{18}\hbar + \frac{i}{18\sqrt{3}}\hbar^{2},$$
$$\sum_{s=0}^{2} \mathcal{W}_{s}^{(\text{hyp})}(\hat{O}_{K_{b}^{2}}; S^{3} \setminus \mathbf{4}_{1})\hbar^{s} = \frac{i\pi + 3\log 3}{3} + \frac{i\sqrt{3} + 2}{3}\hbar + \frac{i\sqrt{3} + 6}{27}\hbar^{2}.$$

The leading part is related to the complex length of γ_b as follows

(83)

$$\begin{aligned}
\mathcal{W}_{0}^{(\text{hyp})}(\hat{O}_{K_{b}}; S^{3} \setminus \mathbf{4}_{1}) &= \log \left(- \left(e^{\ell_{\mathbb{C}}(\gamma_{b})/2} + e^{-\ell_{\mathbb{C}}(\gamma_{b})/2} \right) \right) \\
&= \frac{7i\pi + 3\log 3}{6} \simeq 0.549306 + 3.66519i . \\
\mathcal{W}_{0}^{(\text{hyp})}(\hat{O}_{K_{b}^{2}}; S^{3} \setminus \mathbf{4}_{1}) &= 2\mathcal{W}_{0}(\hat{O}_{K_{b}}; S^{3} \setminus \mathbf{4}_{1}) \pmod{2\pi i} \\
&\simeq 1.09861 + 1.0472i .
\end{aligned}$$

4.2. Test of the length conjecture - analytic proof up to 1-loop. For the figure-8 knot complement, one of the proofs of the classical volume conjecture [23] is based on the Habiro-Lê formula [17, 22] for the Jones polynomial:

(84)

$$J_n(q) := \frac{J_n(K = \mathbf{4}_1; q)}{J_n(\bigcirc; q)}$$

$$= \sum_{k=0}^{n-1} \prod_{j=1}^k \left(q^{(n-j)/2} - q^{-(n-j)/2} \right) \left(q^{(n+j)/2} - q^{-(n+j)/2} \right) = \sum_{k=0}^{n-1} w_{n,k}(q) \,.$$

At $q = e^{2\pi i/n}$, the function $w_{n,k}(q)$ simplifies quite a bit:

(85)
$$w_{n,k}(q = e^{2\pi i/n}) = \prod_{j=1}^{k} 4\sin^2\left(\frac{j\pi}{n}\right) \,.$$

In the remainder of this section, $w_{n,k}$ will mean this this product depending on n and k. Since $w_{n,k}$ is positive definite, it makes sense to regard $J_n^* := J_n(q = e^{2\pi i/n})$ as a sort of partition function, and use it to define a probability distribution:

(86)
$$p_{n,k} := \frac{w_{n,k}}{J_n^*}, \quad p_{n,k} \ge 0, \quad \sum_{k=0}^n p_{n,k} = 1.$$

(The two boundary cases, k = 0 and k = n, have zero probability.)

We provide an analytic proof of the length conjecture for the figure 8-knot complement in the presence of the K_b or K_b^2 . In this context, the length conjecture is a prediction on the large n limit of the series,

(87)
$$A_n^{K_b} = \log\left(\frac{J_{n,2}^*}{J_n^*}\right), \quad A_n^{K_b^2} = \log\left(1 + \frac{J_{n,3}^*}{J_n^*}\right).$$

Our proof relies on simple formulas for $J_{n,2}(q)$ and $J_{n,3}(q)$:

(88)

$$J_{n,2}(q) := \frac{(-1) \times J_{n,\tilde{n}=2}(K = \mathbf{4}_1 \cup K_b; q)}{J_n(\bigcirc; q)}$$

$$= q^{1/2} J_n(q) + q^{-1/2} \sum_{k=0}^{n-1} q^{-k} w_{n,k}(q),$$

$$J_{n,3}(q) := \frac{J_{n,\tilde{n}=3}(K = \mathbf{4}_1 \cup K_b; q)}{J_n(\bigcirc; q)}$$

$$= (1+q) \sum_{k=0}^{n-1} q^{-k} w_{n,k}(q) + q^{-1} \sum_{k=0}^{n-1} q^{-2k} w_{n,k}(q).$$

For
$$J_{n,2}$$
, we include a sign factor to compensate the sign factor appearing in (148).
After the compensation, it becomes identical to the conventional Jones polynomial.
These formulas are *conjectural*. We have verified that they agree with computations
based on braid diagrams and *R*-matrices up to $n = 20$. It would be interesting to
derive these formulas from first principles, for example, using recursion relations [15]
which proved extremely useful for Jones polynomials for single-component knots.

18 PRARIT AGARWAL, DONGMIN GANG, SANGMIN LEE, AND MAURICIO ROMO

With the probability interpretation of (86), the series A_n are rephrased as

(90)
$$A_n^{K_b} = \log\left(q^{1/2} + q^{-1/2}E_{n,1}\right), \quad E_{n,a} = \sum_k q^{-ak}p_{n,k},$$
$$A_n^{K_b^2} = \log\left(1 + (1+q)E_{n,1} + q^{-1}E_{n,2}\right).$$

We recognize $E_{n,a}$ as the expectation value of the function q^{-ak} with respect to the probability distribution $p_{n,k}$.

As explained in *e.g.* [23], $w_{n,k}$ for fixed *n* reaches its maximum value at k/n = 5/6. As *n* increases, the ratio of the maximum value to generic values becomes exponentially large. So, as in a typical problem in statistical mechanics, we can approximate the probability distribution as Gaussian.

At the classical (s = 0) level of the conjecture, the only information we need is the position of the peak of the Gaussian:

(91)
$$E_{n,1} \approx q^{-k}|_{(k/n=5/6)} = e^{-5\pi i/3}, \quad E_{n,2} \approx q^{-2k}|_{(k/n=5/6)} = e^{-4\pi i/3}.$$

It gives the leading contribution to A_n , which agrees with the prediction of the state integral model:

(92)
$$\lim_{n \to \infty} A_n^{K=b} = \log(1 + e^{-5\pi i/3}) = \frac{1}{6} (3\log 3 + \pi i).$$
$$\lim_{n \to \infty} A_n^{K=b^2} = \log(1 + 2e^{-5\pi i/3} + e^{-4\pi i/3}) = \frac{1}{3} (3\log 3 + \pi i)$$

To proceed to the 1-loop (s = 1) level, we need to improve upon the Gaussian approximation. To prepare for a continuum limit, we define

(93)
$$x := \frac{k}{n}, \quad 0 \le x \le 1.$$

A key step is to apply the Euler-Maclaurin formula, including the leading correction term, to approximate $w_{n,k}$ by a continuous function,

(94)
$$\frac{w_{n,k}}{w_{n,n/2}} \approx \widetilde{w}_n(x) := \exp\left(2n \int_{1/2}^x \log[2\sin(\pi y)] dy + \log\left[\frac{\sin(\pi x)}{\sin(\pi/2)}\right]\right).$$

Then, the quantity $E_{n,a}$ can be approximated by the ratio of two integrals,

(95)
$$E_{n,a} \approx \left(\int_0^1 \widetilde{w}_n(x)dx\right)^{-1} \int_0^1 e^{-2\pi i a x} \widetilde{w}_n(x)dx$$

To evaluate these integrals up to the (1/n) order, we change the variable

(96)
$$x = \frac{5}{6} + \frac{1}{\sqrt{n}}t,$$

and expand the exponents to relevant orders,

(97)
$$\widetilde{w}_{n}(x) \approx \exp\left(-\sqrt{3}\pi t^{2} - \frac{3\sqrt{3}\pi t + 4\pi^{2}t^{3}}{3\sqrt{n}}\right),$$
$$e^{-2\pi iax} \approx e^{-5\pi ia/6} \exp\left(-\frac{2\pi ia}{\sqrt{n}}t\right).$$

Treating the $(1/\sqrt{n})$ terms in the exponents perturbatively, we find

(98)
$$A_n^{K_b} = \frac{1}{6} (3\log 3 + \pi i) + \frac{6 + i\sqrt{3}}{18} \left(\frac{2\pi i}{n}\right) + O\left(\frac{1}{n}\right)^2,$$
$$A_n^{K_b^2} = \frac{1}{3} (3\log 3 + \pi i) + \frac{2 + i\sqrt{3}}{3} \left(\frac{2\pi i}{n}\right) + O\left(\frac{1}{n}\right)^2,$$

in perfect agreement with (82) to the one-loop order (s = 1) up to an additive factor of $i\pi$ in $A_n^{K_b}$ which is due to the (-1) in our definition of (88).

