

**MEASURE VALUED SOLUTIONS TO THE SPATIALLY
HOMOGENEOUS BOLTZMANN EQUATION
WITHOUT ANGULAR CUTOFF**

YOSHINORI MORIMOTO, SHUAIKUN WANG, AND TONG YANG

ABSTRACT. A uniform approach is introduced to study the existence of measure valued solutions to the homogeneous Boltzmann equation for both hard potential with finite energy, and soft potential with finite or infinite energy, by using Toscani metric. Under the non-angular cutoff assumption on the cross-section, the solutions obtained are shown to be in the Schwartz space in the velocity variable as long as the initial data is not a single Dirac mass without any extra moment condition for hard potential, and with the boundedness on moments of any order for soft potential.

1. INTRODUCTION

The spatially homogeneous Boltzmann equation;

$$(1.1) \quad \partial_t f(t, v) = Q(f, f)(t, v),$$

describes the time evolution of the velocity distribution of a spatially homogeneous dilute gas of particles. The most interesting and important part of this equation is the collision operator given on the right hand side that captures the change rates of the density distribution through elastic binary collisions:

$$Q(g, f)(v) = \int_{\mathbb{R}^3} \int_{\mathbb{S}^2} B(v - v_*, \sigma) \{g(v'_*)f(v') - g(v_*)f(v)\} d\sigma dv_*,$$

where for $\sigma \in \mathbb{S}^2$

$$v' = \frac{v + v_*}{2} + \frac{|v - v_*|}{2}\sigma, \quad v'_* = \frac{v + v_*}{2} - \frac{|v - v_*|}{2}\sigma,$$

that follow from the conservation of momentum and energy,

$$v' + v'_* = v + v_*, \quad |v'|^2 + |v'_*|^2 = |v|^2 + |v_*|^2.$$

We consider the non-negative cross section B of the form

$$(1.2) \quad B(|v - v_*|, \cos \theta) = \Phi(|v - v_*|)b(\cos \theta), \quad \cos \theta = \frac{v - v_*}{|v - v_*|} \cdot \sigma, \quad 0 \leq \theta \leq \frac{\pi}{2},$$

where

$$(1.3) \quad \Phi(|z|) = \Phi_\gamma(|z|) = |z|^\gamma, \quad \text{for some } \gamma > -3,$$

$$(1.4) \quad b(\cos \theta)\theta^{2+2s} \rightarrow K \text{ when } \theta \rightarrow 0+, \text{ for } 0 < s < 1 \text{ and } K > 0.$$

Date: March 3, 2016.

2010 Mathematics Subject Classification. primary 35Q20, 76P05, secondary 35H20, 82B40, 82C40,

Key words and phrases. Boltzmann equation, homogenous, measure valued solutions, characteristic functions.

For the case $\gamma > 0$, we speak of hard potentials; For the case $-3 < \gamma < 0$, we speak of soft potentials, and for the case $\gamma = 0$, we speak of Maxwellian molecule.

An important example of this singular cross section is derived from the inverse power law potential $\rho^{-\gamma}$, ρ being the distance between two interacting particles, in which $s = \frac{1}{\gamma} \in (0, 1)$ and $\gamma = 1 - 4s \in (-3, 1)$. As usual, the range of θ can be restricted to $[0, \pi/2]$, by replacing $b(\cos \theta)$ by its ‘‘symmetrized’’ version

$$[b(\cos \theta) + b(\cos(\pi - \theta))]\mathbf{1}_{0 \leq \theta \leq \pi/2}.$$

The equation (1.1) is supplemented with an initial datum

$$(1.5) \quad f(0, v) = F_0,$$

which is generally assumed to be a density of a probability measure, after normalizing time variable t suitably. We consider the Cauchy problem (1.1)-(1.5) for an initial datum F_0 with a finite second moment, that is, $F_0 \in P_2(\mathbb{R}^3)$. Here, in general, we denote by $P_\alpha(\mathbb{R}^3)$, $\alpha \in [0, \infty)$, the set of probability measures F on \mathbb{R}^3 , such that

$$\int_{\mathbb{R}^3} |v|^\alpha dF(v) < \infty,$$

and moreover when $\alpha > 1$, it requires that

$$(1.6) \quad \int_{\mathbb{R}^3} v_j dF(v) = 0, \quad j = 1, 2, 3,$$

which is always satisfied by the translation in \mathbb{R}^3 .

The purpose of the present paper is to show that, for both hard ($\gamma > 0$) and moderately soft ($0 > \gamma \geq -2$) potentials, the Cauchy problem (1.1)-(1.5) have a measure valued solution in the following sense.

Definition 1.1 (Measure valued solution). *Let $B(\cdot)$ satisfy (1.2)-(1.4) and let $F_0 \in P_\alpha(\mathbb{R}^3)$ with $0 < \alpha \leq 2$. We say that $F_t \in L^\infty([0, \infty); P_\alpha(\mathbb{R}^3))$ is a measure valued solution to the Cauchy problem (1.1)-(1.5) if it satisfies for every $\psi(v) \in C_b^2(\mathbb{R}^3)$,*

$$(1.7) \quad \int_{\mathbb{R}^3} \psi(v) dF_t = \int_{\mathbb{R}^3} \psi(v) dF_0 + \frac{1}{2} \int_0^t \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} \int_{\mathbb{S}^2} b(\cdot) |v - v_*|^\gamma (\psi(v'_*) + \psi(v'_) - \psi(v_*) - \psi(v)) dF_\tau(v) dF_\tau(v_*) d\sigma d\tau.$$

In particular, when $\alpha = 2$, the energy conservation law holds

$$(1.8) \quad \forall t \geq 0, \quad \int |v|^2 dF_t = \int |v|^2 dF_0,$$

so that $F_t \in C([0, \infty); P_2(\mathbb{R}^3))$.

It should be noted that the moment conservation law holds because of (1.6), and the definition does not contain the finite entropy condition (cf., Definition 1 of [25]). The definition implicitly requires the finiteness of the second term of the right hand side of (1.7), similar as in Definition 1.1 of [13]. We are now ready to state our main results.

Theorem 1.2. *Let $B(\cdot)$ satisfy (1.2)-(1.4) and let $-2 \leq \gamma \leq 2$.*

(1) For any initial datum $F_0 \in P_2(\mathbb{R}^3)$, we have a measure valued solution $F_t \in C([0, \infty); P_2(\mathbb{R}^3))$ to the Cauchy problem (1.1)-(1.5). In the hard potential case

$\gamma > 0$, for any $\ell > 0$ and any $t_0 > 0$ we have a measure valued solution $F_t \in C([0, \infty); P_2(\mathbb{R}^3))$ such that

$$(1.9) \quad \int |v|^\ell dF_t < \infty \text{ for any } t \geq t_0.$$

In the (moderately) soft potential case $0 > \gamma (\geq -2)$, we have a measure valued solution $F_t \in C([0, \infty); P_2(\mathbb{R}^3))$ satisfying the property of the moment propagation:

$$(1.10) \quad \text{for any } \ell > 0, \int |v|^\ell dF_0 < \infty \text{ implies } \int |v|^\ell dF_t \in L_{loc}^\infty([0, \infty)).$$

(2) When $-2 \leq \gamma < 0$, let

$$c_{\gamma,s} = \begin{cases} \max\{\frac{\gamma}{2s} + 1, 0\} & \text{if } 0 < s < 1/2, \\ \max\{\gamma + 2s, 0\} & \text{if } \gamma + 2s < 1 \text{ and } 1/2 \leq s < 1, \\ \frac{\gamma}{2s-1} + 2 & \text{if } \gamma + 2s \geq 1 \text{ (, which implies } 1/2 < s). \end{cases}$$

Then, for any initial datum $F_0 \in P_\alpha(\mathbb{R}^3)$ with $\alpha \in (c_{\gamma,s}, 2)$, we have a measure valued solution $F_t \in C([0, \infty); P_\alpha(\mathbb{R}^3))$ to the Cauchy problem (1.1)-(1.5).

Remark 1.3. Different from the Maxwellian molecule case, $\gamma = 0$ (see [24], [16]), it is unknown that the uniqueness of the (measure valued) solution to the Cauchy problem (1.1)-(1.5) holds (see [14] about counter examples without the energy conservation law). Instead of (1.4), we may assume a slightly general condition

$$(1.11) \quad \exists \varepsilon > 0 \text{ such that } \int_0^{\pi/2} b(\cos \theta) \sin^{3-2\varepsilon} \theta d\theta < \infty,$$

which permits the angular cutoff case. In the angular cutoff case of hard potentials, our measure valued solution is the strong solution in the sense of

$$F_t \in C([0, \infty); P_2(\mathbb{R}^3)) \cap C^1([0, \infty); P_0(\mathbb{R}^3)),$$

because it follows from similar arguments as in [13], section 5.2, proof of part (a) that

$$t \mapsto Q^\pm(F_t, F_t) \in C([0, \infty); P_0(\mathbb{R}^3)),$$

where we write $Q = Q^+ - Q^-$.

Since the weak solution given in Theorem 1.2 is energy conservative, it enjoys the smoothing effect if the initial datum is not a single Dirac mass, by using the support expanding property of the energy conservative solution, studied by Fournier [9] in the non-cutoff case and Pulvirenti-Wennberg [22, 23] in the cutoff case.

Theorem 1.4. Let $B(\cdot)$ satisfy (1.2)-(1.4) and assume $2 \geq \gamma > \max\{-2s, -1\}$.

(1) **Case $\gamma > 0$:** If the initial datum $F_0 \in P_2(\mathbb{R}^3)$ is not a single Dirac mass, then for any $t_0 > 0$ there exists a solution $F_t \in C([0, \infty); P_2(\mathbb{R}^3))$ to the Cauchy problem (1.1)-(1.5) such that $F_t \in C^\infty([t_0, \infty); \mathcal{S}(\mathbb{R}^3))$.

(2) **Case $\gamma < 0$:** If $F_0 \in P_2(\mathbb{R}^3)$ is not a single Dirac mass and its support is compact (or $\int |v|^\ell dF_0 < \infty$ for any $\ell > 0$), then we have a solution $F_t \in C([0, \infty); P_2(\mathbb{R}^3)) \cap C^\infty((0, \infty); \mathcal{S}(\mathbb{R}^3))$.

As for the existence of weak solutions for the measure valued initial datum, the case $0 < \gamma \leq 2$ and $0 < s < 1$ was studied in Lu-Mouhot [13] by approximating the initial measure datum to the one with the density function which satisfies the finite entropy and the finite moments of any order, through the Mehler transform, see also Zhang-Zhang [28] for the case $0 < \gamma \leq 1$, by the probabilistic method. The

gain of the regularity up to the Besov space $B_{1,\infty}^s$ for a certain $s > 0$ was proved in Fournier [9], by the probabilistic method, assuming the existence of solutions with finite moments of any order in hard potential case and a certain order in soft case, if $-1 < \gamma < 1$, $0 < s < 1/2$, and this regularity gives the finiteness of the entropy for any positive time. For the other reference on the spatially homogeneous Boltzmann equation before 2002, one can refer the the classical review paper by Villani [26].

The plan of this paper is as follows: In Section 2, we construct the solution under the cutoff assumption on both kinetic and angular parts of the cross section, by using the Toscani metric ([8, 10, 24, 7]). The existence of the solution under the non cutoff assumption is proved in Section 3, by using the equi-continuity of the Fourier transform of approximation solutions obtained in Section 2, together with the gain of moments in the hard potential case proved in Section 4. Finally, we show the smoothness of measure valued solutions, applying methods developed in [4, 11, 17], and using the expanding of the support to the whole space proved in [9].

2. EXISTENCE UNDER CUT-OFF ASSUMPTION

Following Jacob [12] and Cannone-Karch [7], call the Fourier transform of a probability measure $F \in P_0(\mathbb{R}^3)$, that is,

$$\varphi(\xi) = \hat{f}(\xi) = \mathcal{F}(F)(\xi) = \int_{\mathbb{R}^3} e^{-iv \cdot \xi} dF(v),$$

a characteristic function. Put $\mathcal{K} = \mathcal{F}(P_0(\mathbb{R}^3))$. Inspired by a series of works by Toscani and his co-authors [8, 10, 24], Cannone-Karch defined a subspace \mathcal{K}^α for $\alpha \geq 0$ as follows:

$$(2.1) \quad \mathcal{K}^\alpha = \{\varphi \in \mathcal{K}; \|\varphi - 1\|_\alpha < \infty\},$$

where

$$(2.2) \quad \|\varphi - 1\|_\alpha = \sup_{\xi \in \mathbb{R}^3} \frac{|\varphi(\xi) - 1|}{|\xi|^\alpha}.$$

The space \mathcal{K}^α endowed with the distance

$$(2.3) \quad \|\varphi - \tilde{\varphi}\|_\alpha = \sup_{\xi \in \mathbb{R}^3} \frac{|\varphi(\xi) - \tilde{\varphi}(\xi)|}{|\xi|^\alpha}$$

is a complete metric space (see Proposition 3.10 of [7]). It follows that $\mathcal{K}^\alpha = \{1\}$ for all $\alpha > 2$ and the following embeddings (Lemma 3.12 of [7]) hold

$$\{1\} \subset \mathcal{K}^\alpha \subset \mathcal{K}^\beta \subset \mathcal{K}^0 = \mathcal{K} \quad \text{for all } 2 \geq \alpha \geq \beta \geq 0.$$

With this classification on characteristic functions, the global existence of measure valued solution for the Maxwellian molecules case in \mathcal{K}^α was studied in [7](see also [16]). It was proved in [20] that the smoothing effect of measure valued solutions always occurs except for a single Dirac mass initial datum. After that, the results about the existence and smoothing effect of solutions have been improved in [18, 19] by introducing more precise characteristic function spaces \mathcal{M}^α and $\tilde{\mathcal{M}}^\alpha$, which satisfy relations

$$\begin{aligned} \mathcal{K}^\alpha \supseteq \mathcal{F}(P_\alpha) = \mathcal{M}^\alpha = \tilde{\mathcal{M}}^\alpha \supseteq \mathcal{K}^{\alpha'} & \text{ if } 0 < \alpha < \alpha' \leq 2, \alpha \neq 1, \\ \mathcal{K}^1 \supseteq \mathcal{M}^1 \supseteq \mathcal{F}(P_1) = \tilde{\mathcal{M}}^1 \supseteq \mathcal{K}^{\alpha'} & \text{ if } 1 < \alpha' \leq 2. \end{aligned}$$

For the later use, we recall the space $\widetilde{\mathcal{M}}^\alpha$:

$$(2.4) \quad \widetilde{\mathcal{M}}^\alpha = \{\varphi \in \mathcal{K}; \|\operatorname{Re} \varphi - 1\|_{\mathcal{M}^\alpha} + \|\varphi - 1\|_\alpha < \infty\}, \quad \alpha \in (0, 2),$$

where $\operatorname{Re} \varphi$ stands for the real part of $\varphi(\xi)$ and

$$(2.5) \quad \|\varphi - 1\|_{\mathcal{M}^\alpha} = \int_{\mathbb{R}^3} \frac{|\varphi(\xi) - 1|}{|\xi|^{3+\alpha}} d\xi.$$

It should be noted (see [19], ex.) that $\mathcal{K}^2 = \mathcal{F}(P_2(\mathbb{R}^3))$, which is a key of our construction of the measure valued solution for the finite energy initial datum.

In this part, we consider the Cauchy problem for the cut-off equation, that is,

$$(2.6) \quad f_t = \int_{\mathbb{R}^3} \int_{\mathbb{S}^2} b(\cdot) \Phi_c(|v - v_*|) (f(v'_*) f(v') - f(v_*) f(v)) dv_* d\sigma,$$

with the initial data $F_0 \in P_\beta(\mathbb{R}^3)$, $0 < \beta \leq 2$, where we assume that b is integrable and without loss of generality, we can set

$$\int_{\mathbb{S}^2} b(\cdot) d\sigma = 1.$$

For the hard potential case, Φ_c takes the form $\Phi_c(|z|) = |z|^\gamma \phi_c(|z|)$, where $\phi_c \in C_0^\infty(\mathbb{R}^3)$ supporting on $\{z : |z| \leq 2\}$ satisfies $|\phi_c| \leq 1$ and $\phi_c = 1$ on $\{z : |z| \leq 1\}$. While for the soft potentials case, we introduce a different cut-off way of $|v - v_*|^\gamma$ ($-2 < \gamma < 0$): for any $r \geq 2$, let $\psi_r(|z|) \in \mathcal{D}(\mathbb{R})$ be a smooth function satisfying: (1) ψ_r support on $\{z : 1/2r \leq |z| \leq 2r\}$; (2) $0 \leq \psi_r \leq 1$; (3) $\psi_r = 1$ on $\{z : 1/r \leq |z| \leq r\}$. Then, let

$$\Phi_c(|z|) = |z|^\gamma \psi_r(|z|).$$

To investigate the equation (2.6), we start from the Bobylev formula (see [1] and [5, 6]) which reads as,

$$(2.7) \quad \varphi_t(\xi) = \int_{\mathbb{S}^2} b\left(\frac{\xi \cdot \sigma}{|\xi|}\right) \int_{\mathbb{R}^3} \hat{\Phi}_c(\zeta) (\varphi(\xi^- + \zeta) \varphi(\xi^+ - \zeta) - \varphi(\zeta) \varphi(\xi - \zeta)) d\zeta d\sigma.$$

where $\hat{\Phi}_c = \mathcal{F}(\Phi_c)$. As usual, we have the following estimates for $\hat{\Phi}_c$.

