

ON CERTAIN THEOREMS IN OPERATIONAL CALCULUS.

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The object of this paper is twofold: firstly to establish certain theorems in Operational Calculus and secondly to obtain the Laplace transforms of several functions.

I.

1. Let us suppose [1]

$$\Phi(p) = p \int_0^{\infty} e^{-pt} f(t) dt \quad (1)$$

where p is a positive number (or a number whose real part is positive) and the integral on the right converges. We shall then say that $\Phi(p)$ is operationally related to $f(t)$ and symbolically

$$\Phi(p) \doteq f(t) \text{ or } f(t) \doteq \Phi(p). \quad (2)$$

Many interesting relations involving $\Phi(p)$ and $f(t)$ have been obtained. The following will be required in the sequel.

$$p \Phi(p) \doteq \frac{d}{dt} f(t), \text{ if } f(0) = 0 \quad (3)$$

$$p \frac{d}{dp} [\Phi(p)] \doteq -t \frac{d}{dt} f(t) \quad (4)$$

$$\frac{\Phi(p)}{p} \doteq \int_0^t f(t) dt \quad (5)$$

$$p \int_0^{\infty} \frac{\Phi(p)}{p} dp \doteq \frac{f(t)}{t} \quad (6)$$

$$p \frac{d}{dp} \left[\frac{\Phi(p)}{p} \right] \doteq -t f(t). \quad (7)$$

Also Goldstein [2] has proved that if

$$\Phi(p) = f(t), \quad \psi(p) = g(t),$$

then

$$\int_0^{\infty} \frac{\Phi(t) g(t) dt}{t} = \int_0^{\infty} \frac{\psi(t) f(t)}{t} dt, \quad (8)$$

provided the integrals converge.

It is known that if $h(t)$ is another function which satisfies (1), then

$$f(t) - h(t) = n(t),$$

where $n(t)$ is a null-function, i.e. a function such that

$$\int_0^t n(t) dt = 0, \text{ for every } t \geq 0.$$

If $f(t)$ is a continuous function which satisfies (1), then it is the only continuous function which satisfies (1). This theorem is due to Lerch [3].

2. Our object is to investigate that if either of the two functions $f(t)$ and $\Phi(t)$ has an assigned property, then will that property or an analogous property be true of the other function?

We know that

$$\frac{p}{(p^2 + b^2)^{n + \frac{1}{2}}} = \frac{\sqrt{\pi}}{2^n \Gamma(n + \frac{1}{2})} \left(\frac{t}{b}\right)^n J_n(bt). \quad (9)$$

Applying Goldstein's theorem, we get

$$b^2 \int_0^{\infty} \frac{f(t) dt}{(b^2 + t^2)^{n + \frac{1}{2}}} = \frac{\sqrt{\pi} b}{2^n \Gamma(n + \frac{1}{2})} \int_0^{\infty} \left(\frac{t}{b}\right)^{n-1} \Phi(t) J_n(bt) dt, \quad \Re(n) > -\frac{1}{2}. \quad (10)$$

Let us now put $b^2 = p$ and interpret. Assuming that $\frac{1}{p} = \kappa$, we get

$$\kappa^{n - \frac{1}{2}} \int_0^{\infty} e^{-t^2 \kappa} f(t) dt = \frac{\sqrt{\pi}}{2^n} \frac{1}{p^{\frac{1}{2}n - 1}} \int_0^{\infty} t^{n-1} \Phi(t) J_n(\sqrt{p}t) dt, \quad (11)$$

provided the integrals converge.