4.3. Test of the length conjecture - numerical check. For a numerical check below, we define two sequences from the truncated invariants,

(99)
$$B_{n}^{K} := \begin{cases} \sum_{s=0}^{2} \mathcal{W}_{s}(\hat{O}_{K}; S^{3} \setminus \mathbf{4}_{1}) \hbar^{s} \Big|_{q = \exp(\frac{2\pi i}{n})} - i\pi, \quad K = K_{b} \\ \sum_{s=0}^{2} \mathcal{W}_{s}(\hat{O}_{K}; S^{3} \setminus \mathbf{4}_{1}) \hbar^{s} \Big|_{q = \exp(\frac{2\pi i}{n})}, \qquad K = K_{b}^{2}. \end{cases}$$

On the Jones-polynomial side, we introduce two sequences,

(100)
$$\begin{aligned} A_n^{K_b} &:= \log \frac{(-1) \times J_{n,\tilde{n}=2}(\mathbf{4}_1 \cup K_b; q)}{J_n(\mathbf{4}_1; q)} \Big|_{q=\exp(\frac{2\pi i}{n})}, \\ & (100) \qquad A_n^{K_b^2} := \log \frac{J_{n,\tilde{n}=2}(\mathbf{4}_1 \cup K_b^2; q)}{J_n(\mathbf{4}_1; q)} \Big|_{q=\exp(\frac{2\pi i}{n})} \\ & = \log \frac{J_{n,\tilde{n}=3}(\mathbf{4}_1 \cup K_b; q) + J_n(\mathbf{4}_1; q)}{J_n(\mathbf{4}_1; q)} \Big|_{q=\exp(\frac{2\pi i}{n})} \end{aligned}$$

The numerical values of A_n^K up to n = 30 are given in Table 1.

The numerical data for A_n is plotted against the continuous graph of B_n in Figure 7. The two plots appear to converge as n increases. To estimate the error more precisely, we define the two-loop error coefficients as

(101)
$$C_n^K = n^3 |A_n^K - B_n^K|, \quad K_b \text{ or } K_b^2.$$

Their numerical values are also given in Table 1. After some fluctuations for small n, say $n \leq 15$, they begin to converge slowly to a fixed O(1) constant. This tendency suggests that the error $|A_n - B_n|$ is approximately $O(1/n)^3$. So, as far as the numerical experiment is concerned, we have confirmed that the all order length conjecture holds up to two-loop order (s = 2).

5. 3D INDEX WITH QUANTUM TRACE MAP

We discuss the 3D index with quantum trace map, which we call $\mathcal{I}_{M+\hat{O}}(m,e;q)$ and $\mathcal{I}_{M_{P/Q}+\hat{O}}(q)$, extending the previous works [7, 14, 13].

The 3D index for a single tetrahedron is

(102)
$$\mathcal{I}_{\Delta}(m,e) := \sum_{n=[e]}^{\infty} \frac{(-1)^n q^{\frac{1}{2}n(n+1)-(n+\frac{1}{2}e)}}{(q)_n(q)_{n+e}} ,$$

where $[e] := \frac{1}{2}(|e|-e)$ and $(q)_n := (1-q)(1-q^2)\dots(1-q^n) .$

n	$A_n^{K_b}$	$C_n^{K_b}$	$A_n^{K_b^2}$	$C_n^{K_b^2}$
2	0.4700 + 1.5708i	3.10		
3	0.4345 + 1.1336i	2.74	-0.6844 + 0.9147i	35.4
4	0.4321 + 0.9545i	2.34	-0.4575 + 1.2120i	46.8
5	0.4403 + 0.8650i	3.67	-0.1391 + 1.4199i	49.7
6	0.4515 + 0.8126i	5.42	0.1251 + 1.4814i	49.8
7	0.4627 + 0.7774i	6.71	0.3133 + 1.4831i	48.6
8	0.4725 + 0.7512i	7.38	0.4469 + 1.4640i	46.5
9	0.4808 + 0.7304i	7.56	0.5444 + 1.4387i	44.1
10	0.4877 + 0.7131i	7.41	0.6176 + 1.4126i	41.5
11	0.4935 + 0.6983i	7.06	0.6742 + 1.3878i	39.0
12	0.4983 + 0.6856i	6.63	0.7190 + 1.3651i	36.8
13	0.5023 + 0.6744i	6.19	0.7552 + 1.3446i	34.9
14	0.5058 + 0.6647i	5.78	0.7851 + 1.3261i	33.2
15	0.5087 + 0.6560i	5.42	0.8102 + 1.3096i	31.8
16	0.5113 + 0.6483i	5.11	0.8316 + 1.2947i	30.6
17	0.5136 + 0.6414i	4.85	0.8499 + 1.2813i	29.7
18	0.5156 + 0.6353i	4.63	0.8659 + 1.2693i	28.9
19	0.5174 + 0.6297i	4.46	0.8800 + 1.2583i	28.2
20	0.5190 + 0.6246i	4.31	0.8924 + 1.2484i	27.7
21	0.5205 + 0.6200i	4.19	0.9035 + 1.2393i	27.2
22	0.5218 + 0.6158i	4.09	0.9135 + 1.2310i	26.8
23	0.5230 + 0.6119i	4.01	0.9225 + 1.2233i	26.4
24	0.5241 + 0.6084i	3.94	0.9306 + 1.2163i	26.1
25	0.5251 + 0.6051i	3.88	0.9381 + 1.2098i	25.9
26	0.5260 + 0.6021i	3.83	0.9449 + 1.2038i	25.6
27	0.5269 + 0.5992i	3.78	0.9511 + 1.1982i	25.4
28	0.5277 + 0.5966i	3.74	0.9569 + 1.1929i	25.2
29	0.5284 + 0.5942i	3.71	0.9622 + 1.1881i	25.1
30	0.5291 + 0.5919i	3.67	0.9672 + 1.1835i	24.9
∞	$\frac{1}{2}\log 3 + \frac{\pi i}{6}$		$\log 3 + \frac{\pi i}{3}$	
	$\approx 0.5493 + 0.5236i$	O(1)	$\approx 1.0986 + 1.0472i$	O(1)

TABLE 1. Numerical test of the length conjecture for the $\mathbf{4}_1$ knot with K_b and K_b^2 .

The 3D index for $M=S^3\backslash \mathcal{K}$ with an insertion on quantum loop operator

(103)
$$\hat{O} = \sum_{\alpha} C_{\alpha}(q^{1/4}) \prod_{i=1}^{k} \hat{z}_{i}^{a_{i}^{(\alpha)}}(\hat{z}_{i}^{\prime\prime})^{b_{i}^{(\alpha)}} \in \hat{\mathbb{C}}_{q}[\mathcal{T}]$$



FIGURE 7. Numerical test of the length conjecture for K_b (above) and K_b^2 (below). The lower/upper data on the same plot represent the real/imaginary parts, respectively. The dots are the numerical data from A_n^K , the curves are from the continuous function B_n^K .

is given as

$$\begin{aligned} &(104) \\ &\mathcal{I}_{M+\hat{O}}(m,e;q) \\ &= \sum_{\alpha} C_{\alpha} \left(\sum_{\{m\}} (-q^{-1/2})^{\langle \nu,\gamma \rangle} \prod_{i=1}^{k} \mathcal{I}_{\Delta}(m_{i} + b_{i}^{(\alpha)}, e_{i} - a_{i}^{(\alpha)}) q^{[a,b;m,e]_{i}^{(\alpha)}} \bigg|_{\gamma=g^{-1}\cdot\tilde{\gamma}} \right) ,\\ &\text{with} \quad \{m\} = \{m_{c_{1}}, \dots, m_{c_{k-1}}\} ,\\ &\gamma := (m_{1}, \dots, m_{k}, e_{1}, \dots, e_{k})^{T} , \quad \tilde{\gamma} = (m, m_{c_{1}}, \dots, m_{c_{k-1}}, e, 0, \dots, 0)^{T} ,\\ &\langle \nu, \gamma \rangle := \sum_{i=1}^{k} ((\nu_{p})_{i}m_{i} - \nu_{i}e_{i}) , \quad [a,b;m,e]_{i}^{(\alpha)} = \frac{1}{2} \left(a_{i}^{(\alpha)}m_{i} + b_{i}^{(\alpha)}e_{i} - b_{i}^{(\alpha)}a_{i}^{(\alpha)} \right) \end{aligned}$$