Lemma 2.1. *Let Φ_c be defined as above, then if $\Phi_c(|z|) = |z|^\gamma \phi_c(|z|)$ with $\gamma > -3$, we have*

$$\forall k \geq 0, k \in \mathbb{N}, \quad |\partial_\zeta^k \hat{\Phi}_c(\zeta)| \lesssim \frac{1}{\langle \zeta \rangle^{3+\gamma+k}};$$

if $\Phi_c(|z|) = |z|^\gamma \psi_r(|z|)$, for all $N \geq 1, k \geq 0, N, k \in \mathbb{N}$ and all $\zeta \in \mathbb{R}^3$, there exists $C_{r,N,k} > 0$ such that

$$|\partial_\zeta^k \hat{\Phi}_c(\zeta)| \lesssim \frac{C_{r,N,k}}{\langle \zeta \rangle^{2N}}.$$

Proof. The proof of the first part of the lemma can be found in [3]. When $k = 0$, the second estimate follows

$$\begin{aligned} \left| (1 + |\zeta|^2)^N \hat{\Phi}_c(\zeta) \right| &= \left| \int [(1 - \Delta_z)^N e^{-i\zeta \cdot z}] \Phi_c(|z|) dz \right| \\ &= \left| \int e^{-i\zeta \cdot z} (1 - \Delta_z)^N \Phi_c(|z|) dz \right| \\ &\leq \int_{\{1/2r \leq |z| \leq 1/r\} \cup \{r \leq |z| \leq 2r\}} |(1 - \Delta_z)^N \Phi_c(|z|)| dz + \|\Phi_c\|_{L^1(\mathbb{R}^3)} \\ &\leq C_{r,N,0}. \end{aligned}$$

The case $k \neq 0$ is similar because Φ_c supports on a compact subset of \mathbb{R}^3 which doesn't contain the origin. \square

We reformulate the problem as

$$(2.8) \quad \varphi_t + A\varphi = \mathcal{G}_1(\varphi) + \mathcal{G}_2(\varphi),$$

where $A = \sup_{z \in \mathbb{R}^3} |\Phi_c(|z|)|$ and

$$(2.9) \quad \mathcal{G}_1(\varphi)(\xi) = \int_{\mathbb{S}^2} b \left(\frac{\xi \cdot \sigma}{|\xi|} \right) \int_{\mathbb{R}^3} \hat{\Phi}_c(\zeta) \varphi(\xi^- + \zeta) \varphi(\xi^+ - \zeta) d\sigma d\zeta;$$

$$(2.10) \quad \mathcal{G}_2(\varphi)(\xi) = A\varphi(\xi) - \int_{\mathbb{S}^2} b \left(\frac{\xi \cdot \sigma}{|\xi|} \right) \int_{\mathbb{R}^3} \hat{\Phi}_c(\zeta) \varphi(\zeta) \varphi(\xi - \zeta) d\sigma d\zeta.$$

Then we can write the equation in the following integral form,

$$(2.11) \quad \varphi(t, \xi) = e^{-At} \varphi_0 + \int_0^t e^{-A(t-\tau)} [\mathcal{G}_1(\varphi(\tau, \xi)) + \mathcal{G}_2(\varphi(\tau, \xi))] d\tau.$$

Our first claim is

Lemma 2.2. *For every $\varphi \in \mathcal{K}$, $\mathcal{G}_1(\varphi)$ and $\mathcal{G}_2(\varphi)$ are continuous and positive definite. Moreover, if we assume*

$$(2.12) \quad \begin{cases} 0 < \alpha < \min\{\gamma, 1\} & \text{when } \gamma > 0, \\ 0 < \alpha \leq 2 & \text{when } \gamma < 0, \end{cases}$$

then for any $\varphi, \psi \in \mathcal{K}^\alpha$, there exists $C > 0$ such that

$$(2.13) \quad |\mathcal{G}_1(\varphi) + \mathcal{G}_2(\varphi) - \mathcal{G}_1(\psi) - \mathcal{G}_2(\psi)| \leq (A + C) \|\varphi - \psi\|_\alpha |\xi|^\alpha,$$

holds for all $\xi \in \mathbb{R}^3$.

Proof. To prove $\mathcal{G}_1(\varphi)$ is a characteristic function, let

$$\mathcal{G}_1^m(\xi) = \int_{\mathbb{S}^2} b \left(\frac{\xi \cdot \sigma}{|\xi|} \right) \int_{\mathbb{R}^3} \hat{\Phi}_c(\zeta) \varphi(\xi^- + \zeta) e^{-\frac{|\xi^- + \zeta|^2}{2m}} \varphi(\xi^+ - \zeta) e^{-\frac{|\xi^+ - \zeta|^2}{2m}} d\zeta d\sigma.$$

Since $\hat{\Phi}_c, b$ are integrable, we can apply the dominant convergence theorem to find that

$$\mathcal{G}_1^m(\xi) \rightarrow \mathcal{G}_1(\varphi)(\xi) \text{ pointwisely as } m \rightarrow \infty.$$

So we only need to prove for each m , \mathcal{G}_1^m is characteristic function. The following proof follows Lemma 2.1 of [21].

Denote $F = \mathcal{F}^{-1}(\varphi)$, then $\mathcal{G}_1^m(\xi)$ is the Fourier transform of the positive density

$$G_1^m(F)(v) = \int_{\mathbb{R}^3} \int_{\mathbb{S}^2} b \left(\frac{(v - v_*) \cdot \sigma}{|v - v_*|} \right) \Phi_c(|v - v_*|) f_m(v'_*) f_m(v') dv_* d\sigma$$

where

$$f_m(v) = \int_{\mathbb{R}^3} \omega_m(v - u) dF(u), \quad \omega_m(v) = \left(\frac{m}{2\pi} \right)^{3/2} e^{-\frac{m}{2}|v|^2}.$$

Similarly, to prove $\mathcal{G}_2(\varphi)$ is a characteristic function, let

$$\mathcal{G}_2^m(\xi) = A\varphi(\xi) e^{-\frac{|\xi|^2}{2m}} - \int_{\mathbb{S}^2} b \left(\frac{\xi \cdot \sigma}{|\xi|} \right) \int_{\mathbb{R}^3} \hat{\Phi}_c(\zeta) \varphi(\zeta) e^{-\frac{|\zeta|^2}{2m}} \varphi(\xi - \zeta) e^{-\frac{|\xi - \zeta|^2}{2m}} d\zeta d\sigma.$$

We can find that

$$\mathcal{G}_2^m(\xi) \rightarrow \mathcal{G}_2(\varphi)(\xi) \text{ pointwisely as } m \rightarrow \infty.$$

So we only need to prove for each m , \mathcal{G}_2^m is characteristic function. In fact, $\mathcal{G}_2^m(\xi)$ is the Fourier transform of the positive density

$$G_2^m(F)(v) = \left(A - \int_{\mathbb{R}^3} \int_{\mathbb{S}^2} b \left(\frac{(v - v_*) \cdot \sigma}{|v - v_*|} \right) \Phi_c(|v - v_*|) f_m(v_*) dv_* d\sigma \right) f_m(v)$$

where

$$f_m(v) = \int_{\mathbb{R}^3} \omega_m(v - u) dF(u), \quad \omega_m(v) = \left(\frac{m}{2\pi} \right)^{3/2} e^{-\frac{m}{2}|v|^2}.$$

Next, to show (2.13), we first consider the case $\Phi_c(|z|) = |z|^\gamma \phi_c(|z|)$, that is, hard potential case $\gamma > 0$. We only need to prove

$$(2.14) \quad \int_{\mathbb{S}^2} b \left(\frac{\xi \cdot \sigma}{|\xi|} \right) \int_{\mathbb{R}^3} |\hat{\Phi}_c(\zeta - \xi^-) - \hat{\Phi}_c(\zeta)| \frac{|\zeta|^\alpha + |\xi - \zeta|^\alpha}{|\xi|^\alpha} d\zeta d\sigma \leq C$$

for some constant $C > 0$.

For $|\xi| \leq 1$,

$$\begin{aligned} & \int_{\mathbb{R}^3} |\hat{\Phi}_c(\zeta - \xi^-) - \hat{\Phi}_c(\zeta)| \frac{|\zeta|^\alpha + |\xi - \zeta|^\alpha}{|\xi|^\alpha} d\zeta \\ & \leq \int_{\mathbb{R}^3} \int_0^1 |\partial_\zeta^1 \hat{\Phi}_c(\zeta - \tau \xi^-)| d\tau |\xi^-| \frac{|\zeta|^\alpha + |\xi - \zeta|^\alpha}{|\xi|^\alpha} d\zeta \\ & \lesssim \sin \frac{\theta}{2} \int_0^1 \int_{\mathbb{R}^3} \frac{|\zeta|^\alpha + |\xi - \zeta|^\alpha}{\langle \zeta - \tau \xi^- \rangle^{3+\gamma+1}} d\zeta d\tau \\ & \lesssim \sin \frac{\theta}{2} \int_0^1 \int_{\mathbb{R}^3} \frac{2|\zeta - \tau \xi^-|^\alpha + 1}{\langle \zeta - \tau \xi^- \rangle^{3+\gamma+1}} d\zeta d\tau \lesssim C \sin \frac{\theta}{2}. \end{aligned}$$

For $|\xi| > 1$,

$$\begin{aligned} & \int_{\mathbb{R}^3} |\hat{\Phi}_c(\zeta - \xi^-) - \hat{\Phi}_c(\zeta)| \frac{|\zeta|^\alpha + |\xi - \zeta|^\alpha}{|\xi|^\alpha} d\zeta \\ & \leq \int_{\mathbb{R}^3} |\hat{\Phi}_c(\zeta - \xi^-)| \frac{|\zeta|^\alpha + |\xi - \zeta|^\alpha}{|\xi|^\alpha} d\zeta + \int_{\mathbb{R}^3} |\hat{\Phi}_c(\zeta)| \frac{|\zeta|^\alpha + |\xi - \zeta|^\alpha}{|\xi|^\alpha} d\zeta \\ & \lesssim \int_{\mathbb{R}^3} \frac{2|\zeta - \xi^-|^\alpha + 2}{\langle \zeta - \xi^- \rangle^{3+\gamma}} d\zeta + \int_{\mathbb{R}^3} \frac{2|\zeta|^\alpha + 1}{\langle \zeta \rangle^{3+\gamma}} d\zeta \\ & \leq C. \end{aligned}$$

Therefore, (2.14) is obtained in the case $\gamma > 0$.

For the proof of (2.13) in the case $\Phi_c = |z|^\gamma \psi_r(|z|)$, that is, $\gamma < 0$, notice that $\hat{\Phi}_c(\zeta)$ is real-valued function and satisfies $\hat{\Phi}_c(\zeta) = \hat{\Phi}_c(-\zeta)$. Then we have

$$\begin{aligned} & \int_{\mathbb{S}^2} b \left(\frac{\xi \cdot \sigma}{|\xi|} \right) \int_{\mathbb{R}^3} [\hat{\Phi}_c(\zeta - \xi^-) - \hat{\Phi}_c(\zeta)] \varphi(\zeta) \varphi(\xi - \zeta) d\sigma d\zeta \\ & = \frac{1}{2} \int_{\mathbb{S}^2} b \left(\frac{\xi \cdot \sigma}{|\xi|} \right) \int_{\mathbb{R}^3} [\hat{\Phi}_c(-\zeta + \xi - \xi^-) + \hat{\Phi}_c(-\zeta + \xi^-) \\ & \quad - \hat{\Phi}_c(-\zeta + \xi) - \hat{\Phi}_c(-\zeta)] \varphi(\zeta) \varphi(\xi - \zeta) d\sigma d\zeta. \end{aligned}$$

Instead of (2.14), we show

$$(2.15) \quad \int_{\mathbb{R}^3} \left| \hat{\Phi}_c(-\zeta + \xi - \xi^-) + \hat{\Phi}_c(-\zeta + \xi^-) - \hat{\Phi}_c(-\zeta + \xi) - \hat{\Phi}_c(-\zeta) \right| \times \frac{|\zeta|^\alpha + |\xi - \zeta|^\alpha}{|\xi|^\alpha} d\zeta < C.$$

We only consider the case $|\xi| \leq 1$ (since another case $|\xi| > 1$ is easy). Using

$$\begin{aligned} & \left| \hat{\Phi}_c(-\zeta + \xi^-) - \hat{\Phi}_c(-\zeta) - \nabla \hat{\Phi}_c(-\zeta) \cdot \xi^- \right| \\ &= \left| \int_0^1 (1-\tau)(\xi^-)^T \nabla^2 \hat{\Phi}_c(-\zeta + \tau \xi^-)(\xi^-) d\tau \right| \lesssim |\xi|^2 \int_0^1 \frac{1}{\langle -\zeta + \tau \xi^- \rangle^{2N}} d\tau, \end{aligned}$$

and

$$\begin{aligned} & \left| \hat{\Phi}_c(-\zeta + \xi^+) - \hat{\Phi}_c(-\zeta + \xi) - \nabla \hat{\Phi}_c(-\zeta + \xi) \cdot (-\xi^-) \right| \\ &= \left| \int_0^1 (1-\tau)(-\xi^-)^T \nabla^2 \hat{\Phi}_c(-\zeta + \xi - \tau \xi^-)(-\xi^-) d\tau \right| \\ &\lesssim |\xi|^2 \int_0^1 \frac{1}{\langle -\zeta + \xi - \tau \xi^- \rangle^{2N}} d\tau, \end{aligned}$$

we have

$$\begin{aligned} & \left| \hat{\Phi}_c(-\zeta + \xi - \xi^-) + \hat{\Phi}_c(-\zeta + \xi^-) - \hat{\Phi}_c(-\zeta + \xi) - \hat{\Phi}_c(-\zeta) \right| \\ &\lesssim \left| \nabla \hat{\Phi}_c(-\zeta) \cdot \xi^- + \nabla \hat{\Phi}_c(-\zeta + \xi) \cdot (-\xi^-) \right| \\ &+ |\xi|^2 \int_0^1 \frac{1}{\langle -\zeta + \xi - \tau \xi^- \rangle^{2N}} + \frac{1}{\langle -\zeta + \tau \xi^- \rangle^{2N}} d\tau. \end{aligned}$$

Due to

$$\begin{aligned} \left| \nabla \hat{\Phi}_c(-\zeta) \cdot \xi^- - \nabla \hat{\Phi}_c(-\zeta + \xi) \cdot \xi^- \right| &= \left| \int_0^1 (-\xi)^T \nabla^2 \hat{\Phi}_c(-\zeta + \tau \xi)(\xi^-) d\tau \right| \\ &\lesssim |\xi|^2 \int_0^1 \frac{1}{\langle -\zeta + \tau \xi \rangle^{2N}} d\tau, \end{aligned}$$

we obtain

$$\begin{aligned} & \left| \hat{\Phi}_c(-\zeta + \xi - \xi^-) + \hat{\Phi}_c(-\zeta + \xi^-) - \hat{\Phi}_c(-\zeta + \xi) - \hat{\Phi}_c(-\zeta) \right| \\ &\lesssim |\xi|^2 \int_0^1 \frac{1}{\langle -\zeta + \tau \xi \rangle^{2N}} + \frac{1}{\langle -\zeta + \xi - \tau \xi^- \rangle^{2N}} + \frac{1}{\langle -\zeta + \tau \xi^- \rangle^{2N}} d\tau. \end{aligned}$$

Hence, (2.15) is clear for $|\xi| \leq 1$. \square

The existence and uniqueness of the solution to the cut-off equation can be obtained by applying the Banach contraction principle to the operator

$$\mathcal{F}(\varphi)(\xi, t) \equiv e^{-At} \varphi_0 + \int_0^t e^{-A(t-\tau)} [\mathcal{G}_1(\varphi(\tau, \xi)) + \mathcal{G}_2(\varphi(\tau, \xi))] d\tau.$$

For $T > 0$, which will be determined later, denote $\mathcal{X}_T^\alpha = C([0, T], \mathcal{K}^\alpha)$, then \mathcal{X}_T^α is a complete metric space with respect to the following distance:

$$\|\varphi(\xi, t) - \psi(\xi, t)\|_{\mathcal{X}_T^\alpha} = \sup_{t \in [0, T]} \|\varphi(\xi, t) - \psi(\xi, t)\|_\alpha, \quad \text{for all } \varphi, \psi \in \mathcal{X}_T^\alpha.$$

Then, we will prove

Lemma 2.3. *For all $\varphi \in \mathcal{K}$, $\mathcal{F}(\varphi)$ is continuous and positive definite, and satisfies $\mathcal{F}(\varphi)(t, 0) = 1$. Moreover, for any $\varphi, \tilde{\varphi} \in \mathcal{K}^\alpha$, there exists $T > 0$, which is independent of φ_0 , such that $\mathcal{F} : \mathcal{X}_T^\alpha \rightarrow \mathcal{X}_T^\alpha$ is a contraction mapping.*

Proof. The proof of $\mathcal{F}(\varphi)$ is continuous and positive definite follows Lemma 2.2. We can also see that $\mathcal{F}(\varphi)(t, 0) = 1$ by using the fact $\mathcal{G}_1(\varphi)(t, 0) + \mathcal{G}_2(\varphi)(t, 0) = A$. Since

$$\mathcal{F}(\varphi) - 1 = e^{-At}(\varphi_0 - 1) + \int_0^t e^{-A(t-\tau)}(\mathcal{G}_1(\varphi) + \mathcal{G}_2(\varphi) - \mathcal{G}_1(1) - \mathcal{G}_2(1))d\tau,$$

by Lemma 2.2, we have

$$\sup_{t \in [0, T]} \|\mathcal{F}(\varphi) - 1\|_\alpha \leq \|\varphi_0 - 1\|_\alpha + (A + C)T\|\varphi - \tilde{\varphi}\|_{\mathcal{X}_T^\alpha} < \infty.$$

Hence, we proved \mathcal{F} maps \mathcal{X}_T^α to itself.