Again let us divide both sides of (10) by b and put $b = p$. On interpretation, we get

$$\int_0^\infty \left(\frac{x}{t}\right)^n f(t) J_n(xt) dt \doteq \int_0^\infty \left(\frac{t}{p}\right)^{n-1} \Phi(t) J_n(pt) dt, \quad R(n) > -\frac{1}{2}. \quad (12)$$

This can also be written in the form

$$x^{n-\frac{1}{2}} \int_0^\infty \sqrt{x} t^{n-\frac{1}{2}} f(t) J_n(xt) dt \doteq \frac{1}{p^{n-1}} \int_0^\infty t^{n-1} \Phi(t) J_n(pt) dt. \quad (13)$$

Suppose $t^{-n-\frac{1}{2}} f(t)$ is self-reciprocal in the Hankel transform of order n . Then

$$f(x)/x \doteq \int_0^\infty \left(\frac{t}{p}\right)^{n-1} \Phi(t) J_n(pt) dt. \quad (14)$$

But by (6),

$$p \int_p^\infty \frac{\Phi(p)}{p} dp \doteq \frac{f(x)}{x}.$$

Therefore

$$\int_0^\infty t^{n-1} \Phi(t) J_n(pt) dt = p^n \int_p^\infty \frac{\Phi(p)}{p} dp, \quad (15)$$

provided the integrals converge.

Dividing both sides by p^n and differentiating with respect to p (assuming that differentiation under the sign of integration is permissible and that $\Phi(t)/t$ is a continuous function of t in $(0, \infty)$), we get on writing $n-1$ for n ,

$$\int_0^\infty \sqrt{p} t^{n-\frac{3}{2}} \Phi(t) J_n(pt) dt = p^{n-\frac{3}{2}} \Phi(p), \quad (16)$$

showing that $t^{n-\frac{3}{2}} \Phi(t)$ is self-reciprocal in the Hankel transform of order n , when (16) converges.

Thus we have

Theorem I. If $t^{-n-\frac{1}{2}} f(t)$ is self-reciprocal in the Hankel transform of order n and $\Phi(t)/t$ is continuous in $(0, \infty)$ then $t^{n-\frac{3}{2}} \Phi(t)$ is self-reciprocal in the Hankel transform of order n .

We can also write (12) in the form

$$\int_0^{\infty} \left(\frac{\varkappa}{t}\right)^n f(t) J_n(\varkappa t) dt \doteq \frac{1}{p^{n-\frac{1}{2}}} \int_0^{\infty} \sqrt{pt} t^{n-\frac{3}{2}} \Phi(t) J_n(pt) dt. \quad (17)$$

Let $t^{n-\frac{3}{2}} \Phi(t)$ be self-reciprocal in the Hankel transform of order n . The (17) becomes

$$\int_0^{\infty} \left(\frac{\varkappa}{t}\right)^n f(t) J_n(\varkappa t) dt \doteq \frac{\Phi(p)}{p}. \quad (18)$$

But by (5),

$$\frac{\Phi(p)}{p} \doteq \int_0^{\varkappa} f(t) dt.$$

Hence by Lerch's theorem

$$\int_0^{\infty} \left(\frac{\varkappa}{t}\right)^n f(t) J_n(\varkappa t) dt = \int_0^{\varkappa} f(t) dt. \quad (19)$$

Differentiating both sides with respect to \varkappa (assuming that differentiation under the sign of integration is permissible and $f(t)$ is a continuous function of t), we get on writing $n+1$ for n

$$\int_0^{\infty} \sqrt{\varkappa t} t^{-n-\frac{1}{2}} f(t) J_n(\varkappa t) dt = \varkappa^{-n-\frac{1}{2}} f(\varkappa), \quad (20)$$

showing that $t^{-n-\frac{1}{2}} f(t)$ is self-reciprocal in the Hankel transform of order n . We thus have conversely,

Theorem II. If $t^{n-\frac{3}{2}} \Phi(t)$ is self-reciprocal in the Hankel transform of order n and $f(t)$ is continuous, then $t^{-n-\frac{1}{2}} f(t)$ is self-reciprocal in the Hankel transform of order n . •

In (12) let us put $n = \frac{1}{2}$. We obtain

$$\int_0^{\infty} \frac{f(t)}{t} \sin \varkappa t dt \doteq \int_0^{\infty} \frac{\Phi(t)}{t} \sin pt dt \quad (21)$$

By (4), we get

$$\varkappa \int_0^{\infty} f(t) \cos \varkappa t dt \doteq -p \int_0^{\infty} \Phi(t) \cos pt dt, \quad (22)$$

where we again assume that differentiation under the sign of integration is permissible.