Here $g \in Sp(2k, \mathbb{Z})$ and $(\vec{\nu}, \vec{\nu}_p) \in \mathbb{Z}^k$ are determined by gluing equations of an ideal triangulation \mathcal{T} of M as given in (34). Under the Dehn filling, $M \to M_{P/Q}$, with a slope P/Q, the 3D index becomes [13]

$$\mathcal{L}_{M_{P/Q}+\hat{O}} = \sum_{(m,e)\in\mathbb{Z}^2} \frac{1}{2} (-1)^{Rm+2Se} \left(\delta_{\frac{P}{2}m+Qe,0} (q^{\frac{Rm+2Se}{2}} + q^{-\frac{Rm+2Se}{2}}) - \delta_{\frac{P}{2}m+Qe,-1} - \delta_{\frac{P}{2}m+Qe,1} \right) \mathcal{I}_{M+\hat{O}}(m,e) ,$$

Here two integers (R, S) are chosen to satisfy

 τ

 τ

(106)
$$\begin{pmatrix} R & S \\ P & Q \end{pmatrix} \in SL(2,\mathbb{Z}) .$$

For given (P, Q), two integers (R, S) satisfying the $SL(2, \mathbb{Z})$ relation is not unique but the choice of them does not affect the 3D index.

Conjecture 5.1. The quantum trace map $\operatorname{tr}_q^{\mathcal{T}}$ in the conjecture 2.1 satisfies the following properties

For any even link $K \subset M$,

(107)
$$\begin{aligned} \mathcal{I}_{M+\hat{O}_{K}}(m,e;q) &= \mathcal{I}_{M+\hat{O}}(-m,-e;q) \text{ and} \\ \mathcal{I}_{M_{P/Q}+\hat{O}_{K}}(m,e;q) &= 0 , \text{ if } M_{P/Q} \text{ is a Lens space }. \end{aligned}$$

Not all $\hat{O} \in \hat{\mathbb{C}}_q[\mathcal{T}]$ satisfy the above constraints and it will provide a guideline for determining the quantum trace map.

Example 5.2. As an example, consider the case when $M = S^3 \setminus 4_1$. The 3D index $\mathcal{I}_M(m, e; q)$ is

(108)
$$\mathcal{I}_{M=S^{3}\backslash \mathbf{4}_{1}}(m,e;q) = \sum_{e_{c}\in\mathbb{Z}}\mathcal{I}_{\Delta}(e_{c}-e,e_{c})\mathcal{I}_{\Delta}(e_{c}-m,e_{c}-e-m) .$$

With an insertion of $\hat{O} = \hat{y}^{\alpha} \hat{z}^{\beta} (\hat{y}'')^{\gamma} (\hat{z}'')^{\delta} \in \hat{\mathbb{C}}_q[\mathcal{T}]$, the index becomes (109)

$$\mathcal{I}_{S^{3}\backslash \mathbf{4}_{1}+\hat{O}}(m,e;q) = \sum_{e_{c}\in\mathbb{Z}}q^{\frac{\alpha(e_{c}-e)+\beta(e_{c}-m)+\gamma e_{c}+\delta(e_{c}-e-m)-\alpha\gamma-\beta\delta}{2}} \times \mathcal{I}_{\Delta}(e_{c}-e+\gamma,e_{c}-\alpha)\mathcal{I}_{\Delta}(e_{c}-m+\delta,e_{c}-e-m-\beta) .$$

Using the above expression, one can check that all $\hat{O} \in \operatorname{tr}_q^{\mathcal{T}}(\mathcal{S}_q^{\operatorname{even}}[S^3 \setminus \mathbf{4}_1])$ in (81) satisfy the conditions in the conjecture. For $\hat{O} = C_1 \hat{y}^{-1} \hat{z}'' + C_2 \hat{z}^{-1} \hat{y}'' + C_3 \hat{y}^{-1} \hat{z}^{-1} \in \hat{\mathbb{C}}_q[\mathcal{T}]$, as an example, the 3D index becomes

(110)
$$\mathcal{I}_{S^{3}\backslash \mathbf{4}_{1}+\hat{O}}(m,e) = \sum_{e_{c}\in\mathbb{Z}} \left[C_{1}\mathcal{I}_{\Delta}(e_{c}-e,e_{c}+1)\mathcal{I}_{\Delta}(e_{c}-m+1,e_{c}-e-m)q^{-\frac{m}{2}} + C_{2}\mathcal{I}_{\Delta}(e_{c}-e,e_{c}+1,e_{c})\mathcal{I}_{\Delta}(e_{c}-m,e_{c}-e-m+1)q^{\frac{m}{2}} + C_{3}\mathcal{I}_{\Delta}(e_{c}-e,e_{c}+1)\mathcal{I}_{\Delta}(e_{c}-m,e_{c}-e-m+1)q^{\frac{m-2e_{c}+e}{2}} \right],$$

which implies (to avoid clutter, we suppress the subscript in $\mathcal{I}_{S^3\backslash \mathbf{4}_1+\hat{O}})$

$$\begin{split} \mathcal{I}(0,0) &= C_3 + (-2C_1 - 2C_2 + C_3)q - 2(C_1 + C_2 + C_3)q^2 \\ &+ (2C_1 + 2C_2 - 5C_3)q^3 + \dots, \\ \mathcal{I}(1,0) &= (C_3 - C_1)q^{1/2} - (C_1 + C_2 + C_3)q^{3/2} + (2C_1 - C_2 - 3C_3)q^{5/2} + \dots, \\ \mathcal{I}(-1,0) &= (C_3 - C_2)q^{1/2} - (C_1 + C_2 + C_3)q^{3/2} + (2C_2 - C_1 - 3C_3)q^{5/2} + \dots, \\ \mathcal{I}(1,1) &= (C_3 - C_1) - C_3q + (C_1 - 2C_2 - 2C_3)q^2 + (4C_1 - 3C_3)q^3 + \dots, \\ \mathcal{I}(-1,-1) &= (C_3 - C_2) - C_3q + (C_2 - 2C_1 - 2C_3)q^2 + (4C_2 - 3C_3)q^3 + \dots. \end{split}$$

For the index to be invariant under the Weyl \mathbb{Z}_2 symmetry, $(m, e) \leftrightarrow (-m, -e)$, in the conjecture, the quantum loop operator $\hat{O} = C_1 \hat{y}^{-1} \hat{z}'' + C_2 \hat{z}^{-1} \hat{y}'' + C_3 \hat{y}^{-1} \hat{z}^{-1}$ should satisfy

(111)
$$C_1 = C_2$$
.

So we expect that only the \hat{O} with $C_1 = C_2$ can appear as image of the quantum trace map. For example, from (81),

(112)

$$\operatorname{tr}_{q}^{\mathcal{T}}(K_{b}) = (\hat{y}^{-1} + \hat{z}^{-1} - \hat{y}^{-1}\hat{z}^{-1}) \\ \sim (\hat{y}^{-1}(\hat{z}'' + \hat{z}^{-1}) + \hat{z}^{-1}(1 - \hat{y}^{-1})) \\ \sim (\hat{y}^{-1}(\hat{z}'' + \hat{z}^{-1}) + \hat{z}^{-1}\hat{y}'') = \hat{O}|_{C_{1}=C_{2}=C_{3}=-1}$$

satisfies the condition.