Next, for any $\varphi, \tilde{\varphi} \in \mathcal{X}_T^\alpha$, using Lemma 2.13 again, we can obtain

$$\sup_{t \in [0, T]} \|\mathcal{F}(\varphi) - \mathcal{F}(\tilde{\varphi})\|_\alpha \leq (A + C)T\|\varphi - \tilde{\varphi}\|_{\mathcal{X}_T^\alpha}.$$

This shows \mathcal{F} is a contraction mapping if $T > 0$ is sufficiently small. \square

Lemma 2.3 shows that there exists a unique solution $\varphi(\xi, t) \in C([0, T], \mathcal{K}^\alpha)$ to the problem by the Banach fixed point theory. Since the choice of T doesn't depend on the initial data, we can apply the above argument to the initial data $\varphi(\xi, T)$ and we can obtain the solution in $C([T, 2T], \mathcal{K}^\alpha)$. Continuing to do this, we can obtain the following theorem.

Theorem 2.4. *Let b, Φ_c be with the same setting as the beginning of this section. Then for any $\varphi_0 \in \mathcal{K}^\alpha$ with α satisfying (2.12), there exists a unique classical solution $\varphi(\xi, t) \in C([0, \infty), \mathcal{K}^\alpha)$ to the Cauchy problem of equation (2.7).*

As a result, we can obtain

Theorem 2.5. *Let b, Φ_c be with the same setting as the beginning of this section. Then for any $F_0 \in P_\beta$, $0 < \beta \leq 2$, there exists a unique solution $F_t \in C([0, \infty), P_\alpha)$ with $\alpha \in (0, \beta)$ satisfying (2.12) to the Cauchy problem (2.6) in the distribution sense, that is, for any $\psi(v) \in C_b^2(\mathbb{R}^3)$, the following holds,*

$$(2.16) \quad \int \psi(v)dF_t = \int \psi(v)dF_0 + \frac{1}{2} \int_0^t \iint \int_{\mathbb{S}^2} b(\cdot)\Phi_c(\psi(v'_*) + \psi(v') - \psi(v_*) - \psi(v))dF_\tau(v_*)dF_\tau(v)d\sigma d\tau.$$

Moreover, if $F_0 \in P_2$, then the solution satisfies

$$(2.17) \quad \int_{\mathbb{R}^3} |v|^2 dF_t = \int_{\mathbb{R}^3} |v|^2 dF_0,$$

for all $t > 0$.

Proof. The existence and uniqueness of the solution follow Theorem 2.4 once notice the inclusion $\mathcal{F}^{-1}(\mathcal{K}^\beta) \subset P_\alpha$.

To show (2.17), we first prove $\int |v|^2 dF_t < \infty$. Let $W_\delta(v) = \langle v \rangle^2 \langle \delta v \rangle^{-2}$. Using $W_\delta(v') \leq W_\delta(v) + W_\delta(v_*)$, we can obtain

$$\begin{aligned} \frac{d}{dt} \int W_\delta(v) dF_t(v) &= \iint b(\cdot) \Phi_c W_\delta(v) (dF_t(v'_*) dF_t(v') - dF_t(v_*) dF_t(v)) \\ &\leq \iint b(\cdot) \Phi_c (W_\delta(v) + W_\delta(v')) dF_t(v_*) dF_t(v) \\ &\lesssim \iint b(\cdot) \Phi_c (W_\delta(v) + W_\delta(v_*)) dF_t(v_*) dF_t(v) \\ &\lesssim \int W_\delta(v) dF_t(v), \end{aligned}$$

Therefore,

$$\int W_\delta(v) dF_t(v) < C.$$

Letting $\delta \rightarrow 0$, we have $\int |v|^2 dF_t < \infty$. Then, (2.17) holds by using the fact that $|v|^2$ is collision invariant. \square

3. EXISTENCE UNDER NON-CUTOFF ASSUMPTION

In this section, we consider the Cauchy problem (1.1)-(1.5) under the assumptions $\gamma \geq -2$ and (1.11). For each $n \geq 2$, let $b_n = \min\{b, n\}$. For the hard potential case, let $\Phi_n(|z|) = \Phi(|z|) \phi_c(|z|/n)$; for the soft potential case, let $\Phi_n(|z|) = \Phi(|z|) \psi_n(|z|)$.

3.1. Finite energy initial datum. By Theorem 2.5, for any $F_0 \in P_2$, there exists a unique measure valued solution F_t^n to the equation

$$(3.1) \quad f_t = \int_{\mathbb{R}^3} \int_{\mathbb{S}^2} b_n(\cdot) \Phi_n(|v - v_*|) (f(v'_*) f(v') - f(v_*) f(v)) dv_* d\sigma,$$

and F_t^n satisfies (2.17). On the other hand, the corresponding solution $\varphi^n = \mathcal{F}(dF_t^n) \in C([0, \infty), \mathcal{K}^\alpha)$ solves

$$(3.2) \quad \frac{\partial \varphi}{\partial t} = \int_{\mathbb{R}^3} \int_{\mathbb{S}^2} b_n \left(\frac{\xi \cdot \sigma}{|\xi|} \right) \int_{\mathbb{R}^3} \hat{\Phi}_n(\zeta) (\varphi(\xi^- + \zeta) \varphi(\xi^+ - \zeta) - \varphi(\zeta) \varphi(\xi - \zeta)) d\zeta d\sigma,$$

with the initial data $\varphi_0 = \mathcal{F}(dF_0)$.

Lemma 3.1. *The sequence of solutions $\{\varphi^n(\xi, t)\}_{n=1}^\infty$ is bounded in $C(\mathbb{R}^3 \times [0, \infty))$ and equi-continuous on any compact subsets of $\mathbb{R}^3 \times [0, \infty)$.*

Proof. **Step 1:** Uniform bound. $|\varphi^n| \leq 1$.

Step 2: Equi-continuity with respect to ξ in a compact set in \mathbb{R}_ξ^3 . This follows from

$$(3.3) \quad |\varphi^n(\xi) - \varphi^n(\xi + \eta)| \leq \|1 - \varphi^n\|_\alpha \left(2|\xi|^{\alpha/2} |\eta|^{\alpha/2} + |\eta|^\alpha \right)$$

and

$$\begin{aligned}
\sup_t \sup_{\xi \neq 0} \frac{|\varphi^n - 1|}{|\xi|^\alpha} &= \sup_t \|\varphi^n(t) - 1\|_\alpha \\
&\leq C_\alpha \sup_t \int |v|^\alpha dF_t^n \\
&\leq C_\alpha \sup_t \int (1 + |v|^2) dF_t^n \\
&\leq C_\alpha \int (1 + |v|^2) dF_0.
\end{aligned}$$

For the proof of (3.3), we refer (19) of [16].

Step 3: Equi-continuity with respect to t . First we consider the soft potential case $0 > \gamma \geq -2$. Let $\psi \in C_b^2(\mathbb{R}^3)$ be a test function in the definition of the measure valued solution. Then for any $0 \leq s < t \leq T$, there exists a $C_\psi > 0$ independent of n such that

$$\begin{aligned}
(3.4) \quad &\left| \int \psi dF_t^n(v) - \int \psi dF_s^n(v) \right| \\
&\leq \frac{1}{2} \int_s^t \left| \iiint b_n(\psi'_* + \psi' - \psi_* - \psi) d\sigma \right. \\
&\quad \left. \times |v - v_*|^\gamma \psi_n(|v - v_*|) dF_\tau^n(v) dF_\tau^n(v_*) \right| d\tau \\
&\leq |t - s| C_\psi \max_{s \leq \tau \leq t} \int |v - v_*|^{\gamma+2} dF_\tau^n(v) dF_\tau^n(v_*) \\
&\leq |t - s| C_\psi \int \langle v \rangle^2 dF_0 \text{ since } \gamma \geq -2,
\end{aligned}$$

where we used the usual manners (cf, p.290-291 in [25]) as follows: If $\mathbf{k} = \frac{v - v_*}{|v - v_*|}$, then the Taylor expansion of order 2 for $\psi' - \psi$,

$$\begin{aligned}
\psi' - \psi &= \nabla \psi(v)(v' - v) + \int_0^1 (1 - \tau) \nabla^2 \psi(v + \tau(v' - v)) d\tau (v' - v)^2 \\
&= \frac{|v - v_*|}{2} \nabla \psi(v) \{(\sigma \cdot \mathbf{k}) \mathbf{k} - \sigma\} + \nabla \psi(v) \frac{v - v_*}{2} (1 - \mathbf{k} \cdot \sigma) + O(|v - v_*|^2 \theta^2)
\end{aligned}$$

and it for $\psi'_* - \psi_*$ imply

$$\begin{aligned}
(3.5) \quad &\psi'_* + \psi' - \psi_* - \psi - \frac{|v - v_*|}{2} \{ \nabla \psi(v) - \nabla \psi(v_*) \} \{ (\sigma \cdot \mathbf{k}) \mathbf{k} - \sigma \} \\
&= O(|v - v_*|^2 \theta^2)
\end{aligned}$$

because of $\nabla(\psi - \psi_*) = O(|v - v_*|)$, and we have

$$\left| \int_{\mathbb{S}^2} b_n(\psi'_* + \psi' - \psi_* - \psi) d\sigma \right| \leq C_\psi |v - v_*|^2,$$

since $\int_{\mathbb{S}^2} b_n(\sigma \cdot \mathbf{k}) \{(\sigma \cdot \mathbf{k}) \mathbf{k} - \sigma\} d\sigma = 0$ (where we used only $\int_0^{\pi/2} b(\cos \theta) \sin^3 \theta d\theta < \infty$ instead of (1.11)). Since we only consider the compact subset of $[0, \infty) \times \mathbb{R}_\xi^3$, for any fixed ξ which is bounded, we know $e^{-i\xi \cdot v} \in C_b^2(\mathbb{R}_v^3)$. This ends the proof once we take $\psi = e^{-i\xi \cdot v}$ in (3.4).

The hard potential case requires a subtle device, the Povzner inequality deeply studied in [15], which is briefly reviewed in the next section. Anyway, similar way to the soft potential case, we estimate

$$\begin{aligned} & \left| \int \psi dF_t^n(v) - \int \psi dF_s^n(v) \right| \\ & \leq \frac{1}{2} \int_s^t \left| \iiint b_n(\psi'_* + \psi' - \psi_* - \psi) d\sigma \right. \\ & \quad \left. \times |v - v_*|^\gamma \phi_c(|v - v_*|/n) dF_\tau^n(v) dF_\tau^n(v_*) \right| d\tau. \end{aligned}$$

Since it follows from the first order expansion that

$$\psi'_* + \psi' - \psi_* - \psi - \frac{|v - v_*|}{2} \{ \nabla \psi(v) - \nabla \psi(v_*) \} \{ (\sigma \cdot \mathbf{k}) \mathbf{k} - \sigma \} = O(|v - v_*| \theta)$$

$$\text{hence, } = O(|v - v_*| \theta)^{2-2\varepsilon} \quad \text{for any } 0 < \varepsilon < \frac{1}{2},$$

we have

$$\begin{aligned} & \left| \int \psi dF_t^n(v) - \int \psi dF_s^n(v) \right| \\ & \leq C_\psi \int_s^t \iiint \left(\int_0^{\pi/2} b(\cos \theta) \theta^{3-2\varepsilon} d\theta \right) \\ & \quad \times |v - v_*|^{\gamma+2-2\varepsilon} \phi_c(|v - v_*|/n) dF_\tau^n(v) dF_\tau^n(v_*) d\tau \\ & \leq C_\psi \left\{ |t - s| \max_{s \leq \tau \leq t} \int \langle v \rangle^2 dF_\tau^n(v) \int \langle v_* \rangle^2 dF_\tau^n(v_*) \right. \\ & \quad \left. + \int_s^t \left(\iint_{\mathbb{R}_v^3 \times \mathbb{R}_{v_*}^3} \mathbf{1}_{\{|v| \geq 2|v_*|\} \cup \{|v_*| \geq 2|v|\}} \right. \right. \\ & \quad \left. \left. \times |v - v_*|^{2-2\varepsilon} \min\{|v - v_*|^\gamma, n^\gamma\} dF_\tau^n(v) dF_\tau^n(v_*) \right) d\tau \right\} \\ & \leq C_\psi \left\{ |t - s| \left(\int \langle v \rangle^2 dF_0(v) \right)^2 + \int_s^t \left(\int \langle v \rangle^{2-2\varepsilon} \min\{\langle v \rangle^\gamma, n^\gamma\} dF_\tau^n(v) \right) d\tau \right\}. \end{aligned}$$

Now we use (4.16) in the next section, that is, for any $T > 0$ there exists a $C_T > 0$ independent of n such that

$$\int_0^T \left(\int \langle v \rangle^{2-\varepsilon} \min\{\langle v \rangle^\gamma, n^\gamma\} dF_\tau^n(v) \right) d\tau < C_T.$$

If $0 \leq s < t \leq T$ and if $R > 1$, then we have

$$\int_s^t \left(\int_{\langle v \rangle \geq R} \langle v \rangle^{2-2\varepsilon} \min\{\langle v \rangle^\gamma, n^\gamma\} dF_\tau^n(v) \right) d\tau \leq \frac{C_T}{R^\varepsilon},$$

which leads us to the equi-continuity in t because

$$\int_s^t \left(\int_{\langle v \rangle < R} \langle v \rangle^{2-2\varepsilon} \min\{\langle v \rangle^\gamma, n^\gamma\} dF_\tau^n(v) \right) d\tau \lesssim R^4 |t - s|.$$

□

By Lemma 3.1, the Ascoli-Arzelà theorem, and the Cantor diagonal argument, we can conclude that there exists a subsequence of solutions $\{\varphi^{n_k}\}_{n_k} \subset \{\varphi^n\}_n$,

denoted still by $\{\varphi^n\}_{n=1}^\infty$, which converges uniformly in every compact subset of $\mathbb{R}^3 \times [0, \infty)$. Let

$$\varphi(\xi, t) = \lim_{n \rightarrow \infty} \varphi^n(\xi, t),$$

then $\varphi(\cdot, t)$ is a characteristic function for every $t \geq 0$ as the point-wise limit of characteristic functions. Denote $F_t = \mathcal{F}^{-1}(\varphi(\xi, t))$, then $F_t^n \rightarrow F_t$ in $\mathcal{S}'(\mathbb{R}^3)$. We will prove F_t is a measure valued solution to the equation (1.1). Before this, we first prove the following lemma:

Lemma 3.2. *Let F_t^n, F_t be defined as above.*

(i) *For any $t > 0$, we have*

$$(3.6) \quad \int |v|^2 dF_t \leq \int |v|^2 dF_0.$$

(ii) *For any $T > 0$ and $\psi \in C(\mathbb{R}^3)$, satisfying the growth condition $|\psi(v)| \lesssim \langle v \rangle^l$, for some $0 < l < 2$, we have*

$$(3.7) \quad \lim_{n \rightarrow \infty} \int \psi(v) dF_t^n = \int \psi(v) dF_t \quad \text{uniformly to } t \in [0, T].$$

(iii) *For any $T > 0$ and $\Psi(v, v_*) \in C(\mathbb{R}^3 \times \mathbb{R}^3)$, satisfying the growth condition $|\Psi(v, v_*)| \lesssim \langle v \rangle^l \langle v_* \rangle^l$, for some $0 < l < 2$, we have*

$$(3.8) \quad \lim_{n \rightarrow \infty} \iint \Psi(v, v_*) dF_t^n(v) dF_t^n(v_*) = \iint \Psi(v, v_*) dF_t(v) dF_t(v_*),$$

uniformly to $t \in [0, T]$.

Proof. (i) We first claim that $\int |v|^2 dF_t < \infty$ for any $t > 0$. For any $\delta \in (0, 1)$,

$$\sup_{\delta < |\xi| < \delta^{-1}} \frac{|\varphi^n - 1|}{|\xi|^2} \leq \|\varphi^n - 1\|_2 \leq C \int |v|^2 dF_t^n \leq C \int |v|^2 dF_0.$$

Letting $\delta \rightarrow 0$, we know there exists $C > 0$, such that

$$\|\varphi(t) - 1\|_2 = \lim_{\delta \rightarrow 0} \sup_{\delta < |\xi| < \delta^{-1}} \frac{|\varphi - 1|}{|\xi|^2} = \lim_{\delta \rightarrow 0} \lim_{n \rightarrow \infty} \sup_{\delta < |\xi| < \delta^{-1}} \frac{|\varphi^n - 1|}{|\xi|^2} < C \int |v|^2 dF_0.$$

Then, since it follows from $\mathcal{K}^2 = \mathcal{F}(P_2(\mathbb{R}^3))$ (see [19], ex.) that $\int |v|^2 dF_t$ is bounded, for any $\epsilon > 0$, we can select $R_{\epsilon, t} > 0$ large enough such that

$$\int_{|v| \geq R_{\epsilon, t}} |v|^2 dF_t < \epsilon.$$

Notice that $F_t^n \rightarrow F_t$ in $\mathcal{S}'(\mathbb{R}^3)$, we have

$$\int |v|^2 dF_t \leq \lim_{n \rightarrow \infty} \int |v|^2 \chi_{\{|v| \leq R_{\epsilon, t}\}} dF_t^n + \epsilon \leq \int |v|^2 dF_0 + \epsilon.$$

Letting $\epsilon \rightarrow 0$, we obtain (3.6).