If $\Phi(t)$ is self-reciprocal in the cosine transform, we obtain

$$\sqrt{\frac{2}{\pi}} \kappa \int_0^{\infty} f(t) \cos \kappa t dt \doteq -p \Phi(p). \tag{23}$$

But by (3),

$$p \Phi(p) \doteq f'(\kappa), \text{ if } f(0) = 0.$$

Hence

$$\sqrt{\frac{2}{\pi}} \kappa \int_0^{\infty} f(t) \cos \kappa t dt = -f'(\kappa).$$

Integrating the left hand side by parts, we have

$$\sqrt{\frac{2}{\pi}} \int_0^{\infty} f'(t) \sin \kappa t dt = f'(\kappa), \text{ when } f(\infty) = 0, \tag{24}$$

showing that $f'(t)$ is self-reciprocal in the sine transform. We therefore have

Theorem III. If $\Phi(t)$ is self-reciprocal in the cosine transform and $f(0) = f(\infty) = 0$, then $f'(\kappa)$ is self-reciprocal in the sine transform. Again integrating the left hand side of (22), we have

$$\int_0^{\infty} f'(t) \sin \kappa t dt \doteq p \int_0^{\infty} \Phi(t) \cos p t dt,$$

provided $f(\infty) = 0$.

If $f'(t)$ is self-reciprocal in the sine-transform, we get

$$f'(\kappa) \doteq \sqrt{\frac{2}{\pi}} p \int_0^{\infty} \Phi(t) \cos p t dt. \tag{25}$$

But when $f(0) = 0$, we have by (3), $f'(\kappa) \doteq p \Phi(p)$, so that

$$\sqrt{\frac{2}{\pi}} \int_0^{\infty} \Phi(t) \cos p t dt = \Phi(p), \tag{26}$$

showing that $\Phi(t)$ is self-reciprocal in the cosine transform. Hence the converse theorem follows, viz.,

Theorem IV. If $f(0) = f(\infty) = 0$ and $f'(x)$ is self-reciprocal in the sine transform, then $\Phi(t)$ is self-reciprocal in the cosine transform.

Again in (22) let $f(t)$ be self-reciprocal in the cosine transform. Then

$$xf(x) \doteq - \sqrt{\frac{2}{\pi}} p \int_0^{\infty} \Phi(t) \cos pt dt.$$

But by (7),

$$xf(x) \doteq -p \frac{d}{dp} \left[\frac{\Phi(p)}{p} \right],$$

so that

$$\sqrt{\frac{2}{\pi}} \int_0^{\infty} \Phi(t) \cos pt dt = \frac{d}{dp} \left[\frac{\Phi(p)}{p} \right]. \quad (27)$$

Integrating both sides with respect to p between the limits zero and p and changing the order of integration on the left (if that is permissible), we notice that if $\Phi(p)/p \rightarrow 0$ as $p \rightarrow 0$,

$$\sqrt{\frac{2}{\pi}} \int_0^{\infty} \frac{\Phi(t)}{t} \sin pt dt = \frac{\Phi(p)}{p}, \quad (28)$$

showing that $\Phi(t)/t$ is self-reciprocal in the sine transform. Hence we have

Theorem V. If $f(t)$ is self-reciprocal in the cosine transform and $\Phi(t)/t \rightarrow 0$ as $t \rightarrow 0$, then $\Phi(t)/t$ is self-reciprocal in the sine transform. Conversely, if $\Phi(t)/t$ is self-reciprocal in the sine transform, we have

$$\sqrt{\frac{2}{\pi}} \int_0^{\infty} \frac{\Phi(t)}{t} \sin pt dt = \frac{\Phi(p)}{p} \doteq \int_0^x f(t) dt, \text{ by (5).}$$

Hence by (4),

$$\sqrt{\frac{2}{\pi}} p \int_0^{\infty} \Phi(t) \cos pt dt \doteq -xf(x),$$

provided $f(t)$ is continuous and differentiation under the sign of integration is permissible.