6. QUANTUM TRACE MAP IN COMPLEX CHERN-SIMONS THEORY

In this section, we give a brief review on general aspect of $SL(2, \mathbb{C})$ Chern-Simons theory and its relation the volume conjecture. Refer to [16, 8, 34] for details. By interpreting the quantum trace operator as holomorphic Wilson loop in the complex Chern-Simons theory, we give a physical derivation of the length conjecture.

The action of $SL(2,\mathbb{C})_{(k,\sigma)}$ Chern-Simons theory is given as follows

$$S_{(k,\sigma)}[\mathcal{A},\overline{\mathcal{A}};M] = \frac{k+\sigma}{8\pi} \int \operatorname{tr}\left(\mathcal{A} \wedge d\mathcal{A} + \frac{2}{3}\mathcal{A}^3\right) + \frac{k-\sigma}{8\pi} \int \operatorname{tr}\left(\overline{\mathcal{A}} \wedge d\overline{\mathcal{A}} + \frac{2}{3}\overline{\mathcal{A}}^3\right) \,.$$

The $k \in \mathbb{Z}$ should be quantized for gauge invariance while σ can be either real $(\sigma \in \mathbb{R})$ or purely imaginary $(\sigma \in i\mathbb{R})$. The state-integral and the 3D index for a knot complement $M = S^3 \setminus \mathcal{K}$ compute the partition function Z of the complex Chern-Simons theory with k = 1 and k = 0 respectively [7, 6]:

$$\begin{aligned} Z_{\hbar=2\pi ib^2}(M; X_{\mathbf{m}}) &= Z \left[SL(2, \mathbb{C})_{k,\sigma} \text{ CS theory on } M \text{ with } k = 1, \sigma = \frac{1-b^2}{1+b^2} \right] \\ &= \int \frac{[D\mathcal{A}][D\overline{\mathcal{A}}]}{(\text{gauge})} \Big|_{b.c} \exp\left(iS_{(k,\sigma)}[\mathcal{A}, \overline{\mathcal{A}}; M]\right) \Big|_{k=1,\sigma=\frac{1-b^2}{1+b^2}}, \\ \mathcal{I}_M^{\text{fugacity}}(m, u; q) &= Z \left[SL(2, \mathbb{C})_{k,\sigma} \text{ CS theory on } M \text{ with } k = 0, \sigma = \frac{4\pi i}{\log q} \right] \\ &= \int \frac{[D\mathcal{A}][D\overline{\mathcal{A}}]}{(\text{gauge})} \Big|_{b.c} \exp\left(iS_{(k,\sigma)}[\mathcal{A}, \overline{\mathcal{A}}; M]\right) \Big|_{k=0,\sigma=\frac{4\pi i}{\log q}}. \end{aligned}$$

24 PRARIT AGARWAL, DONGMIN GANG, SANGMIN LEE, AND MAURICIO ROMO

Here $\mathcal{I}_{M}^{\text{fugacity}}(m, u; q)$ is the 3D index in fugacity basis

(114)
$$\mathcal{I}_{M}^{\text{fugacity}}(m, u; q) = \sum_{e \in \mathbb{Z}} \mathcal{I}_{M}(m, e; q) u^{e}$$

The boundary conditions are

$$P \exp\left(\oint_{O_{\mathrm{m}}} \mathcal{A}\right) = \begin{pmatrix} e^{X/2} & *\\ 0 & e^{-X/2} \end{pmatrix}, \quad P \exp\left(\oint_{O_{\mathrm{m}}} \overline{\mathcal{A}}\right) = \begin{pmatrix} e^{\tilde{X}/2} & *\\ 0 & e^{-\tilde{X}/2} \end{pmatrix}$$

(115) where

$$(e^{X}, e^{\tilde{X}}) = \begin{cases} (e^{X_{\mathbf{m}}}, e^{\frac{X_{\mathbf{m}}}{b^{2}}}) & \text{for } Z_{\hbar=2\pi i b^{2}}(M; X_{\mathbf{m}}) \\ (q^{\frac{m}{2}}u, q^{\frac{m}{2}}u^{-1}) & \text{for } \mathcal{I}_{M}^{\text{fugacity}}(m, u; q) \end{cases}$$

The partition functions can be written in the following factorization form [1] (116)

$$Z_{\hbar=2\pi i b^2}(M; X_{\mathbf{m}}) = \frac{1}{2} \sum_{\rho} B_M^{\rho}(q; e^X) B_M^{\rho}(\tilde{q}; e^{\tilde{X}}) \text{ with } (q, \tilde{q}) = (e^{2\pi i b^2}, e^{\frac{2\pi i}{b^2}}) ,$$
$$\mathcal{I}^{\text{fugacity}}(m, u; q) = \frac{1}{2} \sum_{\rho} B_M^{\rho}(q; X) B_M^{\rho}(\tilde{q}; \tilde{X}) \text{ with } (q, \tilde{q}) = (q, q^{-1}) .$$

 $(e^X, e^{\tilde{X}})$ here are identical to that of (115). $B^{\rho}_M(q; e^X)$ is so-called holomorphic block labelled by an (adj)-irreducible $SL(2, \mathbb{C})$ flat connection \mathcal{A}_{ρ} satisfying the boundary condition in (115).¹ It can be defined as following path-integral

(117)
$$B_M^{\rho}(q; e^X) = \int_{\Gamma^{\rho}} [D\mathcal{A}] e^{-\frac{1}{2\hbar} \int_M \operatorname{tr}(\mathcal{A} \wedge d\mathcal{A} + \frac{2}{3}\mathcal{A}^3)} \text{ with } q = e^{\hbar}.$$

Here Γ^{ρ} denotes the Lefschetz thimble in the functional space of \mathcal{A} associated with the flat connection \mathcal{A}_{ρ} [34].

The $Z_{\hbar}(M, X_{\mathbf{m}})$ and $B_M(q = e^{\hbar}; e^X)$ have the same asymptotic expansion in the limit $\hbar \to 0$:

(118)
$$B_M^{\rho_\alpha}(q = e^{\hbar}; e^X) \xrightarrow{\hbar \to 0} \exp\left(\sum_{n=0}^{\infty} S_n^{(\alpha)}(M)\hbar^{n-1}\right)\Big|_{X_{\mathbf{m}}=X}$$

where the $S_n^{(\alpha)}$ is the perturbative expansion of the state-integral around the saddle point $\vec{Z}^{(\alpha)}$ as defined in (43). For each irreducible flat connection ρ_{α} , there is an associated saddle point $\vec{Z}^{(\alpha)}$ in the state-integral model. More precisely, the map is generically in 2-to-1 since two flat connections ρ_{α} and $\rho_{\bar{\alpha}}$ related to each other by the action of $H^1(M, \mathbb{Z}_2) = \mathbb{Z}_2$ correspond to the same saddle point $\vec{Z}^{(\alpha)}$. In the asymptotic limit $\hbar \to 0$, the $B_M^{\rho}(\tilde{q}; e^{\tilde{X}})$ becomes trivially 1 since $\tilde{q} = e^{-\frac{4\pi^2}{\hbar}}$ goes to 0 in the limit and $B^{(\rho)}(q)$ becomes an infinite power series in $q^{1/2}$ starting with $1 + O(q^{1/2})$ as $q \to 0$.