(ii) Since for any $R > 0$,

$$R^{2-l} \int_{|v| \geq R} |\psi(v)| dF_t^n \lesssim \int_{|v| \geq R} \langle v \rangle^2 dF_t^n \leq \int \langle v \rangle^2 dF_t^n = \int \langle v \rangle^2 dF_0,$$

for any $\epsilon > 0$, we can find $R_\epsilon > 0$ large enough such that

$$\int_{|v| \geq R_\epsilon} |\psi(v)| dF_t^n \leq \epsilon.$$

The same argument can be applied to F_t . As a result, we can select $R'_\epsilon > 0$ large enough, such that

$$(3.9) \quad \int_{|v| \geq R'_\epsilon} |\psi(v)| dF_t^n + \int_{|v| \geq R'_\epsilon} |\psi(v)| dF_t \leq 2\epsilon.$$

Let $\chi_{R'_\epsilon}(v)$ be a $C_0^\infty(\mathbb{R}^3)$ function such that $\chi_{R'_\epsilon} = 1$ on $\{|v| \leq R'_\epsilon\}$. Then, by setting $\hat{\psi} = \mathcal{F}(\chi_{R'_\epsilon} \psi) \in \mathcal{S}(\mathbb{R}^3)$ we have

$$(3.10) \quad \left| \int \chi_{R'_\epsilon} \psi dF_t^n - \int \chi_{R'_\epsilon} \psi dF_t \right| = (2\pi)^3 \left| \int \hat{\psi}(\xi) (\varphi^n(\xi, t) - \varphi(\xi, t)) d\xi \right| \\ \lesssim M^{-1} \int_{|\xi| \geq M} \langle \xi \rangle \hat{\psi}(\xi) d\xi + \sup_{|\xi| \leq M} |\varphi^n(\xi, t) - \varphi(\xi, t)|$$

for any arbitrary $M > 0$. Finally, (3.7) follows (3.9) and (3.10).

(iii) For any $R > 0$, denote

$$K_R = \{(v, v_*) : |v| \leq R, |v_*| \leq R\}$$

which is compact in $\mathbb{R}^3 \times \mathbb{R}^3$. For any $\epsilon > 0$, since

$$\iint_{K_R^c} |\Psi(v, v_*)| dF_t^n(v) dF_t^n(v_*) \lesssim \frac{1}{\langle R \rangle^{2-l}} \iint_{K_R^c} \langle v \rangle^2 \langle v_* \rangle^2 dF_t^n(v) dF_t^n(v_*) \\ \leq \frac{1}{\langle R \rangle^{2-l}} \left(\int \langle v \rangle^2 dF_t^n(v) \right)^2 \\ \leq \frac{1}{\langle R \rangle^{2-l}} \left(\int \langle v \rangle^2 dF_0(v) \right)^2,$$

we can find $R_\epsilon > 0$ large enough, such that

$$\iint_{K_{R_\epsilon}^c} |\Psi(v, v_*)| dF_t^n(v) dF_t^n(v_*) + \iint_{K_{R_\epsilon}^c} |\Psi(v, v_*)| dF_t(v) dF_t(v_*) \leq \epsilon.$$

To show (3.8), it remains to prove

$$(3.11) \quad \lim_{n \rightarrow \infty} \iint_{K_{R_\epsilon}} \Psi(v, v_*) dF_t^n(v) dF_t^n(v_*) = \iint_{K_{R_\epsilon}} \Psi(v, v_*) dF_t(v) dF_t(v_*).$$

By the Weierstrass approximation theorem, for any $\epsilon > 0$, there exists a polynomial function $\sum_j p_j(v) \tilde{p}_j(v_*)$ such that

$$\sup_{(v, v_*) \in K_{R_\epsilon}} \left| \Psi(v, v_*) - \sum_{j=1}^{N(\epsilon)} p_j(v) \tilde{p}_j(v_*) \right| < \epsilon.$$

From (ii), using the standard smoothing argument, we can obtain that for each p_j (resp. \tilde{p}_j), we have

$$\lim_{n \rightarrow \infty} \int p_j(v) \chi_{\{|v| \leq R_\epsilon\}} dF_t^n(v) = \int p_j(v) \chi_{\{|v| \leq R_\epsilon\}} dF_t(v), \quad \text{uniformly to } t \in [0, T],$$

that is, for $\epsilon > 0$, we can find $N_{j,\epsilon}$ (resp. $\tilde{N}_{j,\epsilon}$) $\in \mathbb{N}$, such that for all $n > N_{j,\epsilon}$, we have

$$\left| \int_{|v| \leq R_\epsilon} p_j(v) dF_t^n(v) - \int_{\{|v| \leq R_\epsilon\}} p_j(v) dF_t(v) \right| \leq \frac{\epsilon}{\alpha_j N(\epsilon)},$$

where

$$\alpha_j = \sup_{|v| \leq R} |p_j(v)| + \sup_{|v_*| \leq R} |\tilde{p}_j(v_*)|.$$

If

$$N = \max_{1 \leq j \leq N(\epsilon)} \{N_{j,\epsilon}, \tilde{N}_{j,\epsilon}\},$$

then, for all $n > N$, we have

$$\begin{aligned} & \left| \iint_{K_{R_\epsilon}} \Psi(v, v_*) dF_t^n(v) dF_t^n(v_*) - \iint_{K_{R_\epsilon}} \Psi(v, v_*) dF_t(v) dF_t(v_*) \right| \\ & \leq \sum_{j=1}^{N(\epsilon)} \left| \iint_{K_{R_\epsilon}} p_j(v) \tilde{p}_j(v_*) dF_t^n(v) dF_t^n(v_*) - \iint_{K_{R_\epsilon}} p_j(v) \tilde{p}_j(v_*) dF_t(v) dF_t(v_*) \right| + 2\epsilon \\ & \leq \sum_{j=1}^{N(\epsilon)} \left| \int_{|v_*| \leq R_\epsilon} \tilde{p}_j(v_*) dF_t^n(v_*) \left(\int_{|v| \leq R_\epsilon} p_j(v) dF_t^n(v) - \int_{|v| \leq R_\epsilon} p_j(v) dF_t(v) \right) \right| \\ & \quad + \sum_{j=1}^{N(\epsilon)} \left| \int_{|v| \leq R_\epsilon} p_j(v) dF_t(v) \left(\int_{|v_*| \leq R_\epsilon} \tilde{p}_j(v_*) dF_t^n(v_*) - \int_{|v_*| \leq R_\epsilon} \tilde{p}_j(v_*) dF_t(v_*) \right) \right| \\ & \quad + 2\epsilon \\ & \leq 3\epsilon. \end{aligned}$$

Letting $\epsilon \rightarrow 0$ leads to (3.8). \square

We are ready to prove the existence part of Theorem 1.2 in the soft potential case.

Proof of (1) of Theorem 1.2 in the case $-2 \leq \gamma < 0$. As in the proof of Lemma 3.1, note that

$$\int_{\mathbb{S}^2} b(\sigma \cdot \mathbf{k}) \{(\sigma \cdot \mathbf{k}) \mathbf{k} - \sigma\} d\sigma = \lim_{n \rightarrow \infty} \int_{\mathbb{S}^2} b_n(\sigma \cdot \mathbf{k}) \{(\sigma \cdot \mathbf{k}) \mathbf{k} - \sigma\} d\sigma = 0$$

and $\nabla(\psi - \psi_*) = O(|v - v_*|)$. Therefore, if we put

$$(3.12) \quad \begin{aligned} \Delta\psi & := \psi'_* + \psi' - \psi_* - \psi - \frac{|v - v_*|}{2} \{\nabla\psi(v) - \nabla\psi(v_*)\} \{(\sigma \cdot \mathbf{k}) \mathbf{k} - \sigma\} \\ & (= O(|v - v_*|^2 \theta^2)), \end{aligned}$$

then we have

$$\begin{aligned}
& \iint \int_{\mathbb{S}^2} b_n |v - v_*|^\gamma \psi_n(|v - v_*|) (\psi'_* + \psi' - \psi_* - \psi) dF_t^n(v) dF_t^n(v_*) d\sigma \\
& \quad - \iint \int_{\mathbb{S}^2} b |v - v_*|^\gamma (\psi'_* + \psi' - \psi_* - \psi) dF_t(v) dF_t(v_*) d\sigma \\
& = \iint \int_{\mathbb{S}^2} b_n |v - v_*|^\gamma \psi_n(|v - v_*|) \Delta \psi dF_t^n(v) dF_t^n(v_*) d\sigma \\
& \quad - \iint \int_{\mathbb{S}^2} b |v - v_*|^\gamma \Delta \psi dF_t(v) dF_t(v_*) d\sigma \\
& = \iint \left(\Psi_n(v, v_*) - \Psi(v, v_*) \right) dF_t^n(v) dF_t^n(v_*) \\
& \quad + \left(\iint \Psi(v, v_*) dF_t^n(v) dF_t^n(v_*) - \iint \Psi(v, v_*) dF_t(v) dF_t(v_*) \right)
\end{aligned}$$

where $\Psi_n(v, v_*) = |v - v_*|^\gamma \psi_n(|v - v_*|) \int_{\mathbb{S}^2} b_n \Delta \psi d\sigma$ and $\Psi(v, v_*) = |v - v_*|^\gamma \int_{\mathbb{S}^2} b \Delta \psi d\sigma$. Note that Ψ_n, Ψ satisfy the growth condition of Lemma 3.2, (iii) with $\ell = 2 + \gamma < 2$. Since, for any $R > 0$, we have

$$\begin{aligned}
& \left| \iint_{\{|v|>R\} \cup \{|v_*|>R\}} \left(\Psi_n(v, v_*) - \Psi(v, v_*) \right) dF_t^n(v) dF_t^n(v_*) \right| \\
& \lesssim R^{\ell-2} \left(\int \langle v \rangle^2 dF_0(v) \right)^2,
\end{aligned}$$

the first term converges to 0, as $n \rightarrow \infty$, because $\Psi_n(v, v_*)$ converges to $\Psi(v, v_*)$ uniformly on a compact set of $\mathbb{R}_v^3 \times \mathbb{R}_{v_*}^3$. By Lemma 3.2, (iii), we also see the second term tends to 0, as $n \rightarrow \infty$. Thus F_t is a measure valued solution. Following the same method as in p.292-293 of [25], we can show (1.8). In fact, by considering $\psi = \psi_m = |v|^2 \phi_c(|v|/m) \in C_b^2(\mathbb{R}^3)$, we have

$$\begin{aligned}
& \int \psi_m dF_t(v) - \int \psi_m dF_0(v) \\
& = \int_0^t \left(\iint \int_{\mathbb{S}^2} b |v - v_*|^\gamma (\psi'_{m,*} + \psi'_m - \psi_{m,*} - \psi_m) dF_\tau(v) dF_\tau(v_*) d\sigma \right) d\tau.
\end{aligned}$$

Since it follows from (3.5) that the right hand side is uniformly bounded with respect to m , for any $\eta > 0$ there exists a $\delta > 0$ independent of m such that

$$\left| \iiint_{\{|v-v_*|<\delta\} \cup \{|\theta|<\delta\}} b |v - v_*|^\gamma (\psi'_{m,*} + \psi'_m - \psi_{m,*} - \psi_m) dF_t(v) dF_t(v_*) d\sigma \right| < \eta,$$

and hence, the Lebesgue convergence theorem shows

$$\lim_{m \rightarrow \infty} \left| \int \psi_m dF_t(v) - \int \psi_m dF_0(v) \right| \leq \eta,$$

which yields (1.8). It remains to show $F_t \in C([0, \infty); P_2(\mathbb{R}^3))$. Taking the limit of (3.4), we obtain $F_t \in C([0, \infty); P_0(\mathbb{R}^3))$. This, together with (1.8), shows $F_t \in C([0, \infty); P_2(\mathbb{R}^3))$, since it follows from [24, Lemma 1] (see also [27, Theorem 7.12]) that for any $t_0 \in [0, \infty)$

$$(3.13) \quad \lim_{R \rightarrow \infty} \left(\limsup_{t \rightarrow t_0} \int_{|v| \geq R} |v|^2 dF_t(v) \right) = 0.$$

The property of the moment propagation (1.10) is a direct consequence of the Povzner inequality mentioned in the next section. \square

We now consider the hard potential case, admitting some results in the next section.

Proof of (1) of Theorem 1.2 when $\gamma > 0$. We consider the convergence of the difference

$$\begin{aligned} & \int_0^t \iint \int_{\mathbb{S}^2} b_n |v - v_*|^\gamma \psi_n(|v - v_*|) (\psi'_* + \psi' - \psi_* - \psi) dF_\tau^n(v) dF_\tau^n(v_*) d\sigma \\ & \quad - \int_0^t \iint \int_{\mathbb{S}^2} b |v - v_*|^\gamma (\psi'_* + \psi' - \psi_* - \psi) dF_\tau(v) dF_\tau(v_*) d\sigma d\tau \\ & = \int_0^t \iint \Psi_n(v, v_*) dF_\tau^n(v) dF_\tau^n(v_*) d\sigma d\tau \\ & \quad - \int_0^t \iint \Psi(v, v_*) dF_\tau(v) dF_\tau(v_*) d\sigma d\tau, \end{aligned}$$

where Ψ is the same as in the soft potential case and

$$\Psi_n(v, v_*) = |v - v_*|^\gamma \phi_c \left(\frac{|v - v_*|}{n} \right) \int_{\mathbb{S}^2} b_n \Delta \psi d\sigma.$$

Note that

$$\begin{aligned} |\Psi(v, v_*)| & \lesssim |v - v_*|^{\gamma+2-2\varepsilon} \lesssim \langle v \rangle^{2-\varepsilon} \langle v_* \rangle^{2-\varepsilon} \mathbf{1}_{\{|v| < 2|v_*| < 4|v|\}} \\ & \quad + \langle v \rangle^{\gamma+2-2\varepsilon} \mathbf{1}_{\{|v| \geq 2|v_*|\}} + \langle v_* \rangle^{\gamma+2-2\varepsilon} \mathbf{1}_{\{|v_*| \geq 2|v|\}}, \\ |\Psi_n(v, v_*)| & \lesssim \langle v \rangle^{2-\varepsilon} \langle v_* \rangle^{2-\varepsilon} \mathbf{1}_{\{|v| < 2|v_*| < 4|v|\}} \\ & \quad + \langle v \rangle^{2-2\varepsilon} \min\{\langle v \rangle^\gamma, n^\gamma\} \mathbf{1}_{\{|v| \geq 2|v_*|\}} + \langle v_* \rangle^{2-2\varepsilon} \min\{\langle v_* \rangle^\gamma, n^\gamma\} \mathbf{1}_{\{|v_*| \geq 2|v|\}}. \end{aligned}$$

Therefore, it follows from (4.16) in the next section that, if $t \leq T$ and $R > 1$, then we have

$$\begin{aligned} & \left| \int_0^t \iint_{\{|v| > R\} \cup \{|v_*| > R\}} \Psi_n(v, v_*) dF_\tau^n(v) dF_\tau^n(v_*) d\sigma d\tau \right| \\ & \lesssim TR^{-2\varepsilon} \left(\int \langle v \rangle^2 dF_0(v) \right)^2 + R^{-\varepsilon} \int_0^T \left(\int \langle v \rangle^{2-\varepsilon} \min\{\langle v \rangle^\gamma, n^\gamma\} dF_\tau^n(v) \right) d\tau \\ & \lesssim R^{-\varepsilon}. \end{aligned}$$

Thanks to (4.17), we have similarly

$$\left| \int_0^t \iint_{\{|v| > R\} \cup \{|v_*| > R\}} \Psi(v, v_*) dF_\tau(v) dF_\tau(v_*) d\sigma d\tau \right| \lesssim R^{-\varepsilon}.$$

Consequently, it suffices to show the convergence of

$$\begin{aligned}
& \int_0^t \iint_{\{|v| \leq R\} \cap \{|v_*| \leq R\}} \Psi_n(v, v_*) dF_\tau^n(v) dF_\tau^n(v_*) d\tau \\
& \quad - \int_0^t \iint_{\{|v| \leq R\} \cap \{|v_*| \leq R\}} \Psi(v, v_*) dF_\tau(v) dF_\tau(v_*) d\tau \\
& = \int_0^t \iint_{\{|v|, |v_*| \leq R\}} (\Psi_n(v, v_*) - \Psi(v, v_*)) dF_\tau^n(v) dF_\tau^n(v_*) d\tau \\
& \quad + \int_0^t \left(\iint_{\{|v|, |v_*| \leq R\}} \Psi(v, v_*) dF_\tau^n(v) dF_\tau^n(v_*) \right. \\
& \quad \quad \left. - \iint_{\{|v|, |v_*| \leq R\}} \Psi(v, v_*) dF_\tau(v) dF_\tau(v_*) \right) d\tau.
\end{aligned}$$

Estimations of the first and the second terms are quite the same as for the soft potentials. In the hard potential case it remains to show the energy conservation law and the moment gain property, which will be sent to the next section. \square

3.2. Infinite energy initial datum for the soft potential. For the proof of (2) of Theorem 1.2, it first show the following lemma, which is a variant of (i) of Lemma 3.2.