But by (22),

$$x \int_0^{\infty} f(t) \cos xt dt \doteq -p \int_0^{\infty} \Phi(t) \cos pt dt.$$

Hence

$$\sqrt{\frac{2}{\pi}} \int_0^{\infty} f(t) \cos \kappa t dt = f(\kappa), \tag{29}$$

showing that $f(t)$ is self-reciprocal in the cosine transform. Thus we have

Theorem VI. If $\varphi(t)/t$ is self-reciprocal in the sine transform and $f(t)$ is continuous, then $f(t)$ is self-reciprocal in the cosine transform.

Theorem IV can also be extended to reciprocal functions.

$$\text{Let } \Phi(p) \doteq f(\kappa), \quad \psi(p) \doteq g(\kappa)$$

and

$$f(0) = g(0) = f(\infty) = g(\infty) = 0.$$

Then if $\Phi(p)$ is reciprocal to $\psi(p)$; $f'(\kappa)$ is reciprocal to $g'(\kappa)$ in the sine transform.

For, by (22)

$$\begin{aligned} \sqrt{\frac{2}{\pi}} \kappa \int_0^{\infty} f(t) \cos \kappa t dt &\doteq - \sqrt{\frac{2}{\pi}} p \int_0^{\infty} \Phi(t) \cos pt dt \\ &\doteq - p \psi(p) \\ &\doteq - g'(\kappa). \end{aligned}$$

Integrating the left hand side and applying Lerch's theorem, we obtain

$$\sqrt{\frac{2}{\pi}} \int_0^{\infty} f'(t) \sin \kappa t dt = g'(\kappa), \tag{30}$$

showing that $f'(\kappa)$ is reciprocal to $g'(\kappa)$ in the sine transform.

Conversely, let $f'(\kappa)$ be reciprocal to $g'(\kappa)$ in the sine transform, where $g(\kappa)$ is continuous in the arbitrary interval $(0, \kappa)$. Let $G(\kappa) = \int_0^{\infty} g(\kappa) d\kappa$, $\Phi(p) \doteq f(\kappa)$ and $\psi(p) \doteq G(\kappa)$. Then if $f(\infty) = 0$; $\Phi(p)$ is reciprocal to $\psi(p)$ in the cosine transform. We have

$$\sqrt{\frac{2}{\pi}} \int_0^{\infty} f'(t) \sin \kappa t dt = g(\kappa). \tag{31}$$

On integration, the left hand side becomes

$$- \sqrt{\frac{2}{\pi}} \kappa \int_0^{\infty} f(t) \cos \kappa t dt,$$

which, by (22) is equal (\doteq) to

$$\sqrt{\frac{2}{\pi}} p \int_0^{\infty} \Phi(t) \cos pt dt.$$

Therefore

$$\begin{aligned} \sqrt{\frac{2}{\pi}} p \int_0^{\infty} \Phi(t) \cos pt dt &\doteq g(\kappa) \\ &\doteq G'(\kappa) \\ &\doteq p \psi(p). \end{aligned}$$

Hence

$$\sqrt{\frac{2}{\pi}} \int_0^{\infty} \Phi(t) \cos pt dt = \psi(p),$$

showing that $\Phi(t)$ is reciprocal to $\psi(p)$ in the cosine transform.

3. A Functional Relation.

Let us now consider the relation (10). Putting $b^2 = p$ and interpreting, we obtain

$$\frac{2^n}{\sqrt{\pi}} \kappa^{n-\frac{1}{2}} \int_0^{\infty} e^{-t^2 \kappa} f(t) dt \doteq \frac{1}{p^{\frac{n-3}{2}}} \int_0^{\infty} t^{n-\frac{3}{2}} (V\sqrt{p}t)^{\frac{1}{2}} \Phi(t) J_n(V\sqrt{p}t) dt,$$

which is our relation (11).