The quantum trace map \hat{O}_K for a knot $K \subset M$ corresponds to the following holomorphic Wilson loop in the CS theory

(119)
$$W_K(\mathcal{A}) = (-1) \times \operatorname{Tr}\left(P \exp\left(\oint_K \mathcal{A}\right)\right)$$

¹A $SL(2, \mathbb{C})$ flat connection $\rho \in \text{Hom}[\pi_M \to SL(2, \mathbb{C})]$ can be alternatively described by a $SL(2, \mathbb{C})$ gauge connection \mathcal{A}_{ρ} with vanishing curvature, i.e. $d\mathcal{A}_{\rho} + \mathcal{A}_{\rho} \wedge \mathcal{A}_{\rho} = 0$.

and the $Z_{\hbar}(M + \hat{O}_K)$ and $\mathcal{I}_{M+\hat{O}_K}$ are given by following path-integrals (120)

$$Z_{\hbar}(M+\hat{O}_{K};X_{\mathbf{m}}) = (-1) \times \int \frac{[D\mathcal{A}][D\overline{\mathcal{A}}]}{(\text{gauge})} e^{iS_{(k,\sigma)}[\mathcal{A},\overline{\mathcal{A}};M]} W_{K}(\mathcal{A}) \Big|_{k=1,\sigma=\frac{1-b^{2}}{1+b^{2}}} ,$$
$$\mathcal{I}_{M+\hat{O}_{K}}^{\text{fugacity}}(m,u;q) = (-1) \times \int \frac{[D\mathcal{A}][D\overline{\mathcal{A}}]}{(\text{gauge})} e^{iS_{(k,\sigma)}[\mathcal{A},\overline{\mathcal{A}};M]} W_{K}(\mathcal{A}) \Big|_{k=0,\sigma=\frac{4\pi i}{\log q}} .$$

With the insertion of loop operator, we expect that factorization should be modified as follows

(121)
$$Z_{\hbar}(M+\hat{O};X_{\mathbf{m}}) = \frac{1}{2} \sum_{\rho} B^{\rho}_{M+\hat{O}}(q;e^{X}) B^{\rho}_{M}(\tilde{q};e^{\tilde{X}}) ,$$
$$\mathcal{I}^{\text{fugacity}}(m,u;q) = \frac{1}{2} \sum_{\rho} B^{\rho}_{M+\hat{O}}(q;e^{X}) B^{\rho}_{M}(\tilde{q};e^{\tilde{X}}) .$$

Here we define

(122)
$$B^{\rho}_{M+\hat{O}_{K}}(q;X) = (-1) \times \int_{\Gamma^{\rho}} [D\mathcal{A}] e^{\frac{i}{2\hbar} \int_{M} \operatorname{tr}(\mathcal{A} \wedge d\mathcal{A} + \frac{2}{3}\mathcal{A}^{3})} W_{K}(\mathcal{A})$$

In the same argument used in (118), the $Z_{\hbar}(M + \hat{O}_K, X_{\mathbf{m}})$ and $B_{M + \hat{O}_K}(q = e^{\hbar}; e^X)$ is expected to share the same asymptotic expansion in the limit $\hbar \to 0$:

(123)
$$B_{M+\hat{O}_{K}}^{\rho_{\alpha}}(q=e^{\hbar};e^{X}) \xrightarrow{\hbar \to 0} \exp\left(\sum_{n=0}^{\infty} S_{n}^{(\alpha)}(\hat{O}_{K};M)\hbar^{n-1}\right)\Big|_{X_{\mathbf{m}}=X}.$$

where the $S_n^{(\alpha)}$ is the perturbative expansion of the state-integral around the saddle point $\vec{Z}^{(\alpha)}$ as defined in (43).

Volume conjecture for a hyperbolic knot \mathcal{K} relates an asymptotic expansion of Jones polynomial to an asymptotic expansion of the $B^{(hyp)}_{S^3\setminus\mathcal{K}}(q,X)$ as follows

$$\begin{pmatrix} \text{asymptotic expansion of } \frac{J_n\left(\mathcal{K}; q = \exp(\frac{2\pi i}{k})\right)}{J_n\left(\bigcirc; q = \exp(\frac{2\pi i}{k})\right)} \text{ in the limit } k = n \to \infty \end{pmatrix}$$

$$(124) \simeq \left(\text{asymptotic expansion of } B^{\rho(\text{hyp})}_{S^3 \setminus \mathcal{K}}(q = e^{\hbar}, X = 0) \text{ in the limit } \hbar \to 0 \right),$$

$$\text{ with an identification } \hbar = \frac{2\pi i}{k}.$$

Here \bigcirc denote the unknot in S^3 and \simeq means the same asymptotic expansion modulo an overall factor of the form $\frac{C}{k^{3/2}}$. The equivalence of two asymptotic expansions can be understood by comparing the path-integral of the Jones polynomial in (150) and the B^{ρ}_{M} in (117) modulo some subtleties. The most subtle part is that how the path-integral of SU(2) gauge field a can be related to the path-integral of $SL(2,\mathbb{C}) = SU(2)_{\mathbb{C}}$ gauge field \mathcal{A} along the $\Gamma^{(\text{hyp})}$ in the asymptotic limit. To understand it, let us consider a finite dimensional integral $\int_{\mathbb{R}} dx e^{ik(\frac{1}{3}x^3+x)}$ in the limit $k \to \infty$ as an analogy. Although it is integration over real axis, it gets contribution from the saddle point at x = -i in the asymptotic limit. So, the appearance of complex flat connection in the asymptotic expansion of SU(2) Chern-Simons theory is not so strange. But unlike in the finite-dimensional analogy, the contribution from the flat connection $\rho_{(\text{hyp})}$ is exponentially growing in the asymptotic limit. This issue has been addressed in [34] and we will not address the subtle issue here.

26 PRARIT AGARWAL, DONGMIN GANG, SANGMIN LEE, AND MAURICIO ROMO

Our length conjecture can be understood in a similar way. For that, we consider the Jones polynomial $J_{n,\tilde{n}=2}(\mathcal{K}\cup K;q)$. We consider the case when the K is a knot, i.e. $\sharp(K) = 1$, for simplicity. The Jones polynomial can be represented by following path-integral

(125)

$$J_{n,\tilde{n}=2}\left(\mathcal{K}\cup K;q=\exp(\frac{2\pi i}{k})\right)$$

$$=(-1)\times\frac{\int \frac{[da]}{(\text{gauge})}\Big|_{\text{b.c. in (149)}}\exp\left(\frac{i(k-2)}{4\pi}\int_{S^3\setminus\mathcal{K}}\operatorname{Tr}(a\wedge da-\frac{2}{3}a^3)\right)W_K(a)}{\int \frac{[da]}{(\text{gauge})}\exp\left(\frac{i(k-2)}{4\pi}\int_{S^3}\operatorname{Tr}(a\wedge da-\frac{2}{3}a^3)\right)}$$

Then, from comparison with the path-integral in (122), one naturally expects that (126)

$$\left(\begin{array}{l} \text{asymptotic expansion of } \frac{J_{n,\tilde{n}=2} \left(\mathcal{K} \cup K; q = \exp(\frac{2\pi i}{k}) \right)}{J_n \left(\bigcirc; q = \exp(\frac{2\pi i}{k}) \right)} \text{ in the limit } k = n \to \infty \right) \\ \simeq \left(\begin{array}{l} \text{asymptotic expansion of } B_{S^3 \setminus \mathcal{K} + \hat{O}_K}^{\rho_{(\text{hyp})}}(q = e^\hbar, X = 0) \text{ in the limit } \hbar \to 0 \right), \\ 2\pi i \end{array} \right)$$

with an identification $\hbar = \frac{2\pi i}{k}$.

Combining (124), (126) with (118), (123), one derives the length conjecture in (18) where the $\mathcal{Z}_{s}^{(\mathrm{hyp})}(\hat{O}_{K}; M)$ is given in (44). Proposed properties of $\mathcal{I}_{M+\hat{O}}(m, e; q)$ or $\mathcal{I}_{M_{P/Q}+\hat{O}}$ in (107) simply follows from

Proposed properties of $\mathcal{I}_{M+\hat{O}}(m, e; q)$ or $\mathcal{I}_{M_{P/Q}+\hat{O}}$ in (107) simply follows from the path-integral expression (113) or the factorization (116). The invariance of the 3D index under $(m, e) \leftrightarrow (-m, -e)$ comes from the Weyl \mathbb{Z}_2 invariance, $X_{\mathbf{m}} \leftrightarrow$ $-X_{\mathbf{m}}$, in the path-integral. For $M_{P/Q}$ = (Lens space), there is no irreducible flat connection ρ on $M_{P/Q}$ and thus we expect the index vanishes from the factorization.

Acknowledgements

The work of DG is supported in part by the National Research Foundation of Korea grant NRF-2021R1G1A1095318 and by Creative-Pioneering Researchers Program through Seoul National University. The work of SL is supported in part by the National Research Foundation of Korea grant NRF-2019R1A2C2084608. MR acknowledges support from the National Key Research and Development Program of China, grant No. 2020YFA0713000, and the Research Fund for International Young Scientists, NSFC grant No. 11950410500.