Lemma 3.3. *Let $-2 \leq \gamma < 0$ and let $c_{\gamma,s} < \alpha < 2$. If F_t^n is a solution in $C([0, \infty); P_{\alpha'})$, $0 < \alpha' < \alpha$, to the Cauchy problem (2.6) for the initial datum $F_0 \in P_\alpha$ (given by Theorem 2.5), then for any $T > 0$ there exists a $C_T > 0$ independent of n such that*

$$\sup_{t \in [0, T]} \int \langle v \rangle^\alpha dF_t^n \leq C_T \int \langle v \rangle^\alpha dF_0.$$

Proof. Using the weight function $W_\delta(v)$, it is easy to show $\int \langle v \rangle^\alpha dF_t^n < \infty$ for any fixed n . First consider the case $s \geq 1/2$. If $\psi_\alpha = \langle v \rangle^\alpha$ then we have

$$\begin{aligned}
\frac{d}{dt} \int \psi_\alpha dF_t^n(v) & = \frac{1}{2} \iiint B_n(\cdot) \Delta \psi_\alpha d\sigma dF_t^n(v) dF_t^n(v_*) \\
& = \frac{1}{2} \iiint_{|v-v_*| < 1} + \frac{1}{2} \iiint_{|v-v_*| \geq 1},
\end{aligned}$$

where $\Delta \psi_\alpha$ is defined by the same formula as (3.12), and it satisfies $\Delta \psi_\alpha = O(|v - v_*|^2 \theta^2)$ because $|\psi_\alpha''| \lesssim 1$. By using the mean value theorem to $\psi'_{\alpha,*} - \psi'_{\alpha,*}$ and $\psi'_\alpha - \psi_\alpha$, notice that $|\Delta \psi_\alpha| \lesssim |v - v_*| \theta (\langle v \rangle^{(\alpha-1)^+} + \langle v_* \rangle^{(\alpha-1)^+})$, where $a^+ = \max\{a, 0\}$ for $a \in \mathbb{R}$. Thus, we have

$$(3.14) \quad |\Delta \psi_\alpha| \lesssim \theta^{1+\delta} |v - v_*|^{1+\delta} \left(\langle v \rangle^{(\alpha-1)^+} + \langle v_* \rangle^{(\alpha-1)^+} \right)^{(1-\delta)}, \quad \forall \delta \in (0, 1].$$

When $\gamma + 2s \geq 1$, choose an $\varepsilon > 0$ such that $\alpha > \frac{\gamma}{2s-1+\varepsilon} + 2 > c_{\gamma,s} (> 1)$ and use (3.14) with $\delta = 2s - 1 + \varepsilon$ to the integral for $|v - v_*| \geq 1$. Using (3.14) with $\delta = 1$ also to the integral for $|v - v_*| < 1$, we see easily that there exists a $C > 0$ independent of n such that

$$\frac{d}{dt} \int \psi_\alpha dF_t^n(v) \leq C \int \psi_\alpha dF_t^n(v),$$

which leads to the desired estimate. When $\gamma + 2s < 1$, choose an $\varepsilon > 0$ such that $\gamma + 2s - 1 < -\varepsilon$ if the range of $\alpha > 1$ is handled. The range $1 \geq \alpha > c_{\gamma,s}$ is easier, by choosing an $\varepsilon > 0$ such that $\alpha > c_{\gamma,s} + \varepsilon$.

For the case $0 < s < 1/2$, we put simply $\Delta\psi_\alpha = \psi'_{\alpha,*} + \psi'_\alpha - \psi_{\alpha,*} - \psi_\alpha$. Then $\Delta\psi_\alpha = O(|v-v_*|^2\theta)$, $= O(|v-v_*|\theta(\langle v \rangle^{(\alpha-1)^+} + \langle v_* \rangle^{(\alpha-1)^+}))$ and $= O(\langle v \rangle^\alpha + \langle v_* \rangle^\alpha)$. The latter two yield

$$|\Delta\psi_\alpha| \lesssim \theta^\delta |v - v_*|^\delta \left(\langle v \rangle + \langle v_* \rangle \right)^{(1-\delta)\alpha + \delta(\alpha-1)^+}, \quad \forall \delta \in (0, 1].$$

from which the corresponding estimation is obvious. \square

By this lemma we see $F_t^n \in P_\alpha(\mathbb{R}^3)$, noting additionally the fact that v_j is collision invariant, when $\alpha > 1$. Using [19, (1.16)] for $\varphi^n(\xi, t) = \mathcal{F}(F_t^n)$, we have

$$(3.15) \quad \int \frac{|1 - \operatorname{Re} \varphi^n(\xi, t)|}{|\xi|^{3+\alpha}} d\xi = c_\alpha \int |v|^\alpha dF_t^n(v),$$

where $0 < c_\alpha = 2 \int \frac{\sin^2(\zeta \cdot \sigma/2)}{|\zeta|^{3+\alpha}} d\zeta$, $\sigma \in \mathbb{S}^2$. Since $\{\varphi^n(\xi, t)\}_{n=1}^\infty$ converges to $\varphi(\xi, t) = \mathcal{F}(F_t)$ uniformly in every compact subset of $\mathbb{R}^3 \times [0, \infty)$, it follows from [19, (1.16)] that for any $t \in (0, T]$

$$\begin{aligned} c_\alpha \int |v|^\alpha dF_t(v) &= \lim_{\delta \rightarrow +0} \int_{\delta \leq |\xi| \leq \delta^{-1}} \frac{|1 - \operatorname{Re} \varphi(\xi, t)|}{|\xi|^{3+\alpha}} d\xi \\ &= \lim_{\delta \rightarrow +0} \lim_{n \rightarrow \infty} \int_{\delta \leq |\xi| \leq \delta^{-1}} \frac{|1 - \operatorname{Re} \varphi^n(\xi, t)|}{|\xi|^{3+\alpha}} d\xi \\ &\leq C_T c_\alpha \int \langle v \rangle^\alpha dF_0(v), \end{aligned}$$

from which we can see Lemma 3.2 holds with the index 2 replaced by α . Thus, as in the proof of (1) of Theorem 1.2, we see that, for any $0 < \alpha' < \alpha$, $F_t \in C((0, \infty); P_{\alpha'}(\mathbb{R}^3))$ is a measure valued solution. Since F_t belongs to $P_\alpha(\mathbb{R}^3)$ uniformly with respect to t in any finite interval in $[0, \infty)$, we obtain

$$\left| \int \langle v \rangle^\alpha dF_t(v) - \int \langle v \rangle^\alpha dF_s(v) \right| = O(|t - s|)$$

by means of the same estimation for $\Delta\psi_\alpha$ as in the proof of Lemma 3.3, which shows $F_t \in C((0, \infty); P_\alpha(\mathbb{R}^3))$ on account of [27, Theorem 7.12]. Thus, the proof of (2) of Theorem 1.2 is completed.

4. POVZNER INEQUALITY AND THE MOMENT GAIN PROPERTY

4.1. Case: Initial datum F_0 with finite $2 + \kappa$ moment. In order to show the properties of moment gain and the moment propagation, we consider the case where F_0 satisfies, for $\kappa > 0$,

$$(4.1) \quad \int \langle v \rangle^{2+\kappa} dF_0 < \infty.$$

Let F_t^n denote the solution to the cutoff equation (3.1), for both hard and potential cases. Considering $W_\delta(v) = \langle v \rangle^{2+\kappa} \langle \delta v \rangle^{-2-\kappa}$ as in the proof of Theorem 2.5 and letting $\delta \rightarrow 0$, we have

$$(4.2) \quad \int \langle v \rangle^{2+\kappa} dF_t^n(v) < \infty \text{ for any } t > 0.$$

We will show the upper bound is independent of n , more precisely, for any fixed $T > 0$ there exists a $C_T > 0$ independent of n such that

$$(4.3) \quad \sup_{t \in [0, T]} \int_{\mathbb{R}_v^3} \langle v \rangle^{2+\kappa} dF_t^n(v) + \int_0^T \int_{\mathbb{R}_v^3} \langle v \rangle^{2+\kappa} \min\{\langle v \rangle^\gamma, n^\gamma\} dF_\tau^n(v) d\tau \leq C_T,$$

by using the Povzner inequality deeply studied in [15] (see also, the proof of Corollary 1.7 of [19]). Since $\sigma \in \mathbb{S}^2$, it can be written as

$$\sigma = \mathbf{k} \cos \theta + \sin \theta (\mathbf{h} \cos \varphi + \mathbf{i} \sin \varphi), \quad \theta \in [0, \pi), \varphi \in [-\pi, \pi),$$

by an orthogonal basis in \mathbb{R}^3 ,

$$\mathbf{k} = \frac{v - v_*}{|v - v_*|}, \quad \mathbf{i} = \frac{v \times v_*}{|v \times v_*|}, \quad \mathbf{h} = \mathbf{i} \times \mathbf{k} = \frac{((v - v_*) \cdot v)v_* - ((v - v_*) \cdot v_*)v}{|v - v_*||v \times v_*|}.$$

It follows from $(v + v_*) \perp \mathbf{i}$ and the definition of \mathbf{h} that

$$\begin{aligned} |v'|^2 &= \left| \frac{v + v_*}{2} \right|^2 + \left| \frac{v - v_*}{2} \right|^2 + \frac{|v - v_*|}{2} (v + v_*) \cdot \sigma \\ &= \frac{1}{4} (2|v|^2 + 2|v_*|^2) + \frac{|v - v_*|}{2} \left((v + v_*) \cdot (\cos \theta \mathbf{k} + \sin \theta \cos \varphi \mathbf{h}) \right) \\ &= \frac{1}{2} (|v|^2 + |v_*|^2) + \frac{\cos \theta}{2} (|v|^2 - |v_*|^2) \\ &\quad + \frac{\sin \theta \cos \varphi}{2|v \times v_*|} \left\{ (v + v_*) \cdot \left(((v - v_*) \cdot v)v_* - ((v - v_*) \cdot v_*)v \right) \right\} \\ &= \frac{|v|^2(1 + \cos \theta)}{2} + \frac{|v_*|^2(1 - \cos \theta)}{2} + |v||v_*| \sin \alpha \sin \theta \cos \varphi, \end{aligned}$$

where α is the angle between v and v_* . Therefore, we have

$$(4.4) \quad \begin{aligned} |v'|^2 &= |v|^2 \cos^2 \frac{\theta}{2} + |v_*|^2 \sin^2 \frac{\theta}{2} + |v \times v_*| \sin \theta \cos \varphi \\ &:= Y(\theta) + Z(\theta) \cos \varphi. \end{aligned}$$

It should be noted that

$$|Z(\theta)| = |(v - v_*) \times v_*| \sin \theta \leq |v - v_*| |v_*| \sin \theta,$$

which will be useful to absorb the singularity of the kinetic part in the soft potential case. Similarly, we have

$$(4.5) \quad \begin{aligned} |v_*'|^2 &= |v_*|^2 \cos^2 \frac{\theta}{2} + |v|^2 \sin^2 \frac{\theta}{2} - |v \times v_*| \sin \theta \cos \varphi \\ &= Y(\pi - \theta) - Z(\theta) \cos \varphi. \end{aligned}$$

If we set $\Psi(x) = \Psi_\kappa(x) = (1 + x)^{1+\kappa/2}$, it follows from the cutoff equation (3.1) that Furthermore we will show that To this end, we first notice

$$(4.6) \quad \frac{d}{dt} \int \langle v \rangle^{2+\kappa} dF_t^n(v) = \frac{1}{2} \iint \Phi_n(|v - v_*|) K_n(v, v_*) dF_t^n(v) dF_t^n(v_*),$$

where

$$\begin{aligned} K_n(v, v_*) &= \int_{\mathbb{S}^2} b_n \{ \Psi(|v'|^2) + \Psi(|v_*'|^2) - \Psi(|v|^2) - \Psi(|v_*|^2) \} d\sigma \\ &= 2 \int_0^\pi \int_0^\pi b_n(\cos \theta) \{ \Psi(|v'|^2) + \Psi(|v_*'|^2) - \Psi(|v|^2) - \Psi(|v_*|^2) \} \sin \theta d\theta d\varphi. \end{aligned}$$

It should be noted that the right hand side of (4.6) is well defined because of (4.2). Note that

$$\begin{aligned}
 & \int_0^\pi \Psi(Y(\theta) + Z(\theta) \cos \varphi) d\varphi \\
 &= \left(\int_0^{\frac{\pi}{2}} + \int_{\frac{\pi}{2}}^\pi \right) \Psi(Y(\theta) + Z(\theta) \cos \varphi) d\varphi \\
 &= \int_0^{\frac{\pi}{2}} \{ \Psi(Y(\theta) + Z(\theta) \cos \varphi) + \Psi(Y(\theta) - Z(\theta) \cos \varphi) - 2\Psi(Y(\theta)) \} d\varphi \\
 & \qquad \qquad \qquad + \pi\Psi(Y(\theta)),
 \end{aligned}$$

and by using integration by parts twice, we have

$$\begin{aligned}
 (4.7) \quad & \int_0^\pi \Psi(|v'|^2) d\varphi = \int_0^\pi \Psi(Y(\theta) + Z(\theta) \cos \varphi) d\varphi \\
 &= \pi\Psi(Y) + [\varphi\{\Psi(Y + Z \cos \varphi) + \Psi(Y - Z \cos \varphi) - 2\Psi(Y)\}]_0^{\frac{\pi}{2}} \\
 & \quad - \int_0^{\frac{\pi}{2}} \varphi\{\Psi'(Y + Z \cos \varphi) - \Psi'(Y - Z \cos \varphi)\}(-Z \sin \varphi) d\varphi \\
 &= \pi\Psi(Y) + \int_0^{\frac{\pi}{2}} Z\varphi \sin \varphi(\Psi'(Y + Z \cos \varphi) - \Psi'(Y - Z \cos \varphi)) d\varphi \\
 &= \pi\Psi(Y) + Z[(\sin \varphi - \varphi \cos \varphi)\{\Psi'(Y + Z \cos \varphi) - \Psi'(Y - Z \cos \varphi)\}]_0^{\frac{\pi}{2}} \\
 &+ Z^2 \int_0^{\frac{\pi}{2}} (\sin \varphi - \varphi \cos \varphi)\{\Psi''(Y + Z \cos \varphi) + \Psi''(Y - Z \cos \varphi)\} \sin \varphi d\varphi \\
 &= \pi\Psi(Y(\theta)) + Z^2 \int_0^{\frac{\pi}{2}} (\sin \varphi - \varphi \cos \varphi) \sin \varphi \\
 & \quad \times \{(\Psi''(Y(\theta) + Z \cos \varphi) + \Psi''(Y(\theta) - Z \cos \varphi))\} d\varphi.
 \end{aligned}$$

Similarly, we obtain

$$\begin{aligned}
 \int_0^\pi \Psi(|v_*'|^2) d\varphi &= \pi\Psi(Y(\pi - \theta)) + Z^2 \int_0^{\frac{\pi}{2}} (\sin \varphi - \varphi \cos \varphi) \sin \varphi \\
 & \quad \times \{(\Psi''(Y(\pi - \theta) + Z \cos \varphi) + \Psi''(Y(\pi - \theta) - Z \cos \varphi))\} d\varphi.
 \end{aligned}$$

In view of these formula, we can divide $K_n(v, v_*)$ into two parts

$$K_n(v, v_*) = -H_n(v, v_*) + G_n(v, v_*),$$

where it follows from the convexity of Ψ that

$$\begin{aligned}
 -H_n(v, v_*) &= 2\pi \int_0^\pi b_n(\cos \theta) \sin \theta \\
 & \times \left[\left\{ \Psi(|v|^2 \cos^2 \frac{\theta}{2} + |v_*|^2 \sin^2 \frac{\theta}{2}) - \cos^2 \frac{\theta}{2} \Psi(|v|^2) - \sin^2 \frac{\theta}{2} \Psi(|v_*|^2) \right\} \right. \\
 & \left. + \left\{ \Psi(|v_*|^2 \cos^2 \frac{\theta}{2} + |v|^2 \sin^2 \frac{\theta}{2}) - \cos^2 \frac{\theta}{2} \Psi(|v_*|^2) - \sin^2 \frac{\theta}{2} \Psi(|v|^2) \right\} \right] d\theta \leq 0.
 \end{aligned}$$

On the other hand, if $Z_0 = Z(\theta)/(1 + Y(\theta)) \in [0, 1]$, then

$$\begin{aligned} & Z^2 \int_0^{\frac{\pi}{2}} (\sin \varphi - \varphi \cos \varphi) \sin \varphi \{ \Psi''(Y + Z \cos \varphi) + \Psi''(Y - Z \cos \varphi) \} d\varphi \\ & \leq O(|\kappa|) Z^2 (1 + Y)^{\kappa/2-1} \int_0^{\pi/2} \varphi^3 \{ (1 + Z_0 \cos \varphi)^{\kappa/2-1} + (1 - Z_0 \cos \varphi)^{\kappa/2-1} \} d\varphi \\ & \begin{cases} \lesssim Z^2 \lesssim |v|^2 |v_*|^2 \theta^2, & \text{if } \kappa < 2; \\ \lesssim Z^2 (1 + Y)^{\kappa/2-1} \lesssim |v|^2 |v_*|^2 (\langle v \rangle^{\kappa-2} + \langle v_* \rangle^{\kappa-2}) \theta^2, & \text{if } \kappa \geq 2. \end{cases} \end{aligned}$$

Consequently, there exist $C_0 > 0$ and C_1 independent of n such that

$$G_n(v, v_*) \begin{cases} \leq C_0 |v|^2 |v_*|^2 \int_0^\pi b_n(\cos \theta) \sin^3 \theta d\theta \leq C_1 |v|^2 |v_*|^2 & \text{if } \kappa < 2; \\ \leq C_1 (|v_*|^\kappa \langle v \rangle^\kappa + |v|^2 \langle v_* \rangle^\kappa) & \text{if } \kappa \geq 2. \end{cases}$$

When $\gamma < 0$ we have

$$(4.8) \quad \begin{aligned} & \int_{\mathbb{R}^3} \langle v \rangle^{2+\kappa} dF_t^n(v) \\ & \leq \int_{\mathbb{R}^3} \langle v \rangle^{2+\kappa} dF_0(v) + C_3' \left(\int_{\mathbb{R}^3} \langle v \rangle^2 dF_0(v) \right) \int_0^t \int_{\mathbb{R}_v^3} \langle v \rangle^{\max\{2, \kappa\}} dF_\tau^n(v) d\tau, \end{aligned}$$

which shows that for any $T > 0$ there exists a $C_T > 0$ independent of n such that

$$\sup_{0 \leq t \leq T} \int_{\mathbb{R}^3} \langle v \rangle^{2+\kappa} dF_t^n(v) \leq C_T.$$

Taking the limit as $n \rightarrow \infty$, we obtain the property of the moment propagation of the solution F_t for soft potentials. .