Suppose $t^{n-\frac{3}{2}} \Phi(t)$ is self-reciprocal in the Hankel transform of order n . The right hand side is $\Phi(V\sqrt{p})$. But if $\Phi(p) \doteq f(t)$, then

$$\Phi(V\sqrt{p}) \doteq \frac{1}{\sqrt{\pi \kappa}} \int_0^{\infty} e^{-t^2/4\kappa} f(t) dt,$$

so that

$$2^n \kappa^n \int_0^{\infty} e^{-t^2 \kappa} f(t) dt = \int_0^{\infty} e^{-t^2/4\kappa} f(t) dt. \quad (32)$$

If we write

$$F(\kappa) = \int_0^{\infty} e^{-t^2 \kappa} f(t) dt,$$

the functional relation becomes

$$2^n \kappa^n F(\kappa) = F\left(\frac{1}{4\kappa}\right). \quad (33)$$

4. If $\Phi(p)$ is given by (1), then by Mellin's inversion formula [4],

$$f(t) = \frac{1}{2\pi i} \int_{c-i\infty}^{c+i\infty} e^{\lambda t} \frac{\Phi(\lambda)}{\lambda} d\lambda, \quad (c > 0) \quad (34)$$

The question naturally arises: if $f(t)$ and $\Phi(t)$ have these assigned properties, are there formulae for determining them otherwise if either of the two functions is known?

We know that

$$\frac{\kappa^n}{(t + \kappa)^{n+\frac{1}{2}}} \doteq 2^{n+\frac{1}{2}} \Gamma(n+1) V\bar{p} e^{\frac{1}{2}pt} D_{-2n-1}(V\sqrt{2pt}). \quad (35)$$

Applying Goldstein's theorem, we get after slight changes in the variables

$$\frac{1}{2^{n+\frac{1}{2}} \Gamma(n+1)} \int_0^\infty \frac{t^{n-1} \Phi(t) dt}{(t+p)^{n+\frac{1}{2}}} = \int_0^\infty t^{-\frac{1}{2}} e^{\frac{1}{2}pt} D_{-2n-1}(V\sqrt{2pt}) f(t) dt. \quad (36)$$

Writing t^2 for t and p^2 for p , the above relation becomes

$$\frac{1}{2^{n+\frac{1}{2}} \Gamma(n+1)} \int_0^\infty \frac{t^{2n-1} \Phi(t^2) dt}{(p^2+t^2)^{n+\frac{1}{2}}} = \int_0^\infty e^{\frac{1}{2}p^2t^2} D_{-2n-1}(V\sqrt{2pt}) f(t^2) dt. \quad (37)$$

Multiplying both sides by p and interpreting, we have on simplification,

$$\begin{aligned} \frac{\kappa^{n-\frac{1}{2}}}{\Gamma(2n+1)} \int_0^\infty V\sqrt{\kappa t} t^{n-\frac{3}{2}} \Phi(t^2) J_n(\kappa t) dt \\ \doteq V\sqrt{2p} \int_0^\infty e^{\frac{1}{2}p^2t^2} D_{-2n-1}(V\sqrt{2pt}) f(t^2) dt, \quad \Re(n) > -\frac{1}{2}. \end{aligned} \quad (38)$$

If $t^{n-\frac{3}{2}} \Phi(t^2)$ is self-reciprocal in the Hankel transform of order n , we get

$$\Phi(\kappa^2) \kappa^{2n-2} \doteq V\sqrt{2} \Gamma(2n+1) p \int_0^\infty e^{\frac{1}{2}p^2t^2} D_{-2n-1}(V\sqrt{2pt}) f(t^2) dt. \quad (39)$$

If $\Phi(t^2)/t$ is self-reciprocal in the sine transform,

$$\Phi(x^2)/x \doteq V\sqrt{2} p \int_0^\infty e^{\frac{1}{2} p^2 t^2} D_{-2}(V\sqrt{2} p t) f(t^2) dt. \quad (40)$$