APPENDIX A. GLUING EQUATIONS OF AN IDEAL TRIANGULATION

From an ideal triangulation \mathcal{T} of a knot complement $M = S^3 \setminus \mathcal{K}$, we have following gluing equations for the logarithmic edge parameters, $(e^Z, e^{Z'}, e^{Z''}) = (z, z'', z'')$, of tetrahedra (refer to, e.g., [9] for details)

$$C_{I=1,...,|\mathcal{T}|} = \sum_{i=1}^{|\mathcal{T}|} F_{Ii}Z_i + G_{Ii}Z'_i + H_{Ii}Z''_i - (2\pi i + \hbar) \text{ (internal edges) },$$
(127)
$$\mathfrak{M} = \sum_{i=1}^{|\mathcal{T}|} U_iZ_i + V_iZ'_i + W_iZ''_i \text{ (meridian) },$$

$$\mathfrak{L} = \sum_{i=1}^{|\mathcal{T}|} P_iZ_i + Q_iZ'_i + R_iZ''_i \text{ (longitude) }.$$

Here $F_{Ii}, G_{Ii}, H_{Ii} \in \{0, 1, 2\}$ while $U_i, V_i, W_i, P_i, Q_i, R_i \in \mathbb{Z}$. The logarithmic edge parameters (Z_i, Z'_i, Z''_i) satisfy following linear relation

(128)
$$Z_i + Z'_i + Z''_i = i\pi + \frac{\hbar}{2} .$$

Then, the gluing equation variety $\chi[\mathcal{T}]$ is given by

(129)

$$\chi[\mathcal{T}] = \left\{ (z_i, z_i'')_{i=1}^{|\mathcal{T}|} : z_i^{-1} + z_i'' - 1 = 0, \\
e^{C_I} := \prod_{i=1}^{|\mathcal{T}|} (-1)^{G_{I_i}} z_i^{F_{I_i} - G_{I_i}} (z_i'')^{H_{I_i} - G_{I_i}} = 1 \right\}.$$

We define $e^{\hat{C}_I}$, $e^{\hat{\mathfrak{M}}}$ and $e^{-\hat{\mathfrak{M}}} \in \mathbb{C}[q^{\pm 1/4}][\hat{z}_i^{\pm 1}, \hat{z}_i''^{\pm 1}]/\langle \hat{z}_i \hat{z}_j'' - q^{-\delta i j} \hat{z}_j'' \hat{z}_i \rangle$ as

(130)
$$e^{\hat{C}_{I}} := q^{-1} \prod_{i=1}^{|\mathcal{T}|} (-q^{\frac{1}{2}})^{G_{Ii}} q^{\frac{1}{2}(F_{Ii} - G_{Ii})(H_{Ii} - G_{Ii})} \hat{z}_{i}^{F_{Ii} - G_{Ii}} (\hat{z}_{i}'')^{H_{Ii} - G_{Ii}} ,$$

$$e^{\hat{\mathfrak{M}}} := \prod_{i=1}^{|\mathcal{T}|} (-q^{\frac{1}{2}})^{Q_{i}} q^{\frac{1}{2}(P_{i} - Q_{i})(R_{i} - Q_{i})} \hat{z}_{i}^{P_{i} - Q_{i}} (\hat{z}_{i}'')^{R_{i} - Q_{i}} ,$$

$$e^{-\hat{\mathfrak{M}}} := \prod_{i=1}^{|\mathcal{T}|} (-q^{-\frac{1}{2}})^{Q_{i}} q^{\frac{1}{2}(P_{i} - Q_{i})(R_{i} - Q_{i})} \hat{z}_{i}^{Q_{i} - P_{i}} (\hat{z}_{i}'')^{Q_{i} - R_{i}} .$$

They all mutually commute due to a symplectic structures in gluing equations [25]

(132)
$$[e^{\hat{C}_I}, e^{\hat{C}_J}] = [e^{\pm \hat{\mathfrak{M}}}, e^{\hat{C}_I}] = 0$$

As a concrete example, we collect some information on the figure-8 knot complement. The fundamental group of the knot-complement is (see Figure 6)

(133)
$$\pi_1(S^3 \setminus \mathbf{4}_1) = \langle a, b, c : ac^{-1}ba^{-1}c = bc^{-1}b^{-1}a = 1 \rangle.$$

It contains the peripheral subgroup $\mathbb{Z} \times \mathbb{Z}$ to be identified as the fundamental group of the boundary torus,

(134)
$$\pi_1(\partial(S^3\backslash \mathbf{4}_1)) = \pi_1(\mathbb{T}^2) = \mathbb{Z} \times \mathbb{Z} = \langle \mathbf{m}, \mathbf{l} \rangle \subset \pi_1(S^3\backslash \mathbf{4}_1).$$

We choose (\mathbf{m}, \mathbf{l}) canonically to be the meridian and longitude. Then the embedding $i : \pi_1(\partial(S^3 \backslash \mathbf{4}_1)) \to \pi_1(S^3 \backslash \mathbf{4}_1)$ is given by

(135)
$$i(\mathbf{m}) = a, \quad i(\mathbf{l}) = ac^{-1}bca^{-1}b^{-1}$$

The simplest ideal triangulation \mathcal{T} of the knot complement consists of two tetrahedra. The edge variables,

(136)
$$(y, y', y'') = (e^Y, e^{Y'}, e^{Y''}), \quad (z, z', z'') = (e^Z, e^{Z'}, e^{Z''}),$$

are subject to the usual conditions,

(137)
$$y + \frac{1}{y'} = 1, \quad y' + \frac{1}{y''} = 1, \quad y'' + \frac{1}{y} = 1, \quad yy'y'' = -1,$$

and similar ones for (z, z', z''). The two internal edges require that

(138)
$$C_1 = Y' + 2Y + Z' + 2Z - 2\pi i = 0, C_2 = Y' + 2Y'' + Z' + 2Z'' - 2\pi i = 0$$

which implies yz = y''z''. In summary, the gluing equation variety $\chi[\mathcal{T}]$ is (139) $\chi[\mathcal{T}] = \{(y, z, y'', z'') : y^{-1} + y'' - 1 = 0, z^{-1} + z'' - 1 = 0, yz = y''z''\}$. For each $(y, z, y'', z'') \in \chi[\mathcal{T}]$, one can assign a $PSL(2, \mathbb{C})$ -representation ρ as follows

(140)

$$\rho(a) = \begin{bmatrix} \frac{1}{\sqrt{y''z}} \begin{pmatrix} y'' & 0\\ -y^{-1} & z \end{pmatrix} \end{bmatrix},$$

$$\rho(b) = \begin{bmatrix} \frac{1}{\sqrt{y'z'}} \begin{pmatrix} z' & -z'\\ y'z' & y'(1-z') \end{pmatrix} \end{bmatrix},$$

$$\rho(c) = \begin{bmatrix} \frac{1}{\sqrt{yz''}} \begin{pmatrix} y-z^{-1} & 1-y\\ -z^{-1} & 1 \end{pmatrix} \end{bmatrix}.$$

Here we regard $PSL(2, \mathbb{C}) = SL(2, \mathbb{C})/\mathbb{Z}_2$ and [...] is the equivalence class under the $\mathbb{Z}_2 = \{\pm 1\}$. These are consistent with the two relations in (133) and the boundary holonomies deduced from the triangulation,

(141)

$$\rho(\mathbf{m}) = \rho(a) = \left[\begin{pmatrix} e^{\mathfrak{M}/2} & 0\\ * & e^{-\mathfrak{M}/2} \end{pmatrix} \right], \quad \mathfrak{M} = Y'' - Z,$$

$$\rho(\mathbf{l}) = \rho(ac^{-1}bca^{-1}b^{-1}) = \left[\begin{pmatrix} e^{\mathfrak{L}} & 0\\ * & e^{-\mathfrak{L}} \end{pmatrix} \right], \quad \mathfrak{L} = Z - Z''.$$

APPENDIX B. COLORED JONES POLYNOMIAL FOR FRAMED LINK

Let $K = \bigcup_{I=1}^{\sharp(K)} K_I$ be a framed link inside 3-sphere S^3 . Here we provide the definitions of colored Jones polynomial $J_{\vec{n}}(K;q)$ using the Kauffman bracket Skein module $S_q[S^3]$. The subscript $\vec{n} = (n_1, \ldots, n_{\sharp(K)})$ with $n_i \in \mathbb{Z}_{\geq 1}$ represents colors of each components of the link.