From now on we consider the hard potential case. Note that there exist $[\theta_1, \theta_2] \subset (0, \pi/2)$ and $c_0 > 0$ independent of n such that $b_n(\cos \theta) \sin \theta \geq c_0$ on $[\theta_1, \theta_2]$. If $A, B > 0$, $A \neq B$ and if $F(x) = \Psi(A) + \Psi(B) - \Psi(Ax + B(1-x)) - \Psi(A(1-x) + Bx)$, then it follows from the convexity of Ψ that $F(x)$ takes the maximum at $x = 1/2$ and is decreasing with respect to $|x - 1/2|$. Therefore we have

$$\begin{aligned} H_n(v, v_*) & \geq 2\pi c_0 (\theta_2 - \theta_1) \left[\Psi(|v|^2) + \Psi(|v_*|^2) - \Psi(|v|^2 \cos^2 \frac{\theta_1}{2} + |v_*|^2 \sin^2 \frac{\theta_1}{2}) \right. \\ & \quad \left. - \Psi(|v_*|^2 \cos^2 \frac{\theta_1}{2} + |v|^2 \sin^2 \frac{\theta_1}{2}) \right], \end{aligned}$$

which shows that there exists a $C_2 > 0$ independent of n such that

$$(4.9) \quad H_n(v, v_*) \geq C_2 \left(\langle v \rangle^{2+\kappa} \mathbf{1}_{\{\langle v \rangle/2 \geq \langle v_* \rangle\}} + \langle v_* \rangle^{2+\kappa} \mathbf{1}_{\{\langle v_* \rangle/2 \geq \langle v \rangle\}} \right),$$

because, for $x_1 = \cos^2 \theta_1/2$ and $X = \langle v \rangle^2 / (\langle v \rangle^2 + \langle v_* \rangle^2)$,

$$\begin{aligned} & x_1 \Psi(|v|^2) + (1 - x_1) \Psi(|v_*|^2) - \Psi(x_1 |v|^2 + (1 - x_1) |v_*|^2) = (\langle v \rangle^2 + \langle v_* \rangle^2)^{1+\kappa/2} \\ & \quad \times \left\{ x_1 X^{1+\kappa/2} + (1 - x_1)(1 - X)^{1+\kappa/2} - \left(x_1 X + (1 - x_1)(1 - X) \right)^{1+\kappa/2} \right\} \end{aligned}$$

and $4/5 \leq X \leq 1$ if $\langle v \rangle \geq 2\langle v_* \rangle$.

Note that $\int \langle v \rangle^2 dF_t^n(v) = \int \langle v \rangle^2 F_0(v)$ and $\int_{|v| \leq R} dF_t^n(v) \geq 1/2$ for a sufficiently large $R > 0$ independent of n . Integrating the equation (4.6), we have

$$(4.10) \quad \begin{aligned} & \int_{\mathbb{R}^3} \langle v \rangle^{2+\kappa} dF_t^n(v) + C'_2 \int_0^t \int_{\mathbb{R}^3} \langle v \rangle^{2+\kappa} \min\{\langle v \rangle^\gamma, n^\gamma\} dF_\tau^n(v) d\tau \\ & \leq \int_{\mathbb{R}^3} \langle v \rangle^{2+\kappa} dF_0(v) \\ & \quad + C'_3 \left(\int_{\mathbb{R}^3} \langle v \rangle^2 dF_0(v) \right) \int_0^t \int_{\mathbb{R}^3} \langle v \rangle^{\max\{2, \kappa\}} \min\{\langle v \rangle^\gamma, n^\gamma\} dF_\tau^n(v) d\tau, \end{aligned}$$

for suitable positive constants C'_2, C'_3 independent of n , which yields (4.3) since, by taking R_0 such that $C'_2 R_0^{\min\{2, \kappa\}} \geq 2C'_3 \int_{\mathbb{R}^3} \langle v \rangle^2 dF_0(v)$, we have

$$(4.11) \quad \begin{aligned} & C'_2 \int_{|v| \geq R_0} \langle v \rangle^{2+\kappa} \min\{\langle v \rangle^\gamma, n^\gamma\} dF_\tau^n(v) \\ & \geq 2C'_3 \left(\int_{\mathbb{R}^3} \langle v \rangle^2 dF_0(v) \right) \int_{|v| \geq R_0} \langle v \rangle^{\max\{2, \kappa\}} \min\{\langle v \rangle^\gamma, n^\gamma\} dF_\tau^n(v). \end{aligned}$$

Thanks to (4.3), we obtain

$$(4.11) \quad \sup_{t \in [0, T]} \int_{\mathbb{R}^3} \langle v \rangle^{2+\kappa} dF_t(v) + \int_0^T \int_{\mathbb{R}^3} \langle v \rangle^{2+\kappa+\gamma} dF_\tau(v) d\tau \leq C_T,$$

by taking the limit. By using (4.3) and (4.11) we can show that F_t is a weak measure valued solution in the sense of (1.7) when the initial datum satisfies (4.1), similar as in the preceding section.

4.2. Case: the initial datum $F_0 \in P_2(\mathbb{R}^3)$. We recall Proposition A1 of [15].

Proposition 4.1. *Let F belong to $P_2(\mathbb{R}^3)$. Then there exists a concave function $\rho(r) \in C^2$, depending on F , such that $\rho(r) \rightarrow \infty$ as $r \rightarrow \infty$, $r\rho(r)$ is convex,*

$$(4.12) \quad r\rho'(r) \lesssim \log r,$$

and such that for all $\varepsilon > 0$ and $\alpha \in (0, 1)$, $(\rho(r) - \rho(\alpha r))r^\varepsilon \rightarrow \infty$ as $r \rightarrow \infty$, and such that

$$(4.13) \quad \int \langle v \rangle^2 \rho(|v|^2) dF(v) < \infty.$$

Proof. The proposition is the same as Proposition A1 of [15], except for the fact that F is a probability measure and (4.12) is explicitly stated. Without loss of generality we may assume that $\int_{\mathbb{R}^3} \langle v \rangle^2 dF(v) = 2$. Take $0 = r_0 < r_1 < \dots < r_j < \dots$ such that

$$\int_{|v| \geq r_j} \langle v \rangle^2 dF(v) \leq 2^{-(j-1)}.$$

One can take $\{r_j\}_{j=1}^\infty$ such that $r_j > e^j$ and $r_{j+1} - r_j > r_j - r_{j-1}$. Let $\rho_1(r)$ be linear in each $[r_j, r_{j+1})$ and such that $\rho_1(r_j) = j$. It is obvious that $\rho_1(r)$ is concave, increasing to infinity with r , and

$$\int_{\mathbb{R}^3} \langle v \rangle^2 \rho_1(|v|^2) dF(v) \leq \sum_{j=1}^\infty 2^{-(j-1)} (j+1) < \infty.$$

Finally we put

$$\rho(r) = \frac{1}{r} \int_0^r \frac{1}{y} \int_0^y \left(\rho_1(z) + 1 - (\log(e+z))^{-1} \right) dz dy$$

in order to make $r\rho(r)$ convex. Since $\rho(r) \leq \rho_1(r) + 1$, we have (4.13), and moreover (4.12) follows from $\rho_1(e^j) \leq j$. The rest of the proof is quite the same as the one of Proposition A1 of [15]. \square

Applying this proposition to the initial datum $F_0 \in P_2(\mathbb{R}^3)$, we have

$$\int \langle v \rangle^2 \rho(|v|^2) dF_0(v) < \infty,$$

and moreover for each fixed n and $T > 0$

$$\sup_{0 \leq t \leq T} \int \langle v \rangle^2 \rho(|v|^2) dF_t^n(v) < \infty,$$

by considering $W_\delta(v) = \langle v \rangle^2 \rho(|v|^2) (\langle \delta v \rangle^2 \rho(\delta^2 |v|^2))^{-1}$ as in the proof of Theorem 2.5 because it follows from the concavity of $\rho(r)$ that $\rho(2r) + \rho(0) \leq 2\rho(r)$, so that

$$W_\delta(v') \lesssim W_\delta(v) + W_\delta(v_*).$$

We follow the proof of Lemma 2.2 iii) of [15], in noticing that our $b(\cos \theta)$ is unbounded. Replace $\Psi_\kappa(x)$ in the preceding procedure by $\Psi(x) = x\rho(x)$. Then

$$\begin{aligned} \Psi(|v'|^2) &= \left(Y(\theta) + Z(\theta) \cos \varphi \right) \rho(Y(\theta) + Z(\theta) \cos \varphi) \\ &\leq \left(Y(\theta) + Z(\theta) \cos \varphi \right) \left(\rho(Y(\theta)) + \rho'(Y(\theta)) Z(\theta) \cos \varphi \right), \end{aligned}$$

because the first factor is non-negative and ρ is concave. Therefore, we have

$$\int_0^\pi \Psi(|v'|^2) d\varphi \leq \pi \Psi(Y(\theta)) + \frac{\pi}{2} Z(\theta)^2 \rho'(Y(\theta)),$$

instead of (4.7). Similarly we have

$$\int_0^\pi \Psi(|v_*'|^2) d\varphi \leq \pi \Psi(Y(\pi - \theta)) + \frac{\pi}{2} Z(\theta)^2 \rho'(Y(\pi - \theta)).$$

Since $|Z| \leq Y$, it follows from (4.12) that

$$\begin{aligned} (4.14) \quad G_n(v, v_*) &\leq C_0 |v|^{2-2\varepsilon} |v_*|^{2-2\varepsilon} \int_0^\pi \left(\frac{|Z|}{Y} \right)^{2\varepsilon} b_n(\cos \theta) \sin^{3-2\varepsilon} \theta d\theta \\ &\leq C_1 |v|^{2-2\varepsilon} |v_*|^{2-2\varepsilon}. \end{aligned}$$

Notice again that there exist $[\theta_1, \theta_2] \subset (0, \pi/2)$ and $c_0 > 0$ independent of n such that $b_n(\cos \theta) \sin \theta \geq c_0$ on $[\theta_1, \theta_2]$, and $F(x) = \Psi(A) + \Psi(B) - \Psi(Ax + B(1-x)) - \Psi(A(1-x) + Bx)$ takes the maximum at $x = 1/2$ and is decreasing with respect

to $|x - 1/2|$. Therefore, if $\tau = \sin^2 \theta_1/2 \in (0, 1)$ and if $|v_*| \geq 2|v|$, we have

$$\begin{aligned}
H_n(v, v_*) &\geq 2\pi \int_{\theta_1}^{\theta_2} b_n(\cos \theta) \sin \theta \left[\Psi(|v_*|^2) + \Psi(|v|^2) \right. \\
&\quad \left. - \Psi(|v|^2 \cos^2 \frac{\theta}{2} + |v_*|^2 \sin^2 \frac{\theta}{2}) - \Psi(|v_*|^2 \cos^2 \frac{\theta}{2} + |v|^2 \sin^2 \frac{\theta}{2}) \right] d\theta \\
&\geq 2\pi(\theta_2 - \theta_1)c_0 \left[|v|^2 \rho(|v|^2) + |v_*|^2 \rho(|v_*|^2) \right. \\
&\quad \left. - \left((1-\tau)|v|^2 + \tau|v_*|^2 \right) \rho\left((1-\tau)|v|^2 + \tau|v_*|^2 \right) \right. \\
&\quad \left. - \left((1-\tau)|v_*|^2 + \tau|v|^2 \right) \rho\left((1-\tau)|v_*|^2 + \tau|v|^2 \right) \right] \\
&\geq 2\pi(\theta_2 - \theta_1)c_0 \left[|v_*|^2 \rho(|v_*|^2) - (|v|^2 + |v_*|^2) \rho(\alpha|v_*|^2) \right],
\end{aligned}$$

where $\alpha = \max\{\frac{1+3\tau}{4}, 1 - \frac{3\tau}{4}\}$, because ρ is increasing and $(1-\tau)|v|^2 + \tau|v_*|^2 \leq (1+3\tau)|v_*|^2/4$. Since $|v|^2 \rho(|v_*|^2) \lesssim |v|^{2-2\varepsilon} |v_*|^{2-2\varepsilon}$ when $|v_*| \geq 2|v|$ and since the similar result holds when v and v_* are exchanged, it follows from Proposition 4.1 that there exist constants $C_2, C_3 > 0$ independent of n such that

$$(4.15) \quad C_2 \left(|v|^{2-\varepsilon} \mathbf{1}_{\{|v|/2 \geq |v_*|\}} + |v_*|^{2-\varepsilon} \mathbf{1}_{\{|v_*|/2 \geq |v|\}} \right) \leq H_n(v, v_*) + C_3 |v|^{2-2\varepsilon} |v_*|^{2-2\varepsilon},$$

which, together with (4.14), concludes

$$(4.16) \quad \sup_{t \in [0, T]} \int_{\mathbb{R}_v^3} \langle v \rangle^2 \rho(|v|^2) dF_t^n(v) + \int_0^T \int_{\mathbb{R}_v^3} \langle v \rangle^{2-\varepsilon} \min\{\langle v \rangle^\gamma, n^\gamma\} dF_\tau^n(v) d\tau \leq C_T,$$

and by taking $n \rightarrow \infty$, we get

$$(4.17) \quad \sup_{t \in [0, T]} \int_{\mathbb{R}_v^3} \langle v \rangle^2 \rho(|v|^2) dF_t(v) + \int_0^T \int_{\mathbb{R}_v^3} \langle v \rangle^{2-\varepsilon+\gamma} dF_\tau(v) d\tau \leq C_T,$$

similar to (4.11). As proved in the preceding section, the above two estimates show the existence of the weak solution for the initial datum $F_0 \in P_2(\mathbb{R}^3)$ in the hard potential case.

The proof of the energy conservation law (1.8) for the hard potentials is due to the reverse Povzner inequality, as stated in p.293 of [25]. If one considers $\Psi_\kappa(x) = (1+x)^{1+\kappa/2}$ with $\kappa < 0$, then it follows from the concavity of Ψ_κ that $-H(v, v_*) \geq 0$ and

$$G(v, v_*) \leq O(|\kappa|) |v|^2 |v_*|^2,$$

so that

$$\int \langle v \rangle^{2+\kappa} dF_t(v) \geq \int \langle v \rangle^{2+\kappa} dF_0(v) - O(|\kappa|) \left(\int |v|^2 dF_0(v) \right)^2,$$

where H, G are defined in the subsection 4.1 with b_n there replaced by b . Letting $\kappa \rightarrow -0$, we get (1.8), since the opposite inequality is obvious.

Finally, we remark the existence of a weak solution with moments gain of any order at a fixed $T > 0$ for the initial datum $F_0 \in P_2(\mathbb{R}^3)$. Indeed, by (4.17) we see that

$$\int_{\mathbb{R}_v^3} \langle v \rangle^{2-\varepsilon+\gamma} dF_\tau(v) < \infty$$

for almost all $\tau \in [0, T]$. Choose a $t_1 > 0$ arbitrarily close to 0 as an initial time. Then we get

$$\int_{t_1}^T \int_{\mathbb{R}^3} \langle v \rangle^{2-\varepsilon+2\gamma} dF_\tau(v) d\tau < \infty.$$

We can repeat the same procedure again, and so on, such that

$$0 < t_1 < t_2 < \dots < t_n < \dots < T.$$

Using the method in the subsection 4.1 (propagation of the moments), we see that for any $\ell > 0$

$$\int \langle v \rangle^\ell dF_T(v) < \infty.$$

5. SMOOTHING EFFECT OF THE SINGULAR CROSS SECTION

We start the following variant of Theorem 1.2 in Fournier [9].