Let us revert back to relation (10) once more. We can write it in the form

$$\frac{2^n}{V\pi} \Gamma(n + \frac{1}{2}) \int_0^\infty \frac{b^{n+\frac{1}{2}} f(t) dt}{(t^2 + b^2)^{n+\frac{1}{2}}} = \int_0^\infty V\sqrt{b} t^{n-\frac{3}{2}} \Phi(t) J_n(bt) dt. \quad (41)$$

If $t^{n-\frac{3}{2}} \Phi(t)$ is self-reciprocal in the Hankel Transform of order n , then

$$\Phi(b) = \frac{2^n \Gamma(n + \frac{1}{2})}{V\pi} b^2 \int_0^\infty \frac{f(t) dt}{(t^2 + b^2)^{n+\frac{1}{2}}}. \quad (42)$$

Conversely if $\Phi(b)$ is given by (42), then putting $b=p$ and interpreting, we get after a bit of reduction that $t^{-n+\frac{1}{2}} f(t)$ is self-reciprocal in the Hankel transform of order $n-1$, provided $f(t)$ is continuous and $n > 0$. If (42) holds and $t^{-n+\frac{1}{2}} f(t)$ is self-reciprocal in the Hankel transform of order $n-1$, then $\Phi(p) \doteq f(t)$. Again expressing the right hand side of (1) as a double integral and changing the order of integration (if that is permissible) we can prove that if $t^{-n+\frac{1}{2}} f(t)$ is self-reciprocal in the Hankel transform of order $n-1$, then $\Phi(b)$ is always given by (42).

We might also have derived similar relations by considering that [5]

$$f(t^2) \doteq \frac{p}{V\pi} \int_0^\infty e^{-p^2 x^2/4} \Phi\left(\frac{1}{x^2}\right) dx. \quad (43)$$

5. A double Integral theorem for $\Phi(t)$.

Let us consider the relation (12) again, Since by (7)

$$p \frac{d}{dp} \left[\frac{\Phi(p)}{p} \right] \doteq -x f(x)$$

we get on differentiating under the sign of integration (if that is permissible)

$$\int_0^\infty \frac{x^{n+1}}{t^n} f(t) J_n(xt) dt \doteq \int_0^\infty \frac{t^n}{p^{n-1}} \Phi(t) J_{n+1}(pt) dt, \quad \Re(n) > -\frac{1}{2}. \quad (44)$$

Also we know

$$\frac{\sqrt{\pi}}{2^{n+1} \Gamma(n + \frac{3}{2})} \frac{\kappa^{n+1} J_n(c\kappa)}{c^n} \doteq \frac{p^2}{(p^2 + c^2)^{n + \frac{3}{2}}}, \quad \Re(n) > -1. \quad (45)$$

Making use of Goldstein's Theorem, we obtain

$$\int_0^\infty \int_0^\infty \frac{\kappa^{n+2} f(t) J_n(\kappa t)}{t^n (\kappa^2 + c^2)^{n + \frac{3}{2}}} dt d\kappa = \frac{\sqrt{\pi}}{2^{n+1} \Gamma(n + \frac{3}{2}) c^n} \times \int_0^\infty \int_0^\infty \kappa t^n \Phi(t) J_n(c\kappa) J_{n+1}(\kappa t) dt d\kappa. \quad (46)$$

Let $c = \frac{1}{p}$ where we now assume that $\frac{1}{p} \doteq y$. Then on simplification, we have

$$y^{n+1} \int_0^\infty \int_0^\infty \frac{f(t)}{t^n} J_n(\kappa t) J_n\left(\frac{y}{\kappa}\right) \frac{dt d\kappa}{\kappa} \doteq \frac{1}{p^{n+1}} \int_0^\infty \int_0^\infty \kappa t^n \Phi(t) J_n\left(\frac{\kappa}{p}\right) J_{n+1}(\kappa t) dt d\kappa. \quad (47)$$