When $n_i = 2$ for all $i = 1, ..., \sharp(K)$, the invariant is called just Jones polynomial and will be simply denoted by J without subscript. The polynomial J(K;q) is defined as follows

(142)
$$J(K;q) := \frac{Y_K}{Y_{\emptyset}} \in \mathbb{Z}[q^{1/4}, q^{-1/4}].$$

Here Y_K is considered to be an element of $S_q[S^3]$. The definition makes sense since the Skein module $S_q[S^3]$ is one-dimensional module spanned by Y_{\emptyset} . The invariant is slightly different from the conventional Jones polynomial, $J^{\text{oriented}}(K;q)$, which is an invariant for oriented (unframed) link. The two invariants are related to each other in the following way

(143)
$$J^{\text{oriented}}(K;q) := (-1)^{\sharp(K)} (-q^{3/4})^{w(K)} J(K;q)$$

The sign factor is introduced in order for $J^{\text{oriented}}(K = \bigcirc^n; q)$ to be $(q^{1/2} + q^{-1/2})^n$. Here w(K) is the writhe of the framed link

(144)
$$w(K) := \sum_{p: \text{crossings in } K} \epsilon(p) , \quad \epsilon(p) := \begin{cases} +1, & p \text{ is a positive crossing} \\ -1, & p \text{ is a negative crossing} \end{cases}$$

The writhe is an oriented framed link invariant. Since the factor $(-q^{3/4})^{w(K)}$ exactly cancels the $(-q^{3/4})^{\pm 1}$ factor in Y_K under the 1st Reidemeister move, see Figure 2, the J^{oriented} is an oriented unframed invariant.

To generalize the definition to general colors \vec{n} , we introduce the symmetric k-product of a framed knot K

(145)
$$\begin{aligned} & \operatorname{Sym}^{\otimes k=0} K = \emptyset , \quad \operatorname{Sym}^{\otimes k=1} K = K , \quad \operatorname{Sym}^{\otimes k=2} K = K^2 - \emptyset;, \\ & \operatorname{Recursively}, \operatorname{Sym}^{\otimes k} K := (\operatorname{Sym}^{\otimes k-1} K) K - (\operatorname{Sym}^{\otimes k-2} K) . \end{aligned}$$

Here K^n denote the *n*-cabling of a framed knot K, a framed link obtained by displacing *n* parallel copies of *K* along the normal direction of framing. Then, we define the framed colored Jones polynomial for a link $K = \bigcup_{I=1}^{\sharp(K)} K_I$ as follows

(146)
$$J_{\vec{n}}(K;q) := J(\bigcup_{I=1}^{\sharp(K)} \operatorname{Sym}^{\otimes(n_I-1)} K_I;q) = \frac{Y_{\bigcup_{I=1}^{\sharp(K)} \operatorname{Sym}^{\otimes(n_I-1)} K_I}}{Y_{\emptyset}}$$

The framed colored Jones polynomial have following path-integral representation (k > 2) [33]

$$\begin{aligned} &(147)\\ J_{\vec{n}}\left(K;q=q=e^{2\pi i/k}\right)\\ &=\prod_{I=1}^{\sharp(K)}(-1)^{n_{I}-1}\frac{\int\frac{[da]}{(\text{gauge})}\exp\left(\frac{i(k-2)}{4\pi}\int_{S^{3}}\text{CS}(a)\right)\prod_{I=1}^{\sharp(K)}\text{tr}_{R_{n}}\left(P\exp(\oint_{K_{I}}a)\right)}{\int\frac{[da]}{(\text{gauge})}\exp\left(\frac{i(k-2)}{4\pi}\int_{S^{3}}\text{CS}(a)\right)} \text{ ,} \end{aligned}$$
where $\text{CS}(a)=\text{Tr}\left(a\wedge da-\frac{2}{3}a^{3}\right)$.

Here $\int \frac{[da]}{(\text{gauge})}$ means integral over gauge-equivalence classes of SU(2) connections on S^3 . $P \exp(\oint_K a)$ denotes the SU(2) holonomy matrix of SU(2) connection aalong the closed knot K. $R_n = \text{Sym}^{\otimes (n-1)} \mathbf{2} \in \text{Hom}[SU(2) \to SU(n)]$ is the *n*dimensional irreducible representation of SU(2) and we define

$$\operatorname{tr}_R(g) = \operatorname{Tr}(R(g))$$
.

Since $\operatorname{tr}_R(g) = \operatorname{tr}(g^{-1})$, the path-integral does not depend on the orientation choice of K_I . To regularize infinity coming from self-interaction of knots, on the other hand, one need to introduce framing and the quantum Wilson loop expectation value depends on the choice. In the above, we include following overall sign factor

(148)
$$\prod_{I=1}^{\sharp(K)} (-1)^{n_I - 1}$$

which reflects that $J_{K \sqcup \bigcirc} = -(q^{1/2} + q^{-1/2})J_K$ while an additional Wilson loop along the disjoint unknot \bigcirc in R = 2 contribute a multiplicative factor $(q^{1/2} + q^{-1/2})$ to the path-integral of Chern-Simons theory [33].

Wilson loop operator, $\operatorname{tr}_R(P\exp(\oint_K a))$, along a knot $K \subset S^3$ with color $R = \operatorname{Sym}^{\otimes (n-1)} \mathbf{2}$ in the SU(2) Chern-Simons theory can be alternatively described by monodromy defect in the knot complement $S^3 \setminus K$ defined by following boundary condition

(149) boundary condition :
$$P \exp\left(\oint_{\bigcirc_{\mathrm{m}}} a\right) = \begin{pmatrix} e^{\frac{\pi i n}{k}} & 0\\ 0 & e^{-\frac{\pi i n}{k}} \end{pmatrix}$$
.

Using the alternative definition of the Wilson loop operator, the Jones polynomial $J_n(\mathcal{K};q)$ can be given as

$$J_n\left(\mathcal{K}; q = e^{2\pi i/k}\right) = (-1)^{n-1} \frac{\int \frac{[da]}{(\text{gauge})} \bigg|_{\text{b.c. in (149)}} \exp\left(\frac{i(k-2)}{4\pi} \int_{S^3 \setminus \mathcal{K}} \text{CS}(a)\right)}{\int \frac{[da]}{(\text{gauge})} \exp\left(\frac{i(k-2)}{4\pi} \int_{S^3} \text{CS}(a)\right)}$$

Using the physical definition of the Jones polynomial, now let us explain a nontrivial property of the knot invariant summarized in Figure 8. In the Chern-Simons



FIGURE 8. The effect of adding the meridian knot. A relation between $Y_{(\text{Sym}^{\otimes (n-1)}\mathcal{K})\cup \bigcirc_{m}}$ and $Y_{(\text{Sym}^{\otimes (n-1)}\mathcal{K})}$ holds, where \bigcirc_{m} is a unknot linking the heavy knot \mathcal{K} .

path-integral, the addition of the meridian knot $\bigcirc_{\rm m}$ on the top of the knot K corresponds to including following term in the integrand of (150),

(151)
$$(-1) \times \operatorname{tr}_{R=2} \left(P \exp(\oint_{\bigcirc_{\mathrm{m}}} a) \right)$$

Note that the value of the above Wilson loop is always

$$-\left(e^{\frac{\pi i n}{k}} + e^{-\frac{\pi i n}{k}}\right) = -\left(q^{\frac{n}{2}} + q^{-\frac{n}{2}}\right)$$

for arbitrary gauge field a satisfying the boundary condition in (149). It explains the property of the knot invariant depicted in Figure 8.