Proposition 5.1. *Let $B(\cdot)$ satisfy (1.2)-(1.4) and let $-2 \leq \gamma \leq 2$. If $\gamma \leq 0$ and if $F_0 \in P_2(\mathbb{R}^3)$ is not a single Dirac mass, then for any weak solution in the sense of Definition 1.1 to the Cauchy problem (1.1)-(1.5) we have*

$$\text{supp } F_t = \mathbb{R}^3 \text{ for all } t > 0.$$

When $\gamma > 0$, if, for $\kappa > \gamma$, F_t satisfies

$$\forall T > 0, \exists C_T > 0 \text{ such that } \sup_{t \in [0, T]} \int_{\mathbb{R}^3} \langle v \rangle^{2+\kappa} dF_t(v) \leq C_T$$

then we have the same conclusion.

Proof. We first consider the case $\gamma \leq 0$. For $\psi(v) \in C_b^2(\mathbb{R}^3)$ we see that the map

$$t \mapsto \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} \int_{\mathbb{S}^2} b(\cdot) |v - v_*|^\gamma (\psi(v'_*) + \psi(v') - \psi(v_*) - \psi(v)) dF_t(v) dF_t(v_*) d\sigma$$

belongs to $C([0, \infty))$ under the assumption $F_0 \in P_2(\mathbb{R}^3)$. In fact, since $2 + \gamma \leq 2$, it follows from (3.13) and (3.5) that, for any $\eta > 0$ there exists $R > 0$ independent of t such that

$$\left| \iint_{\{|v| \geq R\} \cup \{|v_*| \geq R\}} \int_{\mathbb{S}^2} b(\cdot) |v - v_*|^\gamma (\psi'_* + \psi' - \psi_* - \psi) dF_t(v) dF_t(v_*) d\sigma \right| < \eta.$$

Therefore, the similar way to (3.8) leads us to the desired continuity, from which and (1.7) it follows that the map $t \mapsto \int_{\mathbb{R}^3} \psi(v) dF_t$ belongs to $C^1([0, \infty))$. When $\gamma > 0$ we have the same C^1 regularity under the above additional condition.

The rest of the proof is the similar as the one of [9, Theorem 1.2]. Indeed, it follows from conservations of momentum and energy in Definition 1.1 that

$$\int_{\mathbb{R}^3} |v - v_0|^2 dF_t(v) = \int_{\mathbb{R}^3} |v - v_0|^2 dF_0(v) > 0, \quad \forall v_0 \in \mathbb{R}^3,$$

which just implies Step 1 in the proof of [9, Theorem 1.2], that is, F_t is not a single Dirac mass for all $t \geq 0$. Let $B(v_0, \varepsilon)$ be an open ball with a radius $\varepsilon > 0$ centered

at a $v_0 \in \mathbb{R}^3$ and suppose that $\int_{B(v_0, \varepsilon)} dF_t(v) = 0$ for a $t > 0$. Then, by (1.7) we have, for any $0 \leq \psi \in C_0^2(B(v_0, \varepsilon))$,

$$\begin{aligned} & \iint_{\mathbb{R}^3 \times \mathbb{R}^3} \int_{\mathbb{S}^2} b(\cdot) |v - v_*|^\gamma (\psi(v'_*) + \psi(v')) dF_t(v) dF_t(v_*) d\sigma \\ & = 2\partial_t \left(\int_{\mathbb{R}^3} \psi(v) dF_t(v) \right) = 0, \end{aligned}$$

because the map $s \mapsto \int_{\mathbb{R}^3} \psi(v) dF_s$ belongs to $C^1([0, \infty))$ and $\int_{\mathbb{R}^3} \psi(v) dF_s \geq 0$. Since $b > 0$ in $(0, \pi/2)$ and $0 \leq \psi \in C_0^2(B(v_0, \varepsilon))$ can be taken arbitrarily, we see that

$$\begin{aligned} & \iint_{\mathbb{R}^3 \times \mathbb{R}^3} \int_{\mathbb{S}^2} \mathbf{1}_{\{v'(v, v_*, \sigma) \in B(v_0, \varepsilon)\}} \\ & \quad \times \mathbf{1}_{\{v \neq v_*, 0 < (v - v_*) \cdot \sigma < |v - v_*|\}} dF_t(v) dF_t(v_*) d\sigma = 0. \end{aligned}$$

Thus, Step 2 is also similar. Other Steps are quite the same as in the proof of [9, Theorem 1.2]. \square

The following coercivity estimate estimate is a key for the proof of Theorem 1.4.

Proposition 5.2. *Let $B(\cdot)$ satisfy (1.2)-(1.4) and assume $2 \geq \gamma > -2s$. Let $F_0 \in P_2(\mathbb{R}^3)$ be not equal to a single Dirac mass. For any $t_0 > 0$ there exists a weak solution $F_t \in C([0, \infty); P_2(\mathbb{R}^3))$ to the Cauchy problem (1.1)-(1.5) satisfying the following; there exist $T > 0$ and $c_0, C > 0$ such that for any $f \in \mathcal{S}(\mathbb{R}^3)$ and any $t \in [t_0, t_0 + T]$*

$$(5.1) \quad \begin{aligned} -\left(Q(F_t, f), f\right) &= \iiint \Phi b (f^2 - f'f) dv d\sigma dF_t(v_*) \\ &\geq c_0 \|\langle v \rangle^{\gamma/2} f\|_{H^s}^2 - C \|\langle v \rangle^{\gamma/2} f\|_{L^2}^2. \end{aligned}$$

Proof. When $\gamma > 0$, in view of the moment gain property, we may assume F_t satisfies (4.11) with an interval $[0, T]$ replaced by $[t_0/2, T + t_0]$ with $\kappa > \gamma$. It follows from Proposition 5.1 that

$$(5.2) \quad \text{supp } F_{t_0} = \mathbb{R}^3.$$

For the sake of the simplicity of the notation, we write 0 instead of t_0 , in what follows.

For $G \in P_2$, put

$$\mathcal{C}_\gamma(G, f) = \iiint_{\mathbb{R}^3 \times \mathbb{R}^3 \times \mathbb{S}^2} b(\cdot) |v - v_*|^\gamma (f' - f)^2 dv d\sigma dG(v_*),$$

and note that

$$(Q(G, f), f) = -\frac{1}{2} \mathcal{C}_\gamma(G, f) + \frac{1}{2} \iiint \Phi b (f'^2 - f^2) dv d\sigma dG(v_*).$$

It follows from the Cancellation Lemma and Remark 6 in [1] that

$$(5.3) \quad \begin{aligned} & \left| \iiint b |v - v_*|^\gamma (f^2 - f'^2) dv d\sigma dG(v_*) \right| \lesssim \left| \iiint |v - v_*|^\gamma f^2 dv dG(v_*) \right| \\ & \lesssim \int \langle v_* \rangle^{|\gamma|} dG(v_*) \|f\|_{H_{\gamma/2}^{(-\gamma/2)^+}}^2, \end{aligned}$$

where the last inequality in the case $\gamma \geq 0$ is trivial. While for $\gamma < 0$, this follows from the fact that

$$(5.4) \quad |v - v_*|^\gamma \lesssim \langle v \rangle^\gamma \{ \mathbf{1}_{|v-v_*| \geq \langle v \rangle / 2} + \mathbf{1}_{|v-v_*| < \langle v \rangle / 2} \langle v_* \rangle^{-\gamma} |v - v_*|^\gamma \},$$

and the Hardy inequality $\sup_{v_*} \int |v - v_*|^\gamma |g(v)|^2 dv \lesssim \|g\|_{H^{-\gamma/2}}^2$.

For the proof of the proposition, it now suffices to consider only the quantity $\mathcal{C}_\gamma(G, f)$ with $G = F_t$ because one can apply the interpolation inequality

$$\|f\|_{H^{\langle -\gamma/2 \rangle +}}^2 \leq \varepsilon \|f\|_{H^s}^2 + C_\varepsilon \|f\|_{L^2}^2$$

to (5.3) in view of $\gamma > -2s$. The case $\gamma = 0$ is easy. In fact, it follows from Proposition 1 of [1] that

$$\begin{aligned} (2\pi)^3 \mathcal{C}_0(G, f) &= \iint b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \psi(0) (|h(\xi)|^2 + |h(\xi^+)|^2) d\sigma d\xi \\ &\quad - 2\operatorname{Re} \iint b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) \psi(\xi^-) h(\xi^+) \overline{h(\xi)} d\sigma d\xi \\ &\geq \int \left(\int b\left(\frac{\xi}{|\xi|} \cdot \sigma\right) (\psi(0) - |\psi(\xi^-)|) d\sigma \right) |h(\xi)|^2 d\xi \end{aligned}$$

if $\psi(\xi) = \mathcal{F}(G)$, $h(\xi) = \mathcal{F}(f)$. Consider $G = F_t$, $t \in [0, T]$. Since $\operatorname{supp} F_0 = \mathbb{R}^3$, we may apply the method in the subsection 2.1 of [20], developed under the condition that the initial measure F_0 is not concentrated on the straight line. Then we have the desired estimate for $0 \leq t \leq T$ with a sufficiently small $T > 0$, by using the continuity of $\psi(t, \xi) = \mathcal{F}(F_t)$.

We now consider the case $\gamma \neq 0$, following the argument used in the proof of Proposition 2.1 of [4] (see also the one of Lemma 2 of [1]). Set $B(R) = \{v \in \mathbb{R}^3; |v| \leq R\}$ for $R > 0$ and $B_0(R, r) = \{v \in B(R); |v - v_0| \geq r\}$ for a $v_0 \in \mathbb{R}^3$ and $r \geq 0$. For a moment, we concentrate to handle $\mathcal{C}_\gamma(F_0, f)$. By means of $\operatorname{supp} F_0 = \mathbb{R}^3$, for any positive constants $R > 1 > r_0$ we see that

(5.5) the positive measure $\chi_{B_0(R, r_0)} F_0$ is not concentrated on the straight line,

where χ_A denotes a characteristic function of the set $A \subset \mathbb{R}^3$. Let $\varphi_R(v)$ be a non-negative smooth function not greater than one, which is 1 for $|v| \geq 4R$ and 0 for $|v| \leq 2R$. In view of

$$\frac{\langle v \rangle}{4} \leq |v - v_*| \leq 2\langle v \rangle \quad \text{on } \operatorname{supp} (\chi_{B(R)})_* \varphi_R,$$

we have

$$\begin{aligned} 4^{|\gamma|} \Phi(|v - v_*|) (f' - f)^2 &\geq (\chi_{B(R)})_* (\langle v \rangle^{\gamma/2} \varphi_R)^2 (f' - f)^2 \\ &\geq (\chi_{B(R)})_* \left[\frac{1}{2} \left((\langle v \rangle^{\gamma/2} \varphi_R f)' - \langle v \rangle^{\gamma/2} \varphi_R f \right)^2 - \left((\langle v \rangle^{\gamma/2} \varphi_R)' - \langle v \rangle^{\gamma/2} \varphi_R \right)^2 f'^2 \right]. \end{aligned}$$

It follows from the mean value theorem that for a $\tau \in (0, 1)$

$$\begin{aligned} \left| (\langle v \rangle^{\gamma/2} \varphi_R)' - \langle v \rangle^{\gamma/2} \varphi_R \right| &\lesssim \langle v + \tau(v' - v) \rangle^{\gamma/2 - 1} |v - v_*| \sin \frac{\theta}{2} \\ &\lesssim \langle v_* \rangle^{|\gamma/2 - 1|} \langle v' - v_* \rangle^{\gamma/2} \sin \frac{\theta}{2} \\ &\lesssim \langle v_* \rangle^{|\gamma/2| + |\gamma/2 - 1|} \langle v' \rangle^{\gamma/2} \sin \frac{\theta}{2}, \end{aligned}$$

because $|v - v_*|/\sqrt{2} \leq |v' - v_*| \leq |v + \tau(v' - v) - v_*| \leq |v - v_*|$ for $\theta \in [0, \pi/2]$. Therefore, we have

$$(5.6) \quad \mathcal{C}_\gamma(F_t, f) \geq 2^{-1-2|\gamma|} \mathcal{C}_0(\chi_{B(R)} F_t, \varphi_R \langle v \rangle^{\gamma/2} f) - C_R \int dF_t(v_*) \|f\|_{L^2_{\gamma/2}}^2,$$

for a positive constant $C_R \sim R^{|\gamma|+|\gamma-2|}$. For a set $B(4R)$ we take a finite covering

$$B(4R) \subset \bigcup_{v_j \in B(4R)} A_j, \quad A_j = \{v \in \mathbb{R}^3; |v - v_j| \leq \frac{r_0}{4}\}.$$

For each A_j we choose a non-negative smooth function φ_{A_j} which is 1 on A_j and 0 on $\{|v - v_j| \geq r_0/2\}$. Note that

$$\frac{r_0}{2} \leq |v - v_*| \leq 6R \text{ on } \text{supp}(\chi_{B_j(R, r_0)})_* \varphi_{A_j}.$$

Then we have

$$\begin{aligned} \Phi(|v - v_*|)(f' - f)^2 &\gtrsim \min\{r_0^{\gamma^+}, R^{-(\gamma)^+}\} (\chi_{B_j(R, r_0)})_* \varphi_{A_j}^2 (f' - f)^2 \\ &\gtrsim R^{-\gamma^+} \min\{r_0^{\gamma^+}, R^{-(\gamma)^+}\} (\chi_{B_j(R, r_0)})_* \\ &\quad \times \left[\frac{1}{2} \left((\langle v \rangle^{\gamma/2} \varphi_{A_j} f)' - \langle v \rangle^{\gamma/2} \varphi_{A_j} f \right)^2 - \left((\langle v \rangle^{\gamma/2} \varphi_{A_j})' - \langle v \rangle^{\gamma/2} \varphi_{A_j} \right)^2 f'^2 \right]. \end{aligned}$$

Since $\left| (\langle v \rangle^{\gamma/2} \varphi_{A_j})' - \langle v \rangle^{\gamma/2} \varphi_{A_j} \right| \lesssim R^{|\gamma|+1} \langle v' \rangle^\gamma \sin \theta/2$ if $|v_*| \leq R$, we obtain

$$(5.7) \quad \mathcal{C}_\gamma(F_t, f) \gtrsim \min\{(r_0/R)^{\gamma^+}, R^{-(\gamma)^+}\} \mathcal{C}_0(\chi_{B_j(R, r_0)} F_t, \varphi_{A_j} \langle v \rangle^{\gamma/2} f) - C'_{R, r_0} \int dF_t(v_*) \|f\|_{L^2_{\gamma/2}}^2,$$

for a positive constant $C'_{R, r_0} \sim R^{2+2|\gamma|}$.

Notice that we may replace $\chi_{B(R)}$ and $\chi_{B_j(R, r_0)}$ in (5.6) and (5.7), respectively, by C_0^∞ functions $\tilde{\chi}_j$ which still satisfy (5.5). Then, from the argument in the subsection 2.1 of [20] it follows that

$$\exists \kappa_0 > 0 \text{ such that } \left| \widehat{\tilde{\chi}_j F_0}(0) - \left| \widehat{\tilde{\chi}_j F_0}(\xi^-) \right| \right| \geq \kappa_0$$

if ξ^- belongs to a suitable compact set (see (2.3) of [20]). Since for $\varphi(t, \xi) = \int e^{-iv \cdot \xi} dF_t(v)$ we have

$$\widehat{\tilde{\chi}_j F_t}(\xi) = \int_{\mathbb{R}^3} \varphi(t, \xi - \eta) \widehat{\tilde{\chi}_j}(\eta) d\eta,$$

it follows from the uniform continuity of $\varphi(t, \xi)$ on any fixed compact set and $|\varphi| \leq 1$ that there exists a $T > 0$ such that

$$\left| \widehat{\tilde{\chi}_j F_t}(0) - \left| \widehat{\tilde{\chi}_j F_t}(\xi^-) \right| \right| \geq \frac{\kappa_0}{2} \text{ for any } t \in [0, T]$$

if ξ^- belongs to the same suitable compact set stated above. Therefore we obtain

$$\mathcal{C}_0(\tilde{\chi}_j F_t, f) + \|f\|_{L^2}^2 \gtrsim \|f\|_{H^s}^2,$$

which, together with (5.6) and (5.7), shows that there exist $c'_0, C, C' > 0$ such that

$$(5.8) \quad \begin{aligned} \mathcal{C}_\gamma(F_t, f) &\geq c'_0 \left(\|\langle D \rangle^s \varphi_R \langle v \rangle^{\gamma/2} f\|^2 + \sum_j \|\langle D \rangle^s \varphi_{A_j} \langle v \rangle^{\gamma/2} f\|^2 \right) - C \|f\|_{L^2_{\gamma/2}}^2 \\ &\geq c'_0 \|\langle v \rangle^{\gamma/2} f\|_{H^s}^2 - C' \|f\|_{L^2_{\gamma/2}}^2, \end{aligned}$$

because $\varphi_R^2 + \sum_j \varphi_{A_j}^2 \geq 1$ and commutators $[(D)^s, \varphi_R]$, $[(D)^s, \varphi_{A_j}]$ are L^2 bounded operators.