Writing $\frac{\kappa}{t}$ for κ , we get since κ and t are independent variables,

$$y^{n + \frac{1}{2}} \int_0^\infty \int_0^\infty \frac{f(t)}{t^{n + \frac{1}{2}}} \sqrt{yt} J_n(\kappa) J_n\left(\frac{yt}{\kappa}\right) dt \frac{d\kappa}{\kappa} \doteq \frac{1}{p^{n+1}} \int_0^\infty \int_0^\infty \kappa t^n \Phi(t) J_n\left(\frac{\kappa}{p}\right) J_{n+1}(\kappa t) dt d\kappa. \quad (48)$$

Professor Watson [6] has shown that

$$\tilde{\omega}_{\mu, \nu}(\kappa y) = \sqrt{\kappa y} \int_0^\infty J_\nu(t) J_\mu\left(\frac{\kappa y}{t}\right) \frac{dt}{t}, \quad (48')$$

can be taken as the kernel of a *new transform*. Let $f(\kappa)$ be an arbitrary function, and let $g(\kappa)$ be its transform with the Kernel $\tilde{\omega}_{\mu, \nu}(\kappa y)$, so that

$$g(\kappa) = \int_0^\infty \tilde{\omega}_{\mu, \nu}(\kappa y) f(y) dy.$$

Then assuming that the various changes in the order of integration are permissible, we have

$$\int_0^{\infty} \tilde{\omega}_{\mu, \nu}(\kappa y) g(y) dy = f(\kappa). \quad (49)$$

When $f(\kappa) = g(\kappa)$, we say that $f(\kappa)$ is self-reciprocal under this new transform. Hence in (48), if $t^{-n-\frac{1}{2}} f(t)$ is self-reciprocal under this transform¹, the left hand side is $f(y)$, so that

$$f(y) \doteq \frac{1}{p^{n+1}} \int_0^{\infty} \int_0^{\infty} \kappa t^n \Phi(t) J_n\left(\frac{\kappa}{p}\right) J_{n+1}(\kappa t) dt d\kappa \doteq \Phi(p).$$

Therefore

$$\Phi(p) = \frac{1}{p^{n+1}} \int_0^{\infty} \int_0^{\infty} \kappa t^n \Phi(t) J_n\left(\frac{\kappa}{p}\right) J_{n+1}(\kappa t) dt d\kappa. \quad (50)$$

This can be written in the more symmetrical form, after considerable simplification,

$$\Phi(p) = -p^2 \frac{d}{dp} \frac{1}{p^n} \left[\int_0^{\infty} \int_0^{\infty} t^{n-1} \Phi(t) J_n\left(\frac{\kappa}{p}\right) J_n(\kappa t) dt d\kappa \right], \quad \Re(n) > -\frac{1}{2}. \quad (51)$$

provided $\Phi(p)/p$ is continuous. Conversely if (51) holds, then $t^{-n-\frac{1}{2}} f(t)$ is self-reciprocal under this transform.

II.

6. Laplace transforms of certain functions.

Let us now consider the relation (11). We know that

$$J_\nu\left(\frac{a}{p}\right) \doteq J_\nu(\sqrt{2at}) I_\nu(\sqrt{2at}). \quad (52)$$

Let

$$f(t) = J_\nu(\sqrt{2at}) I_\nu(\sqrt{2at}) \text{ and } \Phi(t) = J_\nu\left(\frac{a}{t}\right).$$

We thus obtain

$$\frac{2^n \kappa^{n-\frac{1}{2}}}{\sqrt{\pi}} \int_0^{\infty} e^{-t^2 \kappa} J_\nu(\sqrt{2at}) I_\nu(\sqrt{2at}) dt \doteq \frac{1}{p^{2n-1}} \int_0^{\infty} J_\nu\left(\frac{a}{t}\right) J_n(\sqrt{p}t) t^{n-1} dt. \quad (53)$$

¹ The senior author has been able to construct certain examples giving functions which are self-reciprocal under this *new transform* and also the formal solutions of (49), when $f(\kappa) = g(\kappa)$.