Appendix C. Quantum dilogarithm function

The quantum dilogarithm function (Q.D.L) $\psi_{\hbar}(Z)$ is defined by [10]

(152)
$$\psi_{\hbar}(Z) := \begin{cases} \prod_{r=1}^{\infty} \frac{1-q^r e^{-Z}}{1-\tilde{q}^{-r+1}e^{-\tilde{Z}}} & \text{if } |q| < 1 , \\ \prod_{r=1}^{\infty} \frac{1-\tilde{q}^r e^{-\tilde{Z}}}{1-q^{-r+1}e^{-Z}} & \text{if } |q| > 1 , \end{cases}$$

with

(153)
$$q := e^{\hbar}, \qquad \tilde{q} := e^{-\frac{4\pi^2}{\hbar}}, \qquad \tilde{Z} := \frac{2\pi i Z}{\hbar}.$$

The function satisfies the following difference equations:

(154)
$$\psi_{\hbar}(Z+\hbar) = (1-e^{-Z})\psi_{\hbar}(Z), \quad \psi_{\hbar}(Z+2\pi i) = (1-e^{-\bar{Z}})\psi_{\hbar}(Z).$$

In the asymptotic limit $\hbar \to 0$,

(155)
$$\log \psi_{\hbar}(Z) \xrightarrow{\hbar \to 0} \sum_{n=0}^{\infty} \frac{B_n \hbar^{n-1}}{n!} \operatorname{Li}_{2-n}(e^{-Z}) .$$

Here B_n is the *n*-th Bernoulli number with $B_1 = 1/2$. The Q.D.L satisfies the following identity

(156)
$$\int \frac{dZ}{\sqrt{2\pi\hbar}} e^{-\frac{ZU}{\hbar}} \psi_{\hbar}(Z) = e^{i\delta} e^{\frac{U^2 - (2\pi i + \hbar)U}{2\hbar} + \frac{i\pi(b^2 + b^{-2})}{12}} \psi_{\hbar}(U)$$

with a constant overall phase factor $e^{i\delta}$.

References

- C. Beem, T. Dimofte, and S. Pasquetti. Holomorphic Blocks in Three Dimensions. JHEP, 12:177, 2014.
- F. Bonahon and H. Wong. Quantum traces for representations of surface groups in sl_2. arXiv preprint arXiv:1003.5250, 2010.
- [3] D. Bullock and W. L. Faro. The kauffman bracket skein module of a twist knot exterior. Algebraic & Geometric Topology, 5(1):107–118, 2005.
- [4] D. Bullock, C. Frohman, and J. Kania-Bartoszyńska. Understanding the kauffman bracket skein module. Journal of knot theory and its ramifications, 8(03):265–277, 1999.
- [5] T. Dimofte. Quantum Riemann Surfaces in Chern-Simons Theory. Adv. Theor. Math. Phys., 17(3):479–599, 2013.
- [6] T. Dimofte. Complex Chern–Simons Theory at Level k via the 3d–3d Correspondence. Commun. Math. Phys., 339(2):619–662, 2015.
- [7] T. Dimofte, D. Gaiotto, and S. Gukov. 3-Manifolds and 3d Indices. Adv. Theor. Math. Phys., 17(5):975–1076, 2013.
- [8] T. Dimofte, S. Gukov, J. Lenells, and D. Zagier. Exact results for perturbative chern-simons theory with complex gauge group. arXiv preprint arXiv:0903.2472, 2009.
- [9] T. D. Dimofte and S. Garoufalidis. The Quantum content of the gluing equations. Geom. Topol., 17:1253–1316, 2013.
- [10] L. D. Faddeev and R. M. Kashaev. Quantum Dilogarithm. Mod. Phys. Lett. A, 9:427–434, 1994.
- M. Gabella. Quantum Holonomies from Spectral Networks and Framed BPS States. Commun. Math. Phys., 351(2):563-598, 2017.
- [12] D. Gang, N. Kim, M. Romo, and M. Yamazaki. Aspects of Defects in 3d-3d Correspondence. JHEP, 10:062, 2016.
- [13] D. Gang and K. Yonekura. Symmetry enhancement and closing of knots in 3d/3d correspondence. JHEP, 07:145, 2018.
- [14] S. Garoufalidis. The 3d index of an ideal triangulation and angle structures. arXiv preprint arXiv:1208.1663, 2012.
- [15] S. Garoufalidis and T. T. Lê. The colored jones function is q-holonomic. Geometry & Topology, 9(3):1253–1293, 2005.
- [16] S. Gukov. Three-dimensional quantum gravity, chern-simons theory, and the a-polynomial. Communications in mathematical physics, 255(3):577–627, 2005.
- [17] K. Habiro. On the quantum sl2 invariants of knots and integral homology spheres. Geom. Topol. Monogr, 4:55–68, 2002.
- [18] K. Hikami. Generalized volume conjecture and the a-polynomials: the neumann-zagier potential function as a classical limit of the partition function. *Journal of Geometry and Physics*, 57(9):1895–1940, 2007.
- [19] R. M. Kashaev. Quantum dilogarithm as a 6j-symbol. Modern Physics Letters A, 9(40):3757– 3768, 1994.
- [20] R. M. Kashaev. A link invariant from quantum dilogarithm. Modern Physics Letters A, 10(19):1409–1418, 1995.
- [21] L. H. Kauffman. Statistical mechanics and the jones polynomial. Contemporary Mathematics, 78:263–297, 1988.
- [22] G. Masbaum. Skein-theoretical derivation of some formulas of habiro. Algebraic & Geometric Topology, 3(1):537–556, 2003.
- [23] H. Murakami. An introduction to the volume conjecture. Interactions between hyperbolic geometry, quantum topology and number theory, 541:1–40, 2011.
- [24] H. Murakami and J. Murakami. The colored jones polynomials and the simplicial volume of a knot. Acta Mathematica, 186(1):85–104, 2001.
- [25] W. D. Neumann and D. Zagier. Volumes of hyperbolic three-manifolds. *Topology*, 24(3):307– 332, 1985.
- [26] R. C. Penner. The decorated teichmüller space of punctured surfaces. Communications in Mathematical Physics, 113(2):299–339, 1987.
- [27] J. H. Przytycki. Skein modules of 3-manifolds. arXiv preprint math/0611797, 2006.
- [28] J. H. Przytycki and A. S. Sikora. On skein algebras and sl.2 (c)-character varieties. arXiv preprint q-alg/9705011, 1997.

- [29] W. P. Thurston. Minimal stretch maps between hyperbolic surfaces. arXiv preprint math/9801039, 1998.
- [30] W. P. Thurston and J. W. Milnor. The geometry and topology of three-manifolds, 1979.
- [31] S. Tillmann. Degenerations of ideal hyperbolic triangulations. Mathematische Zeitschrift, 272(3):793-823, 2012.
- [32] V. G. Turaev. The conway and kauffman modules of the solid torus. Zapiski Nauchnykh Seminarov POMI, 167:79–89, 1988.
- [33] E. Witten. Quantum field theory and the jones polynomial. Communications in Mathematical Physics, 121(3):351–399, 1989.
- [34] E. Witten. Analytic Continuation Of Chern-Simons Theory. AMS/IP Stud. Adv. Math., 50:347–446, 2011.

QUEEN MARY UNIVERSITY OF LONDON, MILE END ROAD, LONDON E1 4NS, UK

ELAITRA LTD Email address: agarwalprarit@gmail.com

DEPARTMENT OF PHYSICS AND ASTRONOMY & CENTER FOR THEORETICAL PHYSICS, SEOUL NATIONAL UNIVERSITY, 1 GWANAK-RO, SEOUL 08826, KOREA

ASIA PACIFIC CENTER FOR THEORETICAL PHYSICS (APCTP), POHANG 37673, KOREA *Email address*: arima275@snu.ac.kr

College of Liberal Studies, Seoul National University, Seoul 08826, Korea

DEPARTMENT OF PHYSICS AND ASTRONOMY & CENTER FOR THEORETICAL PHYSICS, SEOUL NATIONAL UNIVERSITY, 1 GWANAK-RO, SEOUL 08826, KOREA Email address: sangmin@snu.ac.kr

YAU MATHEMATICAL SCIENCES CENTER, TSINGHUA UNIVERSITY, BEIJING, 100084, CHINA *Email address*: mromoj@tsinghua.edu.cn