In the soft potential case $\gamma < 0$, noting

$$\Phi(|v - v_*|)(f' - f)^2 \gtrsim R^{-|\gamma|} (\chi_{B(R)})_* (1 - \varphi_{2R})^2 (f' - f)^2,$$

we can replace $\chi_{B_0(R, r_0)}$ in the condition (5.5) by $\chi_{B(R)}$, because one $1 - \varphi_{2R}$ is enough to cover the set B_{4R} , instead of finite many φ_{A_j} . \square

Proof of Theorem 1.4. As in [4, 17], we use the mollifier as follows: Let $\lambda, N_0 \in \mathbb{R}$, $\delta > 0$ and put

$$(5.9) \quad M_\lambda^\delta(\xi) = \frac{\langle \xi \rangle^\lambda}{(1 + \delta \langle \xi \rangle)^{N_0}}, \quad \langle \xi \rangle = (1 + |\xi|^2)^{1/2}.$$

Then $M_\lambda^\delta(\xi)$ belongs to the symbol class $S_{1,0}^{\lambda-N_0}$ of pseudo-differential operators for each fixed $\delta > 0$ and belongs to $S_{1,0}^\lambda$ uniformly with respect to $\delta \in (0, 1]$.

We consider the case $\gamma > 0$. Since we can take arbitrarily small $t_0 > 0$ in both Proposition 5.2 and (1.9) with a fixed large ℓ , for the brevity we may assume $t_0 = 0$ there. Noting $P_\ell(\mathbb{R}^3) \subset H_\ell^{-3/2-\varepsilon}(\mathbb{R}^3)$ for $\forall \varepsilon > 0$, we can apply the proof of Theorem 5.1 in [4]. Namely it suffices to show, for any $0 < t_1 \leq T$ that

$$(5.10) \quad dF_t(v) = f(t, v)dv, \quad f(t, \cdot) \in L^\infty([t_1, T]; L_\ell^2(\mathbb{R}^3)).$$

As in the proof of Theorem 4.1 of [4], we show this by induction. Assume that for $0 > a \geq -3/2 - \varepsilon$, we have

$$\sup_{[0, T]} \|f(t, \cdot)\|_{H_\ell^a} < \infty.$$

It should be noted that if $a < -3/2$ then it follows from (1.9) that for $t \in [0, T]$

$$\|f(t, \cdot)\|_{H_\ell^a}^2 \leq \int \langle \xi \rangle^{2a} d\xi \left(\int (1 + |v|^\ell) dF_t(v) \right)^2 < \infty.$$

Put $\lambda(t) = Nt + a$ for $N > 0$. When $0 < s \leq 1/2$, choose $N_0 = a + (5 + \gamma)/2 \geq 1 - \varepsilon + (\gamma/2) > 0$ such that [4, (3.8)] is fulfilled. Put $\varepsilon_0 = (1 - 2s')/8 > 0$ and consider $\varepsilon = \varepsilon_0$, where $0 < s' < s$ is chosen to satisfy $\gamma + 2s' > 0$. If we choose $N, T_1 > 0$ such that $NT_1 = \varepsilon_0$ then

$$s + \lambda(T_1) - N_0 - a = s + \varepsilon_0 - N_0 \leq s - 1 + 2\varepsilon_0 - (\gamma/2) < (s' - 1/2) + 2\varepsilon_0 < 0,$$

which shows

$$(5.11) \quad M_{\lambda(t)}^\delta f(t) \in L^\infty([0, T_1]; H_\ell^s(\mathbb{R}^3)),$$

where

$$M_{\lambda(t)}^\delta f(t) = \int e^{iv \cdot \xi} M_{\lambda(t)}^\delta(\xi) \varphi(t, \xi) \frac{d\xi}{2\pi^3}, \quad \varphi(t, \xi) = \int_{\mathbb{R}^3} e^{-iv \cdot \xi} dF_t(v).$$

When $s > 1/2$ we choose $1/2 < s' < s$ such that $\gamma + 2s' > 0, 2s' \geq (2s - 1)$. Choose $N_0 = a + (5 + \gamma + 2s' - 1)/2$ such that [4, (3.9)] is satisfied. Put $\varepsilon_0 = (\gamma + 1)/10 > 0$ and consider $\varepsilon = \varepsilon_0$. Then, we have

$$(5.12) \quad s + \lambda(T_1) - N_0 - a = s + \varepsilon_0 - N_0 \leq s - s' + 2\varepsilon_0 - (1 + \gamma)/2 = s - s' - 3\varepsilon_0.$$

Since we may assume $s - s' \leq \varepsilon_0$, (5.12) also shows (5.11).

The formula (5.11), together with [4, Theorem 3.6 and Proposition 3.8], leads us to the fact that, for any $t \in (0, T_1]$,

$$\begin{aligned}
(5.13) \quad & \frac{1}{2} \int_{\mathbb{R}^3} |M_{\lambda(t)}^\delta f(t)|^2 dv - N \int_0^t \int_{\mathbb{R}^3} \left| (\log \langle D \rangle)^{1/2} M_{\lambda(\tau)}^\delta f(\tau) \right|^2 dv d\tau \\
& = \frac{1}{2} \int_{\mathbb{R}^3} |M_{\lambda(0)}^\delta f_0|^2 dv \\
& + \int_0^t \left(Q(F_\tau, M_{\lambda(\tau)}^\delta f(\tau)), M_{\lambda(\tau)}^\delta f(\tau) \right)_{L^2} d\tau \\
& + \int_0^t \left(M_{\lambda(\tau)}^\delta Q(F_\tau, F_\tau) - Q(F_\tau, M_{\lambda(\tau)}^\delta f(\tau)), M_{\lambda(\tau)}^\delta f(\tau) \right)_{L^2} d\tau,
\end{aligned}$$

by the quite same manipulation as in the proof of [4, Lemma 4.3]. Indeed, similar as [4, (4.7)], from the definition (1.7) we start by the formula

$$\begin{aligned}
& \int_{\mathbb{R}^3} \psi(v) dF_t - \int_{\mathbb{R}^3} \psi(v) dF_{t'} = \frac{1}{2} \int_{t'}^t \int_{\mathbb{R}^3} \int_{\mathbb{R}^3} \int_{\mathbb{S}^2} b(\cdot) |v - v_*|^\gamma \\
& \quad \times (\psi(v'_*) + \psi(v') - \psi(v_*) - \psi(v)) dF_\tau(v) dF_\tau(v_*) d\sigma d\tau, \quad 0 \leq t' \leq t \leq T_1,
\end{aligned}$$

and set $\psi = (M_{\lambda(\bar{t})}^\delta)^2 f(\bar{t})$ with $\bar{t} = t, t'$. All other steps are carried on by regarding $F \in P_0(\mathbb{R}^3)$ as $F \in \mathcal{S}'(\mathbb{R}^3)$, in view of the pseudo-differential operator calculus.

Applying [4, Theorem 3.6] to the last term of (5.13), we have

$$\begin{aligned}
(5.14) \quad & \frac{1}{2} \| (M_{\lambda}^\delta f)(t) \|_{L^2}^2 \leq \frac{1}{2} \| f(0) \|_{H^a}^2 + \int_0^t \left(Q(F_\tau, (M_{\lambda}^\delta f)(\tau)), (M_{\lambda}^\delta f)(\tau) \right) d\tau \\
& + C_F \int_0^t \left(\| (M_{\lambda}^\delta f)(\tau) \|_{H_{\gamma+(2s-1)^+}^{s'}} + \| f(\tau) \|_{H_\gamma^s} \right) \| (M_{\lambda}^\delta f)(\tau) \|_{H^{s'}} d\tau \\
& + CN \int_0^t \| (\log \langle D \rangle)^{1/2} (M_{\lambda}^\delta f)(\tau) \|_{L^2}^2 d\tau.
\end{aligned}$$

Proposition 5.2, together with the interpolation in the Sobolev space, yields

$$\left(Q(F_\tau, (M_{\lambda}^\delta f)(\tau)), (M_{\lambda}^\delta f)(\tau) \right) \leq -c_F \| (M_{\lambda}^\delta f)(\tau) \|_{H_{\gamma/2}^s}^2 + C_F \| f(\tau) \|_{H_{\gamma/2}^{-2}}^2.$$

Use an interpolation inequality concerning weighted type Sobolev spaces with respect to variable v as follows: For any $k \in \mathbb{R}, p \in \mathbb{R}_+, \eta > 0$,

$$(5.15) \quad \| f \|_{H_p^k(\mathbb{R}_v^3)}^2 \leq C_\eta \| f \|_{H_{2p}^{k-\eta}(\mathbb{R}_v^3)} \| f \|_{H_0^{k+\eta}(\mathbb{R}_v^3)}, \text{ see for instance [11].}$$

Then we have, for a suitable large $\ell > 0$,

$$\| (M_{\lambda}^\delta f)(t) \|_{L^2}^2 + c \int_0^t \| (M_{\lambda}^\delta f)(\tau) \|_{H_{\gamma+/2}^s}^2 d\tau \leq \| f(0) \|_{H^a}^2 + C \int_0^t \| f(\tau) \|_{H_\ell^a}^2 d\tau.$$

Taking $\delta \rightarrow +0$ and $t = T_1$, we have $f(T_1) \in H^{\lambda(T_1)} = H^{\varepsilon_0+a}$ because $NT_1 = \varepsilon_0$. Repeating the same procedure with 0 and a replaced by T_1 and $a_1 = \varepsilon_0 + a$, we attain to (5.10), finite times, for any small $t_1 > 0$ since the above $T_1 > 0$ can be arbitrarily small by choosing a large N .

The case $\gamma < 0$ is similarly handled by means of (5.15), for the solution satisfying (1.9). We remark that the restriction $\gamma > -1$ comes from what we putted $\varepsilon_0 = (\gamma + 1)/10$ when $s > 1/2$.

Once we obtain (5.10), it follows from [4, Theorem 4.1] that, for any $0 < t_0 < T$, we have $f \in L^\infty([t_0, T]; \mathcal{S}(\mathbb{R}^3))$. Since, by means of [2, Proposition 2.9], for any $\ell > 0$ we have

$$\|Q(f, g)\|_{L_\ell^2} \lesssim \left(\|f\|_{L_{\ell+\gamma+2s}^1} + \|f\|_{L^2} \right) \|g\|_{H_{\ell+\gamma+2s}^{2s}},$$

the equation (1.1) and the Leibniz formula lead us to $F_t \in C^\infty([t_0, \infty); \mathcal{S}(\mathbb{R}^3))$. \square

Acknowledgements: The research of the first author was supported in part by Grant-in-Aid for Scientific Research No.25400160, Japan Society for the Promotion of Science. The research of the third author was supported in part by the General Research Fund of Hong Kong, CityU No. 11303614.

REFERENCES

- [1] R. Alexandre, L. Desvillettes, C. Villani and B. Wennberg, Entropy dissipation and long-range interactions, *Arch. Rational Mech. Anal.* **152** (2000), 327-355.
- [2] R. Alexandre, Y. Morimoto, S. Ukai, C.-J. Xu and T. Yang, Boltzmann equation without angular cutoff in the whole space: Qualitative properties of solutions, *Arch. Rational Mech. Anal.*, **202** (2011), 599-661.
- [3] R. Alexandre, Y. Morimoto, S. Ukai, C.-J. Xu and T. Yang, The Boltzmann equation without angular cutoff in the whole space: I, Global existence for soft potential, *J. Funct. Anal.* **262** (2012), 915-1010.
- [4] R. Alexandre, Y. Morimoto, S. Ukai, C.-J. Xu and T. Yang, Smoothing effect of weak solutions for the spatially homogeneous Boltzmann equation without angular cutoff, *Kyoto J. Math.* **52** (2012), 433-463.
- [5] A. V. Bobylev, The method of the Fourier transform in the theory of the Boltzmann equation for Maxwell molecules. *Dokl. Akad. Nauk SSSR* **225** (6) (1975), 1041-1044.
- [6] A. V. Bobylev, The theory of the nonlinear spatially uniform Boltzmann equation for Maxwell molecules. In *Mathematical physics reviews, Vol. 7*, volume 7 of *Soviet Sci. Rev. Sect. C Math. Phys. Rev.*, pages 111-233. Harwood Academic Publ., Chur, 1988.
- [7] M. Cannone and G. Karch, Infinite energy solutions to the homogeneous Boltzmann equation, *Comm. Pure Appl. Math.* **63** (2010), 747-778.
- [8] E. A. Carlen, E. Gabetta and G. Toscani, Propagation of smoothness and the rate of exponential convergence to equilibrium for a spatially homogeneous Maxwellian gas, *Comm. Math. Phys.*, **199** (1999) 521-546.
- [9] N. Fournier, Finiteness of entropy for the homogeneous Boltzmann equation with measure initial condition, *Ann. Appl. Probab.* **25** (2015), 860-897.
- [10] E. Gabetta, G. Toscani and B. Wennberg, Metrics for probability distributions and the trend to equilibrium for solutions of the Boltzmann equation, *J. Statist. Phys.*, **81**, 901-934.
- [11] Z.H. Huo, Y. Morimoto, S. Ukai, T. Yang, *Regularity of solutions for spatially homogeneous Boltzmann equation without Angular cutoff*. *Kinetic and Related Models*, **1** (2008) 453-489.
- [12] N. Jacob, *Pseudo-differential operators and Markov processes. Vol 1: Fourier analysis and semigroups*. Imperial College Press, London, 2001.
- [13] X. Lu and C. Mouhot, On measure solutions of the Boltzmann equation, part I: Moment production and stability estimates, *Jour. Diff. Equa.* **252**(2012), 3305–3363.
- [14] X. Lu and B. Wennberg, Solutions with increasing energy for the spatially homogeneous Boltzmann equation, *Nonlinear Anal. Real World Appl.* **3**(2002), 243-258.
- [15] S. Mischler and B. Wennberg, On the spatially homogeneous Boltzmann equation, *Ann. Inst. H. Poincaré Anal. Non Linéaire*, **16** (1999), 467-501.
- [16] Y. Morimoto, A remark on Cannone-Karch solutions to the homogeneous Boltzmann equation for Maxwellian molecules, *Kinetic and Related Models*, **5** (2012), 551-561.
- [17] Y. Morimoto, S. Ukai, C.-J. Xu and T. Yang, Regularity of solutions to the spatially homogeneous Boltzmann equation without angular cutoff, *Discrete and Continuous Dynamical Systems - Series A* **24** (2009), 187–212.

- [18] Y. Morimoto, S. Wang and T. Yang, A new characterization and global regularity of infinite energy solutions to the homogeneous Boltzmann equation, *J. Math. Pures Appl.* **103** (2015), 809–829.
- [19] Y. Morimoto, S. Wang and T. Yang, Moment classification of infinite energy solutions to the homogeneous Boltzmann equation, to appear in *Analysis and Applications(2015)* DOI: 10.1142/S0219530515500232, <http://arxiv.org/abs/1506.06493>.
- [20] Y. Morimoto and T. Yang, Smoothing effect of the homogeneous Boltzmann equation with measure valued initial datum, *Ann. Inst. H. Poincaré Anal. Non Linéaire*, **32** (2015), 429-442
- [21] A. Pulvirenti and G. Toscani, The theory of the nonlinear Boltzmann equation for Maxwell molecules in Fourier representation, *Ann. Mat. Pura Appl.* **171** (1996), 181-204.
- [22] M. Pulvirenti and B. Wennberg, Lower bounds for the solutions to the Kac and the Boltzmann equation. Proceedings of the Second International Workshop on Nonlinear Kinetic Theories and Mathematical Aspects of Hyperbolic Systems (Sanremo, 1994). *Transport Theory Statist. Phys.* 25(1996), 437-446.
- [23] M. Pulvirenti and B. Wennberg, A Maxwellian lower bound for solutions to the Boltzmann equation, *Comm. Math. Phys.* **183**(1997), 145-160.
- [24] G. Toscani and C. Villani, Probability metrics and uniqueness of the solution to the Boltzmann equations for Maxwell gas, *J. Statist. Phys.*, **94** (1999), 619-637.
- [25] C. Villani, On a new class of weak solutions to the spatially homogeneous Boltzmann and Landau equations, *Arch. Rational Mech. Anal.*, **143** (1998), 273–307.
- [26] C. Villani, A review of mathematical topics in collisional kinetic theory. In: Friedlander S., Serre D. (ed.), *Handbook of Fluid Mathematical Fluid Dynamics*, Elsevier Science (2002).
- [27] C. Villani, *Topics in optimal transportation*. Graduate Studies in Mathematics, **58**. American Mathematical Society, Providence, RI, (2003) .
- [28] X. Zhang and X. Zhang, Probability approaches to spatially homogeneous Boltzmann equations. *Stoch. Anal. Appl.* **25** (2007), no. 6, 1129-1150.

YOSHINORI MORIMOTO, GRADUATE SCHOOL OF HUMAN AND ENVIRONMENTAL STUDIES, KYOTO UNIVERSITY, KYOTO, 606-8501, JAPAN

E-mail address: morimoto@math.h.kyoto-u.ac.jp

SHUAIKUN WANG, DEPARTMENT OF MATHEMATICS, CITY UNIVERSITY OF HONG KONG, HONG KONG, P. R. CHINA

E-mail address: shuaiwang4-c@my.cityu.edu.hk

TONG YANG, DEPARTMENT OF MATHEMATICS, CITY UNIVERSITY OF HONG KONG, HONG KONG, P. R. CHINA; & DEPARTMENT OF MATHEMATICS, SHANGHAI JIAO TONG UNIVERSITY, SHANGHAI, P. R. CHINA

E-mail address: matyang@cityu.edu.hk