But

$$J_\nu(a z) I_\nu(a z) = \sum_{m=0}^{\infty} \frac{(-1)^m (\frac{1}{2} a z)^{2\nu+4m}}{\Gamma(m+1) \Gamma(\nu+m+1) \Gamma(\nu+2m+1)}. \quad (54)$$

Integrating term by term and applying a result due to Hanumauta Rao [7], we obtain

$$\begin{aligned} & \frac{2^{n-2\nu-1} \kappa^{n-\frac{1}{2}\nu-1}}{\Gamma(\frac{1}{2}\nu+1) \Gamma(\nu+1)} {}_0F_2\left(\frac{1}{2}\nu+1, \nu+1; -\frac{a^2}{16\kappa}\right) \\ & \div \left\{ \frac{\Gamma(n-\frac{1}{2}\nu)}{2^{2\nu-n+1} \Gamma(\frac{1}{2}\nu+1) \Gamma(\nu+1)} \frac{1}{p^{n-\frac{1}{2}\nu-1}} {}_0F_3\left(\frac{1}{2}\nu+1, \nu+1, \frac{1}{2}\nu-n+1; \frac{a^2 p}{16}\right) \right. \\ & \left. + \frac{\Gamma(\frac{1}{2}\nu-n) a^{2n-\nu} p}{2^{3n+1} \Gamma(n+1) \Gamma(\frac{1}{2}\nu+n+1)} {}_0F_3\left(n+1, n-\frac{1}{2}\nu+1, n+\frac{1}{2}\nu+1; \frac{a^2 p}{16}\right) \right\}; \\ & \{ -\mathcal{R}(n+\frac{3}{2}) < \mathcal{R}(n) < \mathcal{R}(\nu+\frac{3}{2}) \text{ and } a > 0 \}. \end{aligned} \quad (54)$$

Again we know that

$$(2t)^{-1} J_0\left(\frac{y^2}{4t}\right) \div p J_0(y\sqrt{\frac{1}{2}p}) K_0(y\sqrt{\frac{1}{2}p}). \quad (55)$$

Let

$$f(t) = (2t)^{-1} J_0\left(\frac{y^2}{4t}\right) \text{ and } \Phi(t) = t J_0(y\sqrt{\frac{1}{2}t}) K_0(y\sqrt{\frac{1}{2}t}).$$

We get (when $n=0$)

$$\frac{1}{2V\pi\kappa} \int_0^\infty e^{-t^2\kappa} t^{-1} J_0\left(\frac{y^2}{4t}\right) dt \div p \int_0^\infty J_0(y\sqrt{\frac{1}{2}t}) K_0(y\sqrt{\frac{1}{2}t}) J_0(\sqrt{p}t) dt. \quad (56)$$

Putting $t=2z^2/y^2$, the right hand side becomes

$$4p y^{-2} \int_0^\infty z J_0(z) K_0(z) J_0\left(\sqrt{p} \frac{2z^2}{y^2}\right) dz.$$

By a result due to Mitra [8], the integral can be evaluated and we finally obtain

$$\frac{1}{V\pi\kappa} \int_0^\infty e^{-t^2\kappa} t^{-1} J_0\left(\frac{y^2}{4t}\right) dt \div V\bar{p} I_0\left(\frac{y^2}{8V\bar{p}}\right) K_0\left(\frac{y^2}{8V\bar{p}}\right). \quad (57)$$

The integral on the left can be evaluated by expressing it as a contour integral.

Again let

$$\Phi(t) = e^{-1/t} t^{-1/2}; \quad f(t) = (\pi)^{-\frac{1}{2}} \sin 2\sqrt{t}.$$

We get

$$\begin{aligned}
 (\pi)^{-\frac{1}{2}} \int_0^{\infty} e^{-t^2 x} \sin 2\sqrt{t} dt &\doteq p^{1/2} \int_0^{\infty} e^{-1/t} t^{-3/2} \sin \sqrt{pt} dt \\
 &\doteq (\pi)^2 e^{-\sqrt[1]{\frac{1}{2} p^4}} \sin \sqrt[1]{2 p^4}.
 \end{aligned} \tag{58}$$

The integral on the left is easily obtainable by direct term by term integration.